



Article Evaluating the Effects of the Rill Longitudinal Profile on Flow Resistance Law

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Abstract: In this paper, for the first time, the effect of the longitudinal profile shape of the rill (uniform, concave, and convex) on flow resistance law was studied. The first part of the paper is based on a theoretical equation to estimate the Darcy–Weisbach friction factor f, deduced from the power velocity distribution and rill measurements performed on a plot. At first, the equation to estimate the Γ parameter of the velocity profile was calibrated using all available measurements. Then an analysis of the hydraulic characteristics at reach scale, for comparable values of discharge, was carried out, comparing the different profile shapes. To assess the influence of the rill profile shape on flow resistance law, this calibration was also carried out using the data categorized by profile. In the second part of the paper, an analysis of the scour depth and eroded rill volume was developed for four rills of each configuration. The results showed that an accurate estimate of f can be obtained by calibrating the flow resistance equation for each profile shape. The component of the Darcy-Weisbach friction factor that is due to the profile shape varied from 0.68 to 14.6% of the overall friction factor for the concave profile, and from 3.4 to 26.9% for the convex profile. The analysis also showed that, for the convex profile, the scour was concentrated downstream of the slope change, while for the uniform and concave profiles, it was uniformly distributed. Furthermore, the scour depth measured in correspondence of the rill thalweg had a generally increasing trend, with discharge for all the investigated profile shapes. Total eroded rill volume of the concave profile was lower than those detected for the uniform and convex profiles and was characterized by a reduction of 57.9%, as compared to the uniform profile.

Keywords: slope profile; concave profile; convex profile; rill erosion; rill hydraulics

1. Introduction

Many researchers have studied the effects of length and slope gradient on runoff and soil loss using uniform plots, i.e., having a constant slope, while a small number of studies have focused on the effect of profile shape (concave, convex) [1,2].

Young and Mutchler [3,4] investigated experimentally the effect of slope shape (convex, concave, uniform) on runoff and soil loss at plot scale and showed that concave hillslopes tends to have reduced total sediment loss compared to a uniform profile slope.

Rieke-Zapp and Nearing [2] carried out laboratory experiments to determine the relationship between slope shapes and soil erosion using five slope shape treatments (convex-linear, concave-linear, nose slope, head slope and uniform). These experiments demonstrated that rill patterns, sediment yield and runoff, all changed with slope shape. In particular, these authors reported a decrease of 75% of soil loss for concave slopes, as compared to uniform ones with similar surface area.

Sensoy and Kara [5] carried out an experimental investigation using nine experimental field plots with different profile shapes, established on a 30% hillslope, subjected to natural rainfall. This field experiment indicated that slope shape (uniform, concave, convex)



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Copyright: © 2022 by the authors. Licensee MDPI, Basel, Switzerland. This article is an open access article distributed under the terms and conditions of the Creative Commons Attribution (CC BY) license (https:// creativecommons.org/licenses/by/ 4.0/). affected both runoff and soil loss. In particular, the uniform slope was characterized by the highest values of runoff and soil loss, while the concave slope produced the lowest values. The experimental results also indicated that the soil eroded from the uniform plots was mainly composed of fine particles (<2 mm), while both the concave and convex slopes were characterized by a large size variability of the eroded particles.

Jeldes et al. [6] compared concave and uniform hillslopes having the same degree of mechanical stability and concluded that concave slopes produced 15–40% less sediment than uniform slopes. This value was lower than that reported in the literature, which varied from 50% [7] to 80% [8].

Mombini et al. [9] studied the effect of surface roughness on soil loss using different complex hillslopes in terms of plan shapes (convergent, parallel, and divergent) and profiles (convex, concave and uniform). This research, carried out in laboratory conditions using three different soil surface roughness values and a single simulated rainfall intensity, indicated that the highest soil loss due to soil roughness was observed on uniform parallel hillslopes.

