



27 **Abstract**

28 Wind energy is extensively invested as a renewable and clean energy source. Savonius wind  
29 rotor, as an energy converter, has the merit of being appropriate for specific applications. The  
30 current paper centers around the performance enhancement of a helical Savonius wind rotor  
31 through blade shape optimization. The basic objective is to assess numerically and  
32 experimentally, in addition to the helical Savonius rotor, two novel rotor blade shapes named  
33 delta bladed and two-stage delta bladed shape. Numerical study was conducted using Ansys  
34 Fluent software. 3-D unsteady simulations were carried out through the use of the SST  $k-\omega$   
35 turbulence model based on the finite volume method solver. Aerodynamic flow and  
36 performance characteristics were investigated. Static and dynamic experimental tests were  
37 undertaken on a wind tunnel. An improvement in  $C_p$  by 22.58 % and 29.5 % for the two-stage  
38 delta bladed rotor and 19.35 % and 16.4 % for the delta bladed rotor over the helical Savonius  
39 rotor was recorded, respectively, numerically and experimentally. In addition, the self-  
40 starting ability as well as the aerodynamic flow characteristics of the helical rotor were  
41 enhanced with the novel blade shapes.

42 The obtained findings can be regarded as technologically significant in terms of opening  
43 further fruitful lines of investigations and promising future research direction as far as the use  
44 of helical Savonius wind turbine is concerned.

45

46 **Keywords:** Helical Savonius rotor, delta bladed shape, two-stage delta bladed shape,  
47 performance characteristics, aerodynamic flow characteristics.

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## 52 **1. Introduction**

53 The global energy demand is growing, while the fossil fuel sources are likely to run out. As a  
54 result of the extravagant use of the conventional fossil sources, the world has also been facing  
55 environmental threats related to the emission of a huge amount of pollutants. To overcome  
56 these problems, wind energy technology as an alternative to fossil fuels, has been a hot area of  
57 research triggering significant scientific concern and whetting the interest among both  
58 scientists and industrials.

59 As an inexhaustible and a friendly environmental source of energy, wind energy has been  
60 considered among the most outstanding renewable energy used to meet domestic and  
61 industrial demands. Wind turbines, defined as the wind energy converter, are basically  
62 classified into two classes: (HAWT) horizontal axis wind turbines and (VAWT) vertical axis  
63 wind turbines. This classification relies notably on the alignment between the coming wind  
64 direction and the turbine rotation axis. Within the wind energy, the HAWT have been the  
65 most commonly used for electricity requirements over the recent years. Compared to the  
66 VAWT, the horizontal ones display higher power efficiency but they require permanent wind  
67 direction to be more powerful. However, the vertical axis wind turbines have inherent  
68 specifications. As they operate independently from wind direction with a lower cost of  
69 installation and maintenance and without a need for yaw mechanism, VAWT are required in  
70 specific implementations. The operating principle rests on the momentum force difference  
71 that turns into a pressure difference between the convex and concave side of the rotor blades  
72 triggering subsequently the turbine rotation around a vertical shaft. VAWT, based on the  
73 driven force, are categorized into the lift force-based Darrieus rotor and the drag force-based  
74 Savonius rotor<sup>1</sup>. Despite having a lower power efficiency compared to HAWT and Darrieus  
75 turbines, the Savonius rotors have significant merits. As a drag machine, it operates thanks to  
76 the drag force resulting from its blades. It is to be noted that it also adds the lift force that

77 contributes to an amount of power production. The Savonius wind turbine can further operate  
78 under turbulent flow, independently from the wind direction and with a good self-starting  
79 ability<sup>2</sup>. Additionally, it has a simple structure, a lower cost of fabrication and maintenance, a  
80 lower noise emission and maintains the most stable performances. Two blades having a semi-  
81 cylindrical shape mounted from either side of a vertical shaft make up its original design,  
82 which creates an outlook that seems analogous to an “S” shape. This concept is meant to  
83 instill the incoming wind to the turbine blades.

84 The overlap ratio, the aspect ratio, the number of stages and blades, the end plates, and the  
85 blade shape stand for the main geometrical parameters influencing the conventional Savonius  
86 wind turbine performance. Its efficiency can equally be affected by adding accessories.  
87 Several investigations have reported both numerically and experimentally that the Savonius  
88 rotor performance has been improved across parametric optimization. With regard to adding  
89 accessories, the principle relies on harvesting the wind to the concave side of the advancing  
90 blade and precluding the convex side of the returning blade. In this respect, numerous designs  
91 were investigated. In fact, it was about deflecting plates installed upstream the rotor composed  
92 of one<sup>3</sup>, two<sup>4,5,6</sup> or even a set of plates<sup>7</sup>, a whole deflector system<sup>8,9</sup>, a curtain arrangement<sup>10</sup> as  
93 well as a guiding box<sup>11</sup>. It was found that the output performances mainly the torque and the  
94 power coefficients as well as the aerodynamic characteristics of the Savonius rotor were well  
95 improved with the use of those accessories. Regarding the optimization of the rotor  
96 geometrical parameters, the overlap ratio was examined. Numerically, Hassanzadeh et al.<sup>12</sup>  
97 recorded inward and outward overlap ratios and deduced that 0.2 was the optimum one. Patel  
98 et al.<sup>13</sup> investigated numerically negative overlap ratios varying from 0 to 0.4 and revealed that  
99 the zero one provided the maximum power coefficient equal to 0.09 at a tip speed ratio of 0.6.

100 As far as the aspect ratio is concerned, it was inferred that its rise enhances the torque  
101 coefficient, although it lowers the rotational speed since the structure becomes heavier<sup>14</sup>.

102 That's why, for structural constraints, it is set near to 1.0 for most of the extant rotors<sup>15</sup>.  
103 Additionally, using end plates was equally regarded as a way to boost the goodness of fit of  
104 Savonius rotor. Within the same vein, Jeon et al.<sup>16</sup> experimented Savonius rotor with different  
105 configurations of lower and upper end plates. A progress of 36% in the power coefficient was  
106 reported. In fact, the inclusion of end plates stop the air escaping from the concave side of  
107 both blades, forcing it to flow from the concave advancing blade to the concave returning  
108 blade. It maintains a greater pressure difference between the concave and convex sides of the  
109 rotor blades along the rotor height. As a result, it minimizes the negative drag and improves  
110 the energy conversion process. Regarding the number of blades, the rotor performance drops  
111 with its increase due to an important vortex phenomenon created close to the rotor blades<sup>17</sup>. In  
112 the same context, Mahmoud et al.<sup>14</sup> found that the two-blade Savonius rotor highlighted the  
113 highest value of power and torque coefficients compared to the three and four blades rotor.  
114 However, experimentally, Wenehenubun et al.<sup>18</sup> compared two, three and four blades rotors.  
115 They emphasized that the four-blade and the three-blade rotors are the most efficient,  
116 respectively, at a low and high speed ratio. Referring to the number of stages, Shamsuddin et  
117 al.<sup>19</sup> asserted that the double-stage rotor performance exceeded the single stage. Furthermore,  
118 the maximum power coefficient recorded was 138 % higher at 5 m/s and the self starting  
119 ability reached 5m/s while it was equal to 7 m/s with the single stage rotor. As the blade  
120 represents the main part of the Savonius wind turbine having the major impact on wind-  
121 capture capability, diligent efforts have been dedicated to optimize it to obtain better  
122 efficiency. Alom et al.<sup>20</sup> identified four blade profile configurations, named respectively,  
123 semi-circular, Benech, modified-Bach and elliptical then subjugated them into numerical and  
124 experimental tests. The maximum power coefficient was numerically equal to 0.27, 0.29, 0.30  
125 and 0.34, respectively. Hence, the elliptical profile displayed the highest performance.  
126 Talukdar et al.<sup>21</sup> focused on the elliptical form and developed a new design. They carried out