Despite the current availability of techniques able to produce a more accurate and detailed representation of the hillslope surface [10–13], the rill erosion processes on non-linear profiles (concave, convex) have been barely studied by the scientific community.

Rill erosion results from soil particle detachment and transport by channelized flows. When overland flow tends to convert into a concentrated flow, the formation of rills occurs. Rills are eroded channels, characterized by rapid morphological evolution, that represent a relevant sediment source and a transport path for eroded particles from hillslopes [14,15]. The increased flow shear stress and velocity within the rill increase sediment yield promoting the transport of soil particles that are detached by rill flow and delivered from the interrill areas to the rill channels [16–23]. Rill erosion has both on-site (reduction of soil fertility and productivity) and off-site (river over-sedimentation, reservoir sedimentation and degradation of water quality) impacts [24] and is dominant as compared to interrill erosion [21].

In previous papers, a theoretical approach based on the integration of a power-velocity profile, was proposed to deduce the Darcy–Weisbach friction factor [25,26], and its applicability was tested for rill flows [27–35] on uniform plots with different slope gradients.

In this paper, for the first time, the theoretical approach by Ferro [25,26] is applied to rill measurements carried out on three plots having different profiles (concave, convex and uniform).

The specific aims of this study are to: (a) test the reliability of the abovementioned theoretical approach; (b) estimate the component of Darcy–Weisbach friction factor corresponding to the profile shape; (c) investigate the evolution of the rill thalweg for the three examined profile shapes and (d) compare the volumes eroded by flow with similar discharges for the different profiles.

2. Materials and Methods

2.1. The Experimental Plots and Measurements

A plot, 2 m wide and 7 m long (Figure 1), located in the experimental area of the Department of Agriculture, Food and Forest Sciences of the University of Palermo, was used to carry out the experiments.



Figure 1. View of the experimental plot with uniform (a), convex (b), and concave (c) profiles.

The plot was prepared with three different profiles (concave, convex and uniform) (Figure 1) using a mean plot slope s_p of 18%. In particular, for the concave and convex profiles, the plot was divided into two 3.27 m long parts, and two slope values (12 and 24%) were set up to obtain a mean value of 18% (Figure 2). The soil was a clay loam characterized by 32.7%, 30.9% and 36.4% of clay, silt, and sand, respectively.



Figure 2. Scheme of the investigated profiles.

For all the profile shapes, soil was manually hoed and raked and then compacted by a woody bar. The bulk density of the soil was determined by sampling the soil, using a metal cylinder of known volume, and determining its weight after drying. A uniform soil moisture condition for all the experiments was achieved by wetting the soil until infiltration stopped. Then, rills were manually incised along the maximum slope direction of the plot and were shaped using a clear flow discharge of 0.1 L s^{-1} , applied for 3 min. The rills had the same longitudinal profile as the plot. For each rill, nine longitudinal segments, bounded by two cross-sections spaced 0.624 m apart (Figure 3), were determined.



Figure 3. View of the plot area with some tested rills (**a**), scheme of cross-sections (**b**), and examples of rill reaches (**c**).

A constant clear inflow discharge, ranging from 0.18 to 0.63 L s⁻¹, was used for the experiments.

Measurements of water depth h, cross-section area σ , wetted perimeter C, bed slope s, and flow velocity V were carried out for rill reaches (i.e., the channel between a given measured cross-section and the rill end) on uniform (99 reaches), concave (71 reaches), and convex (72 reaches) rills. The indirect measurements of Reynolds number $Re = V h/v_k$, in which v_k is the kinematic viscosity, Froude number $F = V/\sqrt{(g h)}$, in which g is the acceleration due to gravity, and Darcy–Weisbach friction factor $f = 8 gRs /V^2$, where R is the hydraulic radius, were also obtained.

For each experimental run, a 3D-DTM of the plot area was built from a set of 70 photographs using a 3D-photo reconstruction technique with the image-processing software Agisoft Photoscan Professional (Version 1.1.6, Agisoft, St. Petersburg, Russia). The slope gradient of each rill segment was determined using the rill thalweg imaged by the 3D-DTM. The slope of each rill reach was calculated by averaging the values measured in the considered reach.