127 experiments in a subsonic wind tunnel and they revealed a 37% improvement over the  
128 traditional semicircular form. In the same vein, Al Absi et al.<sup>22</sup> explored the blade ellipticity  
129 through firstly changing the rotor overlap ratio then modifying the inner blade surface seeing  
130 from the concave part into a zigzag regular form. Based on wind tunnel tests at a wind  
131 velocity of 9m/s, the optimum modified elliptical blade form provided a power coefficient of  
132 0.29 which proved to be higher than the classical Savonius rotor (0.26). Asadi et al.<sup>23</sup> putted  
133 to test, numerically, two forms of Savonius wind rotor blades: the semi circular and the Bach-  
134 type. They proved that the rotor with the Bach-type form outperformed the traditional  
135 semicircular one with a maximum betterment of 20 % in the power coefficient at a tip speed  
136 ratio equal to 1.5. On the other side, Jia et al.<sup>24</sup> handled the blade curve. They set forward a  
137 new-developed two-blade Savonius rotor using the cubic Bezier curve method. The optimized  
138 blade was checked and compared to the classical one. It highlighted twofold benefits. It  
139 improved the power coefficient by 6.87% on the one hand and boosted the aerodynamic  
140 performance on the other hand. In addition, sine and conical blade forms were suggested by  
141 Abdelaziz et al.<sup>25</sup> and then compared to the classical blade form numerically with the use of 3-  
142 D simulations. The maximum power coefficient reduced by 1.25% with the sine form and  
143 enhanced by 8.6% with the conical form. Not too far, a blade shape inspired by sandeels fish  
144 structure was proposed by Abdelghafar et al.<sup>26</sup> The design form was described through the  
145 blade length, its maximum camber and its location. The optimization procedure was based on  
146 evolutionary algorithms. The optimum design exhibited an important increase in the rotor  
147 output up to 9.21 % at a tip speed ratio of 0.8. Within the same vein, modifying the twist  
148 angle was adopted as a tool to examine the change of the classical form of the Savonius rotor  
149 blades into helical one with twisted blades. Lee et al.<sup>27</sup> tested numerically and experimentally  
150 four helical Savonius rotors with 0°, 45°, 90° and 135° twist angles. The maximum power  
151 coefficient was recorded for 45° twist and dropped by 25.5% for 90° and 135°. Saad et al.<sup>28</sup>

152 explored five twist angles from  $0^\circ$  to  $180^\circ$  in step of  $45^\circ$ . With an increase in twist angle, the  
153 created torque took positive values for different rotor angles. Indeed, the helical rotor yielded  
154 a smoother machine operation with an enhanced starting ability. However, above  $45^\circ$  (from  
155  $90^\circ$  to  $180^\circ$ ), a diminishment in its magnitude was noticed and so maximum value was  
156 discovered with  $45^\circ$ . Nevertheless, a closer look into the helical Savonius rotor corroborates  
157 that there is a lack of analysis related to its design parameters optimization. Owing to its  
158 geometry complexity, it has not been utterly examined even though its substantial advantages  
159 over the traditional Savonius rotor. In this context, investigations focused mainly on the twist  
160 angle optimization<sup>27-29</sup>. As well, few parameters were examined, namely, the overlap ratio and  
161 the end plates size ratio<sup>28</sup>. The findings showed that  $45^\circ$  twisted Savonius rotor with zero  
162 overlap ratio and a 1.1 endplates ratio, as compared to other designs, achieved the maximum  
163 output performances. In addition, the effect of multi-stage was also examined<sup>30</sup>. Accordingly,  
164 it enhanced the performance of the rotor in terms of power coefficient and self-starting ability  
165 and lowered the noise and the mechanical vibrations emitted along its operation. From this  
166 perspective, the originality of the current research study lies in improving the performance of  
167 a helical Savonius wind turbine through the optimizing of its blade shape as this technique  
168 showed an important improvement but it remained restricted to the conventional Savonius  
169 rotor.

170 The basic target was to put in test various blade shapes of helical Savonius rotor. Indeed, in  
171 addition to the helical Savonius rotor, two novel rotor blade shapes, namely the delta bladed  
172 rotor and the two-stage delta bladed rotor were investigated numerically and experimentally.  
173 The three rotors were created using the 3D printing technology. Next, wind tunnel tests were  
174 conducted in order to get the static and dynamic experimental characteristics. Computational  
175 simulations using the ANSYS 17.0 software were performed. At this stage of analysis, we  
176 would assert that the central objective of the work at hand is to explore the blade shape

177 behavior on the characteristic performance of the helical Savonius rotor both numerically and  
178 experimentally.

179

## 180 **2. Experimental set up and methodology**

181 This section focused on the experimental approach. In fact, all needed materials were  
182 mentioned. The experimental methodology was described in details and all required  
183 measuring instruments were presented.

184

### 185 **2.1. Design and fabrication of the rotors**

186 The additive manufacturing method rests on fabricating object additively, layer by layer  
187 starting with a three dimensional numerical model created through the use of a solid modeling  
188 computer-aided design (CAD). Hence, it carries directly the digital presentation into a real  
189 physical model. Basically, the 3D printing can reduce costs while obtaining a longer lifetime  
190 compared to other manufacturing procedures. Thus, it makes the additive manufacturing  
191 method extremely beneficial and encourages its incorporation in the industrial area.

192 Through conducting a numerical and experimental investigation addressing the effect of blade  
193 shape, three designs of helical Savonius rotor exhibiting different blade shapes were  
194 undertaken. Referring to the geometrical complexity of rotor shapes, the 3D printing proved  
195 to be the most suitable as a fabrication technology. The first step consisted in designing the  
196 Savonius rotors utilizing the Solidworks software, as depicted in figure 1.a. Afterwards, the  
197 digital models were brought to a 3D printer through the use of a PLA (poly lactic acid)  
198 filament having a diameter of 1.75 mm. Based on the geometry section, the rotors have been  
199 extruded layer by layer. When all layers were handled, the rotors were manufactured and were  
200 removed from the tray for required tests, as illustrated in figure 1.b. The three proposed rotors  
201 portrayed in figure 1 are named, respectively, the helical Savonius rotor, the delta bladed rotor

202 and the two-stage delta bladed rotor. They, all, have no gap distance and a common blade  
203 portion “P” equal to P=48.5 mm. The overall geometrical parameters displayed in Table 1 and  
204 figures 1.a and 1.c are retained for the three rotors.

205

## 206 **2.2. Wind tunnel test facility**

207 The experimental tests were undertaken in a wind tunnel, which was designed and fabricated  
208 to establish the aerodynamic characteristics of the investigated helical rotors<sup>31</sup>. The wind  
209 tunnel exhibited in figure 2 is fabricated from galvanized iron sheets and Plexiglas. It mainly  
210 includes: (1) a settling chamber, (2) a diffuser, (3) a test section consisting of Plexiglas in  
211 which the turbines are placed, (4) a collector and (5) a drive section. The shape of the test  
212 section is a parallelepiped having a width of 0.4 m as well as an equal height and a length of  
213 0.8 m, where wind velocity varies from 0 to 9 m.s<sup>-1</sup>. The air speed was controlled by varying  
214 the frequency of the ventilation fan of the vacuum cleaner mounted in the drive section. The  
215 settling chamber was used to make the airflow straight in the test section. Under these  
216 conditions, it was entirely turbulent and became uniform with turbulence intensity under 1%.  
217 To ascertain this fact, various points in different location along the test section, have been  
218 chosen to measure the air velocity in each point with the use of a hot wire anemometer AM-  
219 4204. It was found that the air velocity value was nearly constant in the test vein. In addition,  
220 the more the points approach the test vein walls, the more the value of the velocity decreased  
221 until reaching 0. The specifications of the used anemometer are shown in Table 2.

222

## 223 **2.3. Experimental measurements and instrumentations**

224 The present experimental study handles three shapes of helical Savonius rotor. It relies on  
225 putting the proposed rotors amidst the test section. Two ball bearings were fixed on either side  
226 of the shaft to ensure its free rotation guidance and to support the rotors. At a selected velocity

227 accordingly at a selected Reynolds number, the incoming wind force generates the blade  
228 rotation and then a rotating power emerges as a result. To identify the proposed helical  
229 Savonius rotor shapes performance and to tackle the interaction between the flow and the  
230 blades, it is significant to express the performance using well-determinate non-dimensional  
231 parameters. Static and dynamic measurements were taken for a wind velocity equal to  $U_{\infty}=9$   
232  $\text{m.s}^{-1}$  accordingly for Reynolds number equal to  $Re= 9.8\times 10^4$ .

233 For the static measurements, relative to the wind direction, with the use of a static digital TQ-  
234 8800 torque meter,  $T_s$  the static torque was assessed. It is measured for different angular  
235 positions lying from  $0^\circ$  to  $360^\circ$  in a step of  $10^\circ$ , using a gradual disk, as highlighted in figure  
236 3. In fact, for each angular position, the rotation of the helical rotor was jammed and the  
237 corresponding static torque value was shown through the digital display. The specifications of  
238 the used static torque meter are shown in Table 3.

239 Then, using equation 1, the static torque coefficient was specified. This non-dimension  
240 parameter reflects the self-starting criteria of the rotor.

$$C_{T_s} = \frac{T_s}{\frac{1}{2}\rho R A U_{\infty}^2} \quad (1)$$

241 where  $T_s$  stands for the static torque (N),  $U_{\infty}$  indicates the incoming wind velocity ( $\text{m.s}^{-1}$ ),  $\rho$   
242 refers to the air density ( $\text{kg.m}^{-3}$ ),  $R$  corresponds to the rotor radius (m), and  $A$  expresses the  
243 projected area calculated as follows:

$$A = DH \quad (2)$$

244 where  $H$  and  $D$  are the height of the rotor and its diameter, respectively.

245 For the dynamic measurements, the rotor rotational speed  $\Omega$  was determined by a CA-1727  
246 digital tachometer model. The specifications of the tachometer are shown in Table 4.

247 Then, the tip speed ratio  $\lambda$  was estimated based on the following equation.

$$\lambda = \frac{R\Omega}{U_{\infty}} \quad (3)$$

248 Gradual loads were applied to the rotor shaft to realize the braking operation. Indeed, a 1mm  
 249 of diameter inelastic nylon wire was rolled up on the shaft maintained on the one end and on  
 250 the other end a load was suspended. For each load, the dynamic torque T was measured with a  
 251 torque meter sensor. Its instantaneous variation was displayed across computer acquisition  
 252 with the use of the installed EMPERIOR LITE software. Then the average value was recorded  
 253 and associated with the corresponding measured rotational speed (subsequently a tip speed  
 254 ratio). The used facilities are outlined in figure 4 and all required dynamic torque  
 255 specifications are shown in Table 5.

256 The experimental performances of the studied helical rotors are expressed in terms of non-  
 257 dimensional coefficients: the power coefficient ( $C_p$ ) and the torque coefficient ( $C_T$ ) as a  
 258 function of the tip speed ratio ( $\lambda$ ).

259 The torque coefficient  $C_T$  defined as the ratio between the effective produced and the  
 260 available torque can be computed in terms of:

$$C_T = \frac{T}{\frac{1}{2}\rho R A U_{\infty}^2} \quad (4)$$

261 The power coefficient  $C_p$  corresponds to the ratio between ( $P_r$ ) the rotor mechanical power  
 262 and ( $P_w$ ) the wind power.

$$C_p = \frac{P_r}{P_w} \quad (5)$$

263 The wind power is indicated as:

$$P_w = \frac{1}{2}\rho A U_{\infty}^3 \quad (6)$$

264 Equation (7) indicates the mechanical power ( $P_r$ ):

$$P_r = T\Omega \quad (7)$$

265 Hence,  $C_p$ ,  $C_T$  and  $\lambda$  are governed by equation (8)

$$C_p = C_T \lambda \quad (8)$$

266

### 267 **3. Numerical methodology**

268 The Computational Fluid dynamics (CFD) have been applied to settle flow problems for  
269 multiple applications. In aerodynamics, the wind flow around turbines can be well predicted,  
270 analyzed and interpreted through solving the governing equations describing the flow  
271 phenomena. For this reason, various codes and softwares have been used (SOLIDWORKS,  
272 COSMOS-FLOWWORKS, ADINA, FLOW SIMULATIONS and ANSYS). The solving  
273 procedure uses basically three primordial elements: (i) the preprocessor which consists in  
274 defining the overall domain, the generation of the mesh and the boundary conditions (ii) the  
275 solver that handles the choice of discretization method and (iii) the post processor which  
276 consists in outputting results.

277 As far as this research work is concerned, the numerical simulations for the proposed rotor  
278 shapes were carried out deploying ANSYS Fluent 17.0 software. It was applied to calculate  
279 the flow characteristics around complicated structures. In fact, it uses the Finite Volume  
280 Method (FVM) to resolve the Navier Stokes equations. This method is well-suited for the  
281 problem interaction between the rotating blades of the turbine and the inlet wind flow.

282

#### 283 **3.1. Simulation domain and setting of boundary conditions**

284 The overall domain and boundary conditions considered in this work are summarized in  
285 figure 5. For suitable analysis, the overall domain was split into two sub-domains which are  
286 separated using an interface created at the level of the adjacent faces. The first surrounding  
287 domain, representing the test section, is considered as steady. The second cylindrical domain,

288 containing the rotor, is considered as a rotating domain with a 185 mm of diameter with an  
289 angular velocity defined with regard to the tip speed ratio. Both domains were modeled in  
290 the ANSYS design modeler. To undergo the same circumstances as those with the wind  
291 tunnel, the inlet boundary condition was set as a uniform velocity in terms of  $U_{\infty}= 9 \text{ m.s}^{-1}$ .  
292 The outlet was set to be an atmospheric pressure. A no-slip moving walls condition was  
293 imposed in the rotor blades with a rotational speed that is conform to the tip speed ratio.  
294 Turbulence viscosity and intensity were set respectively equal to 10% and 5% for both outlet  
295 and inlet. Indeed, those values go in good consistency with those reported in the literature  
296 review<sup>32,33</sup> and are relative to the characteristics of the wind tunnel invested in the  
297 experiments.