The mean water depth h and the mean hydraulic radius R in the reach were determined using the method developed by Di Stefano et al. [36]. This method couples an accurate ground survey of the rill channel, obtained by close-range photogrammetry, with the survey of the water tracks inside the channel marked by a dye solution. The measured water depths and geometric cross-sections obtained by the 3D-DTM, at a distance interval of 6.2 cm within the reach, were used to calculate the values of the hydraulic cross-section area σ and wetted perimeter *C* in the reach and the values of h and *R* using the following equations [36]:

$$h = \frac{\sum_{i=2}^{N-1} \left(\frac{\sigma_1}{2} + \sigma_i + \frac{\sigma_N}{2}\right)}{\sum_{i=2}^{N-1} \left(\frac{w_1}{2} + w_i + \frac{w_N}{2}\right)}$$
(1)

$$R = \frac{\sum_{i=2}^{N-1} \left(\frac{\sigma_1}{2} + \sigma_i + \frac{\sigma_N}{2}\right)}{\sum_{i=2}^{N-1} \left(\frac{C_1}{2} + C_i + \frac{C_N}{2}\right)}$$
(2)

in which *w* is the surface width and *N* is the number of the cross-sections in the reach.

A methylene blue solution was applied as a dye tracer to measure the flow velocity [28,37-40]. The measured surface velocity was converted to the mean flow velocity *V* by a correction factor of 0.8 [41–43].

For four rills of each profile shape, Digital Elevation Models (DEMs) at the ends of the shaping phase (*D*1) and the experimental runs (*D*2) were obtained from the 3D-DTM. The DEM of the difference (DoD) between *D*2 and *D*1 was used to measure the scour depth values inside the whole rill channel and calculate the total eroded rill volume *RV*. Considering that the minimum level of detection of the used DEMs was ± 0.005 m, variations of bed elevation in the range ± 0.005 m were not detectable.

Finally, for the same rills, the scour depth values *SD* were measured, at 0.005 m intervals along the rill thalweg, as the difference in height between the longitudinal profiles at the end of the shaping phase (*D1*) and at the end of the experimental runs (*D2*) (an example for one rill with a convex profile shape is plotted in Figure 4). For each rill, the mean scour depth SD_m was also calculated as the mean of the *SD* values.



Figure 4. Comparison, using a rill with a convex profile, between the rill thalweg at the end of the shaping phase (D1) and at the end of the experimental run (D2).

2.2. The Rill Flow Resistance Equation

For an open channel flow, the local flow velocity profile v(y) along a given vertical is represented by the following functional relationship [44–46]:

$$\varphi\left(\frac{dv}{dy}, y, h, d, u_*, s, \rho, \mu, g\right) = 0 \tag{3}$$

where φ is a functional symbol, v is the local velocity, y is the distance from the bottom, d is a characteristic bed particle diameter, $u_* = \sqrt{(g R s)}$ is the shear velocity, ρ is the water density and μ is the dynamic water viscosity.

$$\frac{y}{u_*}\frac{dv}{dy} = \varphi_1\left(\frac{h}{d}, s, \frac{u_*y}{v_k}, Re, F\right)$$
(4)

in which φ_1 is a functional symbol.

Hypothesizing the Incomplete Self-Similarity in u_*y/v_k [25,47–50], neglecting the influence of *Re* (turbulent flow regime) and considering that the flow Froude number accounts also for the ratio h/d [26], the integration of Equation (4) yields to the following velocity profile:

$$\frac{v}{u_*} = \Gamma(s, F) \left(\frac{u_* \ y}{v_k}\right)^{\delta} \tag{5}$$

where $\Gamma(s, F)$ is a function to be defined by velocity measurements and the exponent δ is calculated by the following theoretical equation [51,52]

$$\delta = \frac{1.5}{\ln Re} \tag{6}$$

The following expression of the Darcy–Weisbach friction factor f is deduced [25,48,53] by integrating Equation (5):

$$f = 8 \left[\frac{2^{1-\delta} \Gamma R e^{\delta}}{(\delta+1) (\delta+2)} \right]^{-2/(1+\delta)}$$
(7)