298

### 299 **3.2. Grid generation**

300 In figure 6, the grid form of the rotor and the outer local rotating domain are illustrated. In  
301 fact, the 3D overall domain which was imported to the ANSYS meshing user interface, was  
302 discretized with the use of unstructured tetrahedral grids. Several body sizings, faces and  
303 edges were used to create finer mesh in the rotating domain compared to the steady domain.  
304 To boost the mesh quality and to trace the boundary layers built up close to the blades of the  
305 rotor, inflation was created. Hence, a prismatic layer was generated at the boundary layer and  
306 particularly near the rotor blades. So that the created mesh is suitable, the distance ( $y_n$ )  
307 defined as the first mesh node distance from the rotor walls should be well chosen. To do that,  
308  $y^+$  is the appropriate parameter to set based on equation 9. The determination of the exact  
309 value of  $y^+$  is based on the turbulence model walls law. As the SST k-w was the used one, a  $y^+$   
310 nearly equal to 1 was associated<sup>34</sup>. Twenty prismatic layers with 1.2 growth rate were created.

$$y^+ = \frac{\rho u_t y_n}{\mu} \quad (9)$$

311 where  $\mu$  stands for the fluid dynamic viscosity (Pa.s) and  $u_t$  presents the friction velocity (m.s<sup>-1</sup>).  
312 <sup>1</sup>).

313 A grid sensitivity test was performed to select the mesh that yields the suitable results.

314

### 315 **3.3. Turbulence modeling and solver setting**

316 The flow surrounding the helical Savonius rotors is naturally turbulent. For these criteria,  
317 transient simulations were undertaken. The time step, calculated as follows, should be wisely  
318 and precisely chosen.

$$\Delta t = \frac{\Delta \theta}{\Omega} \quad (10)$$

319 where  $\theta$  indicates the rotor shaft angular position (rad). It is, in fact, the degree whereby the  
320 rotating domain, subsequently the rotor, rotates per time step.  $\Omega$  denotes the rotor rotational  
321 speed (rad s<sup>-1</sup>). .

322 A time step independency test was performed by putting in test four time steps varying from  
323 0.5°/step to 3°/step. The opted one in this study corresponds to 1° rotation of the helical rotor.  
324 50 numbers of iterations were taken per time step so that the wind flow would be readjusted  
325 and renewed until being stabilized. Calculations were undertaken for six complete rotor  
326 revolutions. The consistent state of the flow was obtained after four revolutions. The details  
327 were, therefore, saved from the fourth cycle. To reach the convergence, the values of the  
328 obtained results for each iteration need to be closer to the previous one. In fact, relaxation  
329 factors are fundamental to amortize the solution from the previous one so that the stability  
330 would be reached. Thus, the momentum, turbulence and continuity residuals were settled  
331 equal to 10<sup>-5</sup> so as to monitor the convergence.

332 The resolution of the fluid dynamics problems relies on the Reynolds-Averaged Navier-  
333 Stokes equations<sup>35,36</sup>. Finite Volume Method was adopted to discretize them. The convective

334 terms as well as the temporal ones were discretized using the Second order Up-wind scheme.  
 335 The SIMPLE scheme (Semi-Implicit Methods for Pressure-Linked Equation) was applied to  
 336 infer the relation between the pressure and the velocity in the overall domain. The Navier-  
 337 Stokes governing equations according to a Newtonian fluid are prescribed under the  
 338 continuity equation indicated in (11) and the momentum equation indicated in (12).

$$\frac{\partial \rho}{\partial t} + \frac{\partial(\rho u_i)}{\partial x_i} = 0 \quad (11)$$

339 where t corresponds to the time (s) and  $u_i$  expresses the velocity component defined in the  $x_i$   
 340 coordinate direction  $x_i=(x,y,z)$ .

$$\frac{\partial(\rho u_i)}{\partial t} + \frac{\partial(\rho u_i u_j)}{\partial x_j} = -\frac{\partial p}{\partial x_i} + \frac{\partial}{\partial x_j} \left[ \mu \left( \frac{\partial u_j}{\partial x_j} + \frac{\partial u_j}{\partial x_i} - \frac{2}{3} \delta_{ij} \frac{\partial u_i}{\partial x_i} \right) \right] + \frac{\partial(-\overline{\rho u_i' u_j'})}{\partial x_j} + F_i \quad (12)$$

341 where  $F_i$  indicates the external forces applied (N) and p denotes the pressure (Pa).  
 342 The components of the Reynolds stress tensor are determined mathematically through  
 343 equation (12).

$$-\overline{\rho u_i' u_j'} = \mu_t \left( \frac{\partial u_i}{\partial x_j} + \frac{\partial u_j}{\partial x_i} \right) - \frac{2}{3} \rho k \delta_{ij} \quad (13)$$

344 where  $x_i, x_j$  corresponds to the Cartesian coordinate,  $\mu_t$  refers to the turbulent viscosity (Pa.s),  
 345  $\delta_{ij}$  stands for the Chronecker indices and k expresses the turbulent kinetic energy (Pa.s).  
 346 Four turbulence models, namely, the Standard k- $\epsilon$ , the RNG k- $\epsilon$ , the Realizable k- $\epsilon$  and the  
 347 SST k- $\omega$  were put to test under this study to clarify their effects on the numerical results. The  
 348 Shear Stress Transport (SST) k- $\omega$  model was adopted for the current simulations. Many  
 349 researchers demonstrated previously its good capability to modelize the wind turbines as well  
 350 as its good ability to predict the flow patterns<sup>37,39</sup>. In fact, the SST k- $\omega$  model combines the k-  
 351  $\omega$  model accuracy and robustness in the near wall regions and the k- $\epsilon$  model free flow in the

352 far wall regions. In this line, Nasef et al.<sup>40</sup> tested the four turbulence models and proved that  
 353 the SST k- $\omega$  model yielded optimal results.

354 The turbulence kinetic energy (k) is expressed in terms of:

$$\frac{\partial(\rho k)}{\partial t} + \frac{\partial(\rho u_i k)}{\partial x_i} = \frac{\partial}{\partial x_i} \left[ \left( \mu + \frac{\mu_t}{\sigma_k} \right) \text{grad}(k) \right] + 2\mu_t \frac{\partial u_i}{\partial x_j} \cdot \frac{\partial u_i}{\partial x_j} - \frac{2}{3} \rho k \frac{\partial u_i}{\partial x_j} \delta_{ijk} - \beta^* \rho k \omega \quad (14)$$

355 The specific dissipation rate ( $\omega$ ) is indicated by:

$$\frac{\partial \rho \omega}{\partial t} + \frac{\partial(\rho u_i \omega)}{\partial x_i} = \frac{\partial}{\partial x_i} \left[ \left( \mu + \frac{\mu_t}{\sigma_{\omega,1}} \right) \text{grad}(\omega) \right] + \gamma_2 \left( 2\rho \frac{\partial u_i}{\partial x_j} \cdot \frac{\partial u_i}{\partial x_j} - \frac{2}{3} \rho \omega \frac{\partial u_i}{\partial x_j} \delta_{ij} \right) - \beta_2 \rho \omega^2 + 2 \frac{\rho}{\sigma_{\omega,2} \omega} \frac{\partial k}{\partial x_k} \frac{\partial \omega}{\partial x_k} \quad (15)$$

356 where  $\beta^*$ ,  $\beta_2$ ,  $\sigma_k$ ,  $\sigma_{\omega,1}$ ,  $\sigma_{\omega,2}$  and  $\gamma_2$  indicate the constants of the SST k- $\omega$  turbulence model and

357  $\mu$  expresses the dynamic viscosity (Pa.s).

358

## 359 4. Results and discussion

### 360 4.1. Experimental results

361 The rotors with the proposed shapes were put to test to control the static torque coefficient  
 362 ( $C_{Ts}$ ), the power coefficient ( $C_p$ ) and the torque coefficient ( $C_T$ ). The static torque is often  
 363 used to describe the starting feature of the Savonius wind turbine. Figure 7 illustrates the  $C_{Ts}$   
 364 versus the azimuths for the tested blade shapes. It is quite obvious that the variation of the  
 365 static torque coefficient is periodic with a 180° cycle. The static torque coefficient increases  
 366 progressively until achieving its peak. The highest  $C_{Ts}$  holds the value of 0.161, 0.178 and  
 367 0.184 for an azimuth of 100°, 70° and 60° for the helical, the delta bladed and the two-stage  
 368 delta bladed rotors, respectively. Concerning the remaining cycle, the static torque coefficient  
 369 drops to reach its minimum at 180°. Regarding these observations, the maximum value of the  
 370 static torque coefficient is determined from a maximum quantity of the air that has attacked

371 the concave side of the blade for the corresponding angular position. Yet, its minimum value  
372 is specified from the air force that has not suitably moved to the concave side of the blade at  
373 the matching azimuth. Hence, the graph indicates that the two-stage delta bladed rotor  
374 displays an improvement with regard to the starting performance over the delta bladed and the  
375 helical Savonius rotors.

376 Referring to the tip speed ratio in a range varying from 0.4 to 1.2, ( $C_T$ ), the torque coefficient  
377 and ( $C_p$ ) the power coefficient were computed before being plotted in figures 8.a and 8.b,  
378 respectively. Departing from figure 8.a, it is inferred that the torque coefficient is inversely  
379 proportional to the tip speed ratio. This fact is related to the application of gradual loads  
380 which not only minimizes the rotational speed of the turbine but also increases the value of  
381 the rotor torque. The highest value of the torque coefficient is obtained at a tip speed ratio  
382 around  $\lambda=0.44$ . It seems to be equal to 0.214, 0.224 and 0.251 for the helical Savonius rotor,  
383 the delta bladed rotor and the two-stage delta bladed rotor, respectively. Relying upon figure  
384 8.b, the power coefficient increments with the augmentation in the tip speed ratio until  
385 reaching its peak over which it decays even with the augmentation in the tip speed ratio. The  
386 peaks of the power coefficient value proved to be 0.122, 0.142 and 0.158 at a tip speed ratio  
387 of 0.7, 0.78 and 0.86 for the helical Savonius rotor, the delta bladed rotor and the two-stage  
388 delta bladed rotor. Therefore, the two-stage delta bladed rotor revealed experimentally an  
389 enhancement of  $C_p$ ,  $C_T$  and even  $C_{Ts}$ . Interestingly, an improvement in  $C_p$  by 29.5 % for the  
390 two-stage delta bladed rotor and 16.4 % for the delta bladed rotor over the helical Savonius  
391 rotor was recorded departing from experimental tests. It is noteworthy that both proposed  
392 novel shapes displayed a better interaction between the projected area of the rotor blades and  
393 the incoming flow. Hence, a higher  $C_p$  was obtained over the helical Savonius rotor.

394 Table 6 exhibits the experimental uncertainties on  $\lambda$ ,  $C_{Ts}$ ,  $C_p$  and  $C_T$  for the suggested rotors at  
395 a Reynolds number equal to  $Re= 9.8\times 10^4$ . Table 7 displays the experimental error for the  
396 different used measurement instruments in this study.

397

#### 398 **4.2. Numerical model validation**

399 A grid study was performed for the generated mesh to offer well estimated numerical results  
400 with respect to the corresponding experiments. Indeed, a mesh refinement was undertaken  
401 wherein the size of the generated mesh progressively increased until the computational results  
402 didn't undergo a significant alteration. The simulations were carried out for the helical  
403 Savonius rotor with four grid sizes, named coarse, medium, fine and extra fine having  
404 260000, 570000, 880000 and 1.3 million nodes, respectively. At a tip speed ratio  $\lambda=0.73$  and  
405 for the proposed mesh sizes, the torque coefficient versus the azimuth is plotted in figure 9. A  
406 good agreement was detected between the extra fine and the fine mesh. However, the error  
407 was more important for the coarse and the medium grid. At this stage, the fine grid could be  
408 chosen for the following simulations considering numerical precision and time saving.

409 Furthermore, the turbulence model effect was investigated. Four turbulence models, namely,  
410 the SST  $k-\omega$ , the RNG  $k-\epsilon$  the standard  $k-\epsilon$  and the realizable  $k-\epsilon$  were examined. Figure 10  
411 depicts the effect of the turbulence model on the torque coefficient of the helical Savonius  
412 rotor at  $\lambda=0.73$ . A good agreement between the SST  $k-\omega$  and the realizable  $k-\epsilon$  was recorded  
413 in comparison with experimental data. However, the RNG  $k-\epsilon$  and the standard  $k-\epsilon$  are far  
414 away from them. Resting upon these and since the SST  $k-\omega$  was regarded as an appropriate  
415 model for many researchers as clarified above, it was adopted for the following simulations.

416 Moreover, the time step independency test was performed by putting in test four time steps  
417 varying from  $0.5^\circ/\text{step}$  to  $3^\circ/\text{step}$ . Figure 11 illustrates the effect of the time step on the torque  
418 coefficient of the helical Savonius rotor at  $\lambda=0.73$  in a steady phase. It is clear that with higher

419 degree per step exactly  $3^\circ/\text{step}$  and  $2^\circ/\text{step}$ , the average value of the torque coefficient  
420 decreases. This is because the flow mechanics, close to the rotor blades, are not effectively  
421 captured. This fact increases the error between numerical and experimental findings.  
422 Considering that the two time steps  $1^\circ/\text{step}$  and  $0.5^\circ/\text{step}$  gave approximately similar torque  
423 coefficient variation, the  $1^\circ/\text{step}$  time was selected under this study to save time calculation.  
424 We enacted a comparison between experimental and numerical results in order to assess the  
425 goodness of fit and reliability of the CFD model. Figure 12 depicts both experimental and  
426 numerical power coefficients of the helical Savonius rotor. The average error is about 5% for  
427 a tip speed ratio ranging from 0.4 to 1.2. From this perspective, it is noteworthy that the  
428 simulation findings go in good conformity with the experimental ones. For this reason, we  
429 may assert that the CFD model is suitable to assess the effect of the proposed blade shapes on  
430 the efficiency of the helical Savonius wind rotor.