From the velocity profile (Equation (5)) the following estimate Γ_v of Γ function [25,53] is obtained by setting equal to α *h* the distance *y* at which the local velocity is equal to the cross-section average velocity *V*:

$$\Gamma_{v} = \frac{V}{u_{*} \left(\frac{u_{*} \alpha h}{v_{k}}\right)^{\delta}}$$
(8)

The coefficient α , which is less than 1, takes into account that both *V* is located below the water surface and the mean velocity profile in the cross-section is considered. The coefficient α has to be calculated by the following theoretical relationship deduced by Ferro [25]:

$$\alpha = \left[\frac{2^{1-\delta}}{(\delta+1)(\delta+2)}\right]^{1/\delta} \tag{9}$$

For this investigation α varied in a narrow range and the mean value of 0.124 was assumed. Considering that, according to Equation (5), Γ theoretically depends only on *s* and *F* [26,54], Γ_v can be estimated using the following power equation:

$$\Gamma_v = a \, \frac{F^b}{s^c} \tag{10}$$

where *a*, *b* and *c* are coefficients to be determined from experimental measurements.

3. Results

3.1. Rill Flow Resistance for Different Profile Shapes

The experimental runs were carried out for Reynolds numbers corresponding to turbulent flow (2764 $\leq Re \leq 11015$ for the uniform profile, 2282 $\leq Re \leq 14346$ for the concave one, and 2486 $\leq Re \leq 9487$ for the convex one) and for Froude number values corresponding to subcritical and supercritical flow conditions ($0.54 \leq F \leq 1.70$ for the uniform profile, $0.64 \leq F \leq 1.63$ for the concave one, and $0.44 \leq F \leq 1.71$ for the convex

one) (Figure 5). In other words, the investigated ranges of *Re* and *F* are very similar for the three different profiles.



Figure 5. Experimental values of Froude and Reynolds number for each profile shape.

Using all the 242 measurements, i.e., neglecting the influence of the profile shape, Equation (10) was calibrated obtaining the following relationship:

$$\Gamma_v = 0.4322 \, \frac{F^{1.0857}}{s^{0.5235}} \tag{11}$$

characterized by a coefficient of determination equal to 0.9894. The comparison between the 242 measured Γ_v values, obtained by Equation (8) and those calculated applying Equation (11) is plotted in Figure 6a.

Introducing Equation (11) into Equation (7), the following flow resistance law was obtained:

$$f = 8 \left[\frac{(\delta+1) (\delta+2)}{2^{1-\delta} Re^{\delta} 0.4322} \frac{s^{0.5235}}{F^{1.0857}} \right]^{2/(1+\delta)}$$
(12)

A good agreement, characterized by a root mean square error *RMSE* equal to 0.1743, between the measured Darcy–Weisbach friction factor values and those calculated by Equation (12) was detected (Figure 6b), even if Equation (12) tended to slightly overestimate the *f* values for the uniform profile. The friction factor values calculated by Equation (12) were characterized by errors less than or equal to $\pm 20\%$ for 96.3% of cases and less than or equal to $\pm 10\%$ for 78.5% of cases. Figure 6 highlights that the *f* measurements obtained for the convex profile covered a different range than those obtained from the uniform and concave profiles. Different characteristics of the convex profile are also revealed by the frequency distributions of the hydraulic variables *F*, *f*, *s* and *V* plotted in Figure 7.





As this analysis showed that profile shape influences flow characteristics, the theoretical flow resistance law was calibrated for each investigated condition. Table 1 lists the a, band c coefficients of Equation (10) obtained for the uniform, concave, and convex profiles.