431

### 432 **4.3. Numerical results**

#### 433 **4.3.1 Aerodynamic flow characteristics**

##### 434 **Magnitude Velocity**

435 Figure 13 illustrates the velocity distributions for the two-stage delta bladed, the delta bladed  
436 and the helical rotors at a tip speed ratio  $\lambda=0.73$ . The velocity upstream the rotor seems to be  
437 uniform and takes the value of  $V=9 \text{ m.s}^{-1}$ , as put in the entry boundary conditions. A  
438 deceleration near the rotor blades turns out to be in the range of  $4\text{-}5 \text{ m.s}^{-1}$  for the suggested  
439 shapes. The vortices formation in the concave side of the returning blades is more  
440 considerable for the two-stage delta bladed rotor. At this stage, the flow passage is more  
441 significant which lessens the negative drag and enhances the turbine rotation anti-clockwise.  
442 The maximum velocity region is situated on the convex side of the returning blade (figure  
443 13.a). It proves to be lower in the two-stage delta bladed rotor than the others. It is equal to

444  $V=14.6 \text{ m.s}^{-1}$ ,  $V=15.3 \text{ m.s}^{-1}$  and  $V=16.7 \text{ m.s}^{-1}$ , respectively, for the two-stage delta bladed, the  
445 delta bladed and the helical Savonius rotor. Indeed, it indicates a less important negative drag  
446 and, as a matter of fact, a higher power coefficient  $C_p$  compared to the others. A wake zone of  
447 velocity, downstream the rotor, is clearly detected. In figure 13.b, the wake shape is lucid and  
448 influenced by the rotor blade shape. The wake formation refers basically to the centrifugal  
449 force that is responsible for the blades rotation. Moreover, a symmetrical aspect is depicted  
450 for all configurations, but proves to be more significant for the two-stage delta bladed rotor.  
451 With regard to the same axial position  $z=0.097 \text{ m}$  and for a  $y$  direction lying from  $y=-0.27 \text{ m}$   
452 to  $y=0.27 \text{ m}$ , various points are produced for each transversal plane identified respectively by  
453  $x=-0.3 \text{ m}$ ,  $x=-0.08 \text{ m}$ ,  $x=0.08 \text{ m}$  and  $x=0.3 \text{ m}$ . The velocity profiles for those specific points  
454 and for all the examined rotor shapes are plotted in figure 14. The velocity evolution upstream  
455 and downstream the rotor reinforces what was mentioned before. In fact, at upstream  $x=-0.3$   
456  $\text{m}$  (figure 14.a), the velocity holds the value of the inlet condition  $V=9 \text{ m.s}^{-1}$ . A sharply  
457 deceleration is detected in figure 14.b which stands for the area near the rotors. Figures 14.c  
458 and 14.d reveal a drop in the velocity downstream. This drop introduces the wake zone. It  
459 becomes more significant while moving away from the rotor and proves to be more notable  
460 for the two-stage delta bladed rotor (figure 14.d). Indeed, it is clearly inferred that the two-  
461 stage delta bladed rotor is more efficient than other rotor shapes.

462

### 463 **Total pressure**

464 The total pressure distributions for the two-stage delta bladed, the delta bladed and the helical  
465 Savonius rotor at a tip speed ratio  $\lambda=0.73$  are highlighted in figure 15 at a tip speed ratio  
466  $\lambda=0.73$ . The pressure distribution for all proposed configurations exhibits two basic  
467 characteristic features: a high and a low-pressure zone. In fact, the high pressure zone  
468 manifests upstream the rotor. With a uniform aspect, it holds almost the same value for all the

469 rotor shapes (around 65 Pa). As for the low pressure zone, it appears around the rotor and  
470 more clearly downstream the advanced blade. By zooming on the rotating area, one can notice  
471 that the distribution of the total pressure on both the advanced and the returning blade is  
472 completely different for the proposed shapes. The pressure difference is responsible for the  
473 turbine anti-clockwise rotation. In fact, the higher total pressure value on the convex side of  
474 the returning blade is obtained for the helical Savonius rotor ( $p=72.56$  Pa). It is relatively less  
475 significant for the two-stage delta bladed rotor ( $p=70.73$  Pa) and the delta bladed rotor  
476 ( $p=68.45$  Pa). Hence, the negative drag decreases with the novel blade shapes. As for the  
477 advancing blade concave side, the total pressure is more significant for the two-stage delta  
478 bladed rotor ( $p=-3.2$  Pa) than for the delta bladed rotor ( $p=-5.9$  Pa) as well as the helical  
479 Savonius rotor ( $p=-11.85$  Pa) even in negative values. At this level, the positive drag is  
480 improved with the novel two-stage delta bladed shape. This, therefore, results in a higher  
481 torque for the two-stage delta bladed Savonius rotor. Additionally, it is obvious departing  
482 from figure 15.b that there is a significant change in the total pressure behind the rotors. It is  
483 related in fact to the vortices formation which is more significant for the two-stage delta  
484 bladed rotor over the others. The blade shape promotes in a straight forward way the pressure  
485 recovery. Indeed, here the wake phenomenon is clear again. Consequently, the drag force  
486 would be more significant for the two-stage delta bladed Savonius rotor. Hence, the novel  
487 two-stage delta bladed shape performance proves to be optimal compared to the other blade  
488 shapes.

489 To gain a better and deeper insight into the pressure comparison in different rotor positions,  
490 the pressure coefficient distribution upon the concave side of the advancing blade and on the  
491 convex side of the returning blade is outlined in figures 16 and 17, respectively, during one  
492 revolution for the proposed rotor shapes at  $\lambda=0.73$ . Following the power coefficient ( $C_{pr}$ )  
493 evolution versus the angular position ( $\Theta$ ), it is worth noting that the  $C_{pr}$  drops with the rise of

494  $\Theta$  in the range of (0-180°). Notably, with the turbine rotation, the blade projected area varies  
495 compared to the incoming wind direction. As a matter of fact, the topmost value of  $C_{pr}$  is  
496 obtained for  $\Theta=0^\circ$  for all configurations and for both concave and convex side (maximum  
497 projected area). With respect to the maximum pressure coefficient on the concave side of the  
498 advanced blade (figure 16), the two-stage delta bladed rotor displays a higher value of ( $C_{pr}$   
499 =0.53) than the delta bladed ( $C_{pr}=0.36$ ) and the helical Savonius rotor ( $C_{pr}=0.21$ ). With  
500 respect to the maximum value of the power coefficient on the returning blade convex side  
501 (figure 17), the two-stage delta bladed exhibits equally a better value of ( $C_{pr}=1.94$ ) compared  
502 to the delta bladed ( $C_{pr}=1.84$ ) as well as the helical Savonius rotor ( $C_{pr}=1.74$ ). Moreover, the  
503 lower value of  $C_{pr}$  is recorded for  $\Theta=120^\circ$  (lower projected area). Furthermore, the two-stage  
504 delta bladed rotor indicates the optimum values in this angular position on both concave side  
505 of the advanced blade ( $C_{pr}=-1.09$ ) and the convex side of the returning blade ( $C_{pr}=-2.85$ ). At  
506 this stage of synthesis, it is clearly inferred that the novel two-stage delta bladed rotor stands  
507 for the optimal shape regarding the pressure coefficient for almost all angular positions.

508

### 509 **The turbulent intensity**

510 Figure 18 illustrates the turbulent intensity distribution for the introduced rotor shapes at a tip  
511 speed ratio amounting to  $\lambda=0.73$ . Upstream, the turbulent intensity seems to be weak for all  
512 the configurations. Near from the rotating domain and more precisely behind the returning  
513 blade concave side, the turbulent intensity for the two-stage delta bladed rotor is lower lying  
514 within the range of 0.051-0.115. Nevertheless, it lies in the range of 0.067-0.121 and 0.073-  
515 0.146, respectively, for the delta bladed and the helical Savonius rotor, which collaborates a  
516 better flow passage to the concave returning blade. Downstream, the turbulent intensity is  
517 more pronounced. The symmetrical aspect is more observed for the novel two-stage delta  
518 bladed rotor than for others (figure 18.b). However, for the helical Savonius rotor the

519 maximum value is close to 0.2. The value of 0.189 and 0.179 recorded for the delta bladed  
520 and the two-stage delta bladed rotor, respectively, are indicative of the wake flow zone as  
521 stated before. Consequently, the turbulent intensity is less important in the case of the two-  
522 stage delta bladed rotor. A lower value of the turbulent intensity involves less fluctuations  
523 suggesting, therefore, a better performance in terms of flow passage and, hence, a better  
524 rotation anti-clockwise.

525

#### 526 **4.3.2 Performance characteristics**

527 The variation of the torque coefficient versus the angular position  $\Theta$  for one revolution is  
528 foreground in figure 19 for the examined rotor blade shapes at  $\lambda=0.73$ . Departing from these  
529 results, it is deduced that the modification of the blade shape influences positively the  
530 variation of the torque coefficient of the helical Savonius wind rotor. Indeed, its average value  
531 moves from  $C_T=0.17$  for the helical Savonius rotor to  $C_T=0.19$  and  $C_T=0.2$  for the delta  
532 bladed and the two-stage delta bladed rotor, respectively. Accordingly, a significant  
533 enhancement at the level of the predicted produced torque is recorded which is justified by the  
534 enhancement of the positive drag applied on the advancing blade on the one hand and the  
535 reduction in terms of the negative drag applied on the returning blade on the other hand,  
536 resulting in an improved net torque thanks to the novel blade shapes.

537 The effect of the blade shape on the helical Savonius wind rotor efficiency was assessed. A  
538 superposition of the experimental as well as the numerical power coefficients for a wide  
539 variety of tip speed ratios ranging from  $\lambda=0.4$  to  $\lambda=1.2$  is highlighted in figure 20 for the two-  
540 stage delta bladed, the delta bladed and the helical rotors. Relying upon the numerical results,  
541 the maximum power coefficients prove to equal  $C_{p,max}= 0.124$ ,  $C_{p,max}= 0.148$  and  $C_{p,max}=$   
542  $0.152$  at a tip speed ratio of  $\lambda= 0.73$ ,  $\lambda= 0.8$  and  $\lambda= 0.88$  respectively for the helical Savonius  
543 rotor, the delta bladed rotor and the two-stage delta bladed rotor. Thus, numerically an

544 improvement in  $C_p$  by 22.58 % for the two-stage delta bladed rotor and 19.35 % for the delta  
545 bladed rotor over the helical Savonius rotor is reported.

546 In the same vein, a remarkable enhancement in terms of the power coefficient is obtained  
547 experimentally and validated numerically with the novel blade shapes. The two-stage delta  
548 bladed rotor displays the highest efficiency over the other rotors. A better interaction between  
549 the projected area of the rotor blades and the incoming flow with a better starting capability  
550 yielding therefore a higher  $C_p$  are reported over the helical Savonius rotor.

551

### 552 **4.3.3. Comparison with previous studies**

553 The current study results comprising of the blade shape impact on the helical Savonius wind  
554 rotor are compared to results accessible in the literature. Table 8 gives all necessary details.

555 The comparison denotes that, the proposed researches were realized under similar wind speed  
556 conditions and with neighboring rotors geometric parameters mainly an aspect ratio near to 1  
557 and 0 overlap distances. The present study highlights important and well-ranked value of  
558 power coefficients compared with others. It shows a  $C_{p \max} = 0.152$  with the optimum blade  
559 shape design which is slightly lower than some works and rather higher than other.

560

## 561 **5. Conclusions**

562 In the present study, in view of performance betterment of a helical Savonius wind turbine,  
563 the blade shape optimization was explored both experimentally and numerically. For this,  
564 two-novel blade shape rotors, named delta bladed and two-stage delta bladed in addition to a  
565 helical Savonius rotor were investigated. Experimental tests were conducted in a wind tunnel  
566 for a Reynolds number equal to  $Re = 9.8 \times 10^4$ . Numerical tests were conducted using Ansys  
567 Fluent software through 3-D unsteady simulations. The main conclusions obtained under this  
568 study are:

- 569 • Departing from the wind tunnel tests, a considerable improvement in  $C_p$  by 29.5 % for  
570 the two-stage delta bladed rotor and 16.4 % for the delta bladed rotor over the helical  
571 Savonius rotor was recorded.
- 572 • The torque coefficient was equally enhanced. At  $\lambda=0.44$ , highest value reached  
573  $C_{T,max}=0.251$  and  $C_{T,max}=0.224$  for the two-stage delta bladed and the delta bladed  
574 rotor, respectively, while it was equal to  $C_{T,max}=0.214$  for the helical Savonius rotor.
- 575 • The starting ability was evenly boosted as a significant amelioration in the static  
576 torque coefficient. The highest  $C_{Ts}$  holds the value of  $C_{Ts,max}= 0.184$  and  $C_{Ts,max}=  
577 0.178$  for the two-stage delta bladed and the delta bladed rotor, respectively, over the  
578 helical rotor  $C_{Ts,max}= 0.161$ .
- 579 • Numerically, the aerodynamic flow characteristics were improved with the novel  
580 blade shapes as a considerable betterment in the flow velocity and the pressure  
581 recovery were observed mainly around the rotor blades.
- 582 • The maximum power coefficient was higher by 22.58 % for the two-stage delta bladed  
583 rotor and 19.35 % for the delta bladed rotor over the helical Savonius rotor. It reached  
584  $C_{p,max}= 0.152$  for the two-stage delta bladed rotor,  $C_{p,max}= 0.148$  for the delta bladed  
585 rotor until it was equal to  $C_{p,max}= 0.124$  for the helical rotor at a tip speed ratio of  $\lambda=  
586 0.88$ ,  $\lambda= 0.8$  and  $\lambda= 0.73$  respectively.
- 587 • At  $\lambda=0.73$ , the torque coefficient of the helical Savonius wind rotor was positively  
588 influenced with the modification of its blade shape. Indeed, its average value moved  
589 from  $C_T=0.17$  for the helical Savonius rotor to  $C_T=0.19$  and  $C_T=0.2$  for the delta  
590 bladed and the two-stage delta bladed rotor, respectively.

591 Resting on the obtained results, the novel two stage delta bladed rotor shape, exhibited an  
592 important improvement in the aerodynamic performance over the delta bladed and the helical

593 Savonius rotor. Its use will therefore bring such a positive impact that it may revolutionize  
594 technological as well as industrial applications.

595 As a final note, we would assert that this synthesis can be regarded as promising as it may  
596 constitute a powerful impetus for researchers to explore further this area. During the research  
597 phase, it is discovered that there is still much opportunity for improvement. In future work,  
598 we plan to improve the performance of the studied helical Savonius rotor using deflector  
599 system.

600

### 601 **Acknowledgement**

602 The authors are thankful to the Laboratory of Electro Mechanic Systems (LASEM) members  
603 for the valuable financial assistance.