Table 1. Values of *a*, *b* and *c* coefficients of Equation (10) obtained for the uniform, concave, and convex profiles.

| Profile Shape | а | b | С |
|---------------|--------|--------|--------|
| Uniform | 0.4419 | 1.0127 | 0.5224 |
| Concave | 0.4025 | 1.1084 | 0.5469 |
| Convex | 0.5187 | 1.0437 | 0.3795 |



Figure 7. Frequency distributions of Froude number (**a**), friction factor (**b**), reach slope (**c**) and flow velocity (**d**).

Figure 8a shows the comparison between the measured Γ_v values obtained by Equation (8) and those calculated applying Equation (10), with the specific *a*, *b* and *c* coefficients (Table 1) for the different profiles. The coefficient of determination of the obtained relationships is equal to 0.9812, 0.9979 and 0.9882 for the uniform, concave, and convex profiles, respectively.

Figure 8b shows good agreement, characterized by a RMSE equal to 0.1234, between the measured values of the Darcy–Weisbach friction factor and those calculated by Equations (7) and (10), with *a*, *b* and *c* coefficients obtained for each profile shape. In this case, the calculated friction factor values are characterized by errors always less than or equal to $\pm 20\%$ and less than or equal to $\pm 10\%$ for 86.8% of cases.



Figure 8. Comparison between the measured Γ_v values obtained by Equation (8) and those calculated by Equation (10) with the shape profile-specific *a*, *b* and *c* coefficients (**a**) and measured Darcy–Weisbach friction factor values and those calculated by Equations (7) and (10) with the shape profile-specific *a*, *b* and *c* coefficients (**b**).

Moreover, to study the influence of profile shape on rill flow resistance, the additive method of shear stresses [55], already used to investigate the effects of sediment transport for rills [29,56], and of riparian vegetation for a vegetated channel [57], was applied. In particular, the Darcy–Weisbach friction factor, f_{mor} , corresponding to the effect of the profile on flow resistance, was calculated as the difference f- f_u between the measured f value of the concave or convex profile and that of, f_u , calculated for the uniform profile in the same hydraulic conditions (Equations (7) and (10), with a, b and c coefficients obtained for the uniform profile). This f_{mor} value is almost equal to zero for 9.9% (concave) and 4.2% (convex) of the measured values. The ratio f_{mor}/f varies from 0.68 to 14.6% for concave profile and from 3.4 to 26.9% for convex profile (Figure 9).



Figure 9. Frequency distribution of the ratio f_{mor}/f .

3.2. Analysis of Rill Scour Depth and Eroded Volume for Different Profile Shapes

Figure 10 shows the DoD obtained for the uniform (a), concave (b) and convex (c) profiles were characterized by comparable ranges of flow discharge ($0.26-0.58 \text{ L s}^{-1}$ for uniform, $0.18-0.46 \text{ L s}^{-1}$ for concave, and $0.21-0.60 \text{ L s}^{-1}$ for convex profiles). This figure clearly shows that the concave profile is characterized by scour depth values (ranging from 0.005 to 0.08 m) lower than those of uniform (0.005-0.153 m) and convex (0.005-0.176 m) profiles.

For all three investigated profile shapes, Figure 10 also shows that rills are generally scoured by flow. The blue areas of this figure identify rill zones where the variation of bed elevation is lower than the minimum level of detection (± 0.005 m), or lack variation. Moreover, for the convex profile, the scour is concentrated downstream of the slope change, while for uniform and concave profiles, it is uniformly distributed.

The knowledge of the scour depth values obtained by the DoDs allows calculation of the total eroded volume RV. Figure 11 shows the comparison between the applied discharge Q and RV using an example of four rills analysed for each profile. The evident trend, as expected, is that RV increases with Q for all the investigated profiles. Finally, for the uniform, concave, and convex profiles, the sum of the four RV values resulted in 0.0644, 0.0271 and 0.0654 m³, respectively.

For the forementioned four experimental runs for each profile shape, the frequency distribution of the scour depth values *SD* was plotted in Figure 12. This figure highlights that scour depth, measured in corresponding rill thalwegs, has a generally increasing trend with discharge for all investigated profile shapes.



Figure 10. DoD obtained for the uniform (a), concave (b) and convex (c) profiles.

Figure 13 shows the frequency distribution of the SD/SD_m ratio for the investigated discharges. This figure demonstrates that the variable SD/SD_m can be considered independent of discharge.

Figure 14 shows the empirical frequency distribution of SD/SD_m for each profile shape, using an example for a single discharge value ($\approx 0.35 \text{ L s}^{-1}$). The overlapping of the three distributions suggested that SD/SD_m does not depend on profile shape.



Figure 11. An example of the comparison between the applied discharge *Q* and the total eroded rill volume *RV*, for four rill channels for each profile shape.



Figure 12. An example of the frequency distribution of the scour depth *SD* for four discharge values, for uniform (**a**), concave (**b**), and convex (**c**) profiles.



Figure 13. An example of the frequency distribution of the SD/SD_m ratio, for four discharge values, for uniform (**a**), concave (**b**), and convex (**c**) profiles.



Figure 14. Frequency distribution of the SD/SD_m ratio for a fixed value of discharge ($\approx 0.35 \text{ L s}^{-1}$).

The relationship between SD_m and the discharge was determined for each profile shape (Figure 15). The mean scour depth increased with discharge according to a power relationship depending on the profile.



Figure 15. Relationship between mean scour depth SD_m and discharge for each profile shape.

4. Discussion

4.1. Rill Flow Resistance for Different Profile Shapes

Figure 5 shows that the rill flows for the three examined profiles were characterized by comparable Reynolds numbers, while for the convex profile, they were characterized, on average, by Froude numbers lower than those related to the uniform and concave profiles, and less than 1 for 93% of cases. For the latter, the pairs (Re, F) were generally overlapped. The detected trend for Froude number confirmed that the profile shape influences the kinematic flow characteristics (Figure 7a), while the comparable Reynolds number values (Figure 5) suggest that the profile does not affect the flow regime.

Figure 7 confirmed this overlapping trend for the measured flow velocity (Figure 7d) and Froude number (Figure 7a) on uniform and concave profiles. The slope values of the concave profile were less than those of the uniform profile, while the *s* values of the convex profile were greater than those measured on the uniform profile (Figure 7c). For the convex profile, the occurrence of the highest slope values and the lowest velocity values can be explained by the relevant erosion phenomena. Indeed, the highest *RV* values detected for the convex profile at the single rill and plot scale, demonstrate that for this profile, the erosion phenomena determine a major increase of grain roughness and sediment transport that lower flow velocity and increase the Darcy–Weisbach friction factor (Figure 7b). The overlapping values of flow velocity and friction factor *f* for uniform and concave profiles can be explained by the similar distribution of erosion phenomena occurring in the rill channels (Figure 10a,b).

The comparison between Figures 6b and 8b and the error statistics highlight that a single flow resistance equation (Equation (12)) is applicable to all investigated profiles, despite the differences highlighted by the analysis of hydraulic characteristics at reach scale. However, the f estimate is more accurate if specific calibration coefficients (a, b and c) are applied for each profile.

Figure 9 demonstrates that the ratio f_{mor}/f was greater than 10% for 39.1% and 69.6% of the investigated cases for the concave and convex profiles, respectively. Thus, for the examined conditions, the effect of profile shape on flow resistance can be considered relevant.

The convex profile was characterized by the highest values of the ratio f_{mor}/f (3.4–26.9%) (Figure 9), which demonstrates that the contribution of the component of total flow resistance due to the profile was more pronounced than for the concave profile. This result can be explained as higher eroded volumes were detected for the convex profile and f_{mor} accounts for sediment transport and roughness effects on the flow resistance.

4.2. Analysis of Rill Scour Depth and Eroded Volume for Different Profile Shapes

Previous studies [23,58] considered a slope value of 18% as a threshold to discriminate between the gentle (GS) and the steep slope (SS) condition. This distinction is explained as differences in hydraulic (flow depth, velocity, Reynolds number, Froude number) and sediment transport (flow transport capacity, actual sediment load) variables occur. For the SS condition, results by Nicosia et al. [58] supported the conclusion that the sediment transport capacity T_c was so high that it was not a limiting factor for the actual sediment transport, while for the GS case, the actual sediment transport could be limited by both T_c and the soil particle detachability and transportability.

For the convex profile, Figure 10c shows that the scour was concentrated, while it was evenly distributed for uniform (Figure 10a) and concave profiles (Figure 10b). In particular, the noticeable scour of the rills in convex profiles was located downstream of the slope change from a gentle (12%) to a steep (24%) slope condition, because in the upstream half of the rill, the actual sediment transport was limited by T_c , while in the downstream half, it was not. Notwithstanding, the spatial distribution of the scour was similar for uniform and concave profiles, the scour depth values of concave profiles were lower than those of uniform and convex profiles, which instead were similar (Figure 10). Consequently, also the eroded volume of the rill for the concave profile was lower than the eroded volumes for the uniform and convex profiles (Figure 11). In particular, the concave profile resulted in 57.9% less soil loss than the uniform profile. This result confirms the literature findings [2,6–8], which have reported 50 to 80% less soil loss from concave hillslopes than from uniform slopes.

Figure 11 also shows that increasing values of flow discharge determine a higher *RV*. This result agrees with that obtained by Di Stefano et al. [56] who found that when discharge increased, the contribution of sediment transport to flow resistance and the eroded rill volume also increased.

The results shown in Figures 13 and 14 suggest that SD/SD_m was independent of discharge and profile shape, respectively. In other words, the mean value SD_m accounted for the effects of both discharge and profile shape on scour depth.

Finally, Figure 15 demonstrates that profile shape determined a different increasing relationship between mean scour depth along the rill thalweg and discharge. The mean scour depth for the uniform profile was always higher than that obtained for the convex profile. Instead, the total eroded volume obtained for the uniform profile was slightly lower than that obtained for the convex profile. This result can be explained as the DoD information accounts for the spatial distribution of scour depth, while SD_m represents the mean value of measurements performed along a single linear flow path.

5. Conclusions

In this paper, the effect of the rill profile on flow resistance law was evaluated using measurements obtained from uniform, concave, and convex profiles.

At first, the equation to estimate the velocity profile parameter Γ was calibrated by all the available measurements and an accurate estimate of the Darcy–Weisbach friction factor (Equation (12)) was obtained. Then, to assess the influence of each profile on flow resistance law, the equation to estimate Γ was also calibrated when the data for each profile shape was distinguished. This choice determined an appreciable improvement of the friction factor estimate.

The Darcy–Weisbach friction factor corresponding to the effect of the profile on flow resistance, varied from 0.68 to 14.6% of the overall f factor for the concave profile, and from 3.4 to 26.9% of f for the convex profile.

The developed analysis also showed that for convex profiles, scour was concentrated downstream of the slope change, while for uniform and concave profiles, it was uniformly distributed. Furthermore, the scour depth measured in the corresponding rill thalwegs, had a generally increasing trend with discharge for all investigated profiles. The total eroded volume for the concave profile was lower than those obtained for uniform and convex profiles and was characterized by being 57.9% less than for the uniform profile.

Finally, the analysis indicated that the mean value of scour depth in the rill thalweg accounted for the effects of both discharge and profile shape, and increases with increased discharge according to a relationship varying with the profile shape.

The main limitations of the study are the use of a single soil type and one mean plot slope. These can be overcome by other experiments using different soils and values of mean plot slope.

The present investigation highlights a markedly less soil loss from concave slopes than from either uniform or convex slopes. Moreover, the proposed approach allows estimating rill flow velocity, which is an essential component of process-based soil erosion models.

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