604

### 605 **Nomenclature**

606 A Projected area,  $m^2$

607  $A_r$  Aspect ratio, dimensionless

608  $C_{Ts}$  Coefficient of the rotor static torque, dimensionless

609  $C_{Ts,max}$  Maximum coefficient of the rotor static torque, dimensionless

610  $C_p$  Coefficient of the rotor power, dimensionless

611  $C_{p,max}$  Maximum coefficient of the rotor power, dimensionless

612  $C_T$  Coefficient of the rotor torque, dimensionless

613  $C_{T,max}$  Maximum coefficient of the rotor torque, dimensionless

614  $C_{pr}$  Coefficient of the rotor blade surface pressure, dimensionless

615 D Rotor diameter, m

616  $D_e$  End plate diameter, m

617 d Blade chord length, m

618	$e$	Overlap distance, m
619	$F_i$	External applied forces, N
620	$G_k$	Production term of turbulence, $\text{kg}\cdot\text{m}^{-1}\cdot\text{s}^{-3}$
621	$H$	Rotor height, m
622	$k$	Turbulent kinetic energy, $\text{m}^2\cdot\text{s}^{-2}$
623	$p$	Pressure, Pa
624	$P_r$	Mechanical power, W
625	$P_w$	Power of the wind, W
626	$R$	Rotor radius, m
627	$Re$	Reynolds Number, dimensionless
628	$s$	Rotor shaft diameter, m
629	$t$	Time, s
630	$T$	Dynamic torque, N.m
631	$T_s$	Static torque, N.m
632	$u_i$	Velocity component specified in the $x_i$ coordinate direction, $\text{m}\cdot\text{s}^{-1}$
633	$u'_i$	Fluctuating velocity components, $\text{m}\cdot\text{s}^{-1}$
634	$\overline{u'_i u'_j}$	Components of the Reynolds stress tensor, $\text{m}^2\cdot\text{s}^{-2}$
635	$U_\infty$	Velocity Incoming wind, $\text{ms}^{-1}$
636	$V$	Velocity of the flow around the rotor, $\text{ms}^{-1}$
637	$x$	Cartesian coordinate, m
638	$x_i$	Cartesian coordinate, m
639	$y$	Cartesian coordinate, m
640	$y^+$	Non dimensional parameter
641	$z$	Cartesian coordinate, m
642	$\beta^*$	Constant of k- $\omega$ turbulent model

643  $\beta_2$  Constant of k- $\omega$  turbulent model  
644  $\sigma_k$  Constant of k- $\omega$  turbulent model  
645  $\sigma_{\omega,1}$  Constant of k- $\omega$  turbulent model  
646  $\sigma_{\omega,2}$  Constant of k- $\omega$  turbulent model  
647  $\theta$  Rotor shaft angular position, rad  
648  $\mu$  Dynamic viscosity (Pa.s)  
649  $\mu_t$  Turbulent viscosity, Pa.s  
650  $\rho$  Air density, kg.m<sup>-3</sup>  
651  $\Omega$  Rotor rotational speed, rad.s<sup>-1</sup>  
652  $\lambda$  Tip sped ratio, dimensionless  
653  $\psi$  Twist angle ,<sup>o</sup>  
654  $\delta_{ij}$  Chronecker indices  
655  $\gamma_2$  Constant of k- $\omega$  turbulent model  
656  $\Delta t$  Time step, s

657

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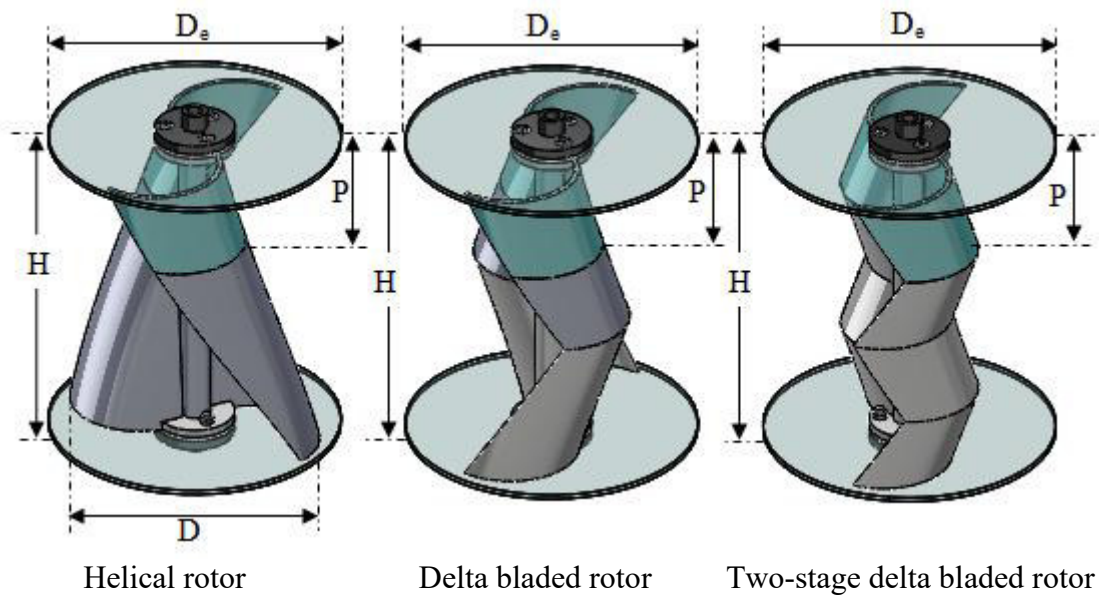
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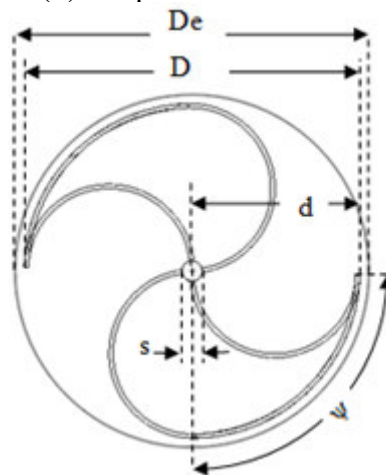
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(a) Digital models

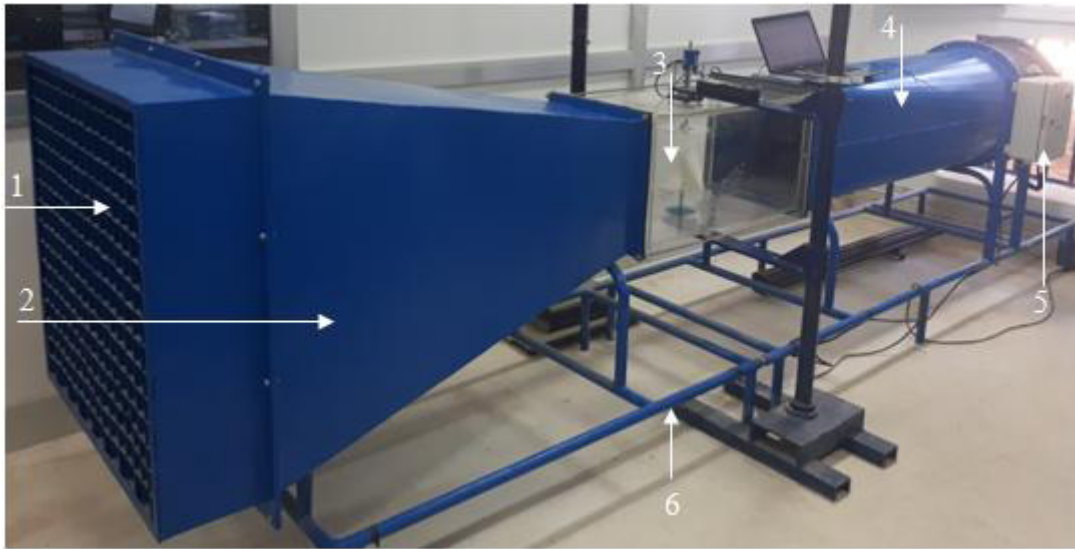


(b) 3D printed rotors



(c) Geometrical parameters

**Figure 1.** The rotor blade shapes covered in the current study.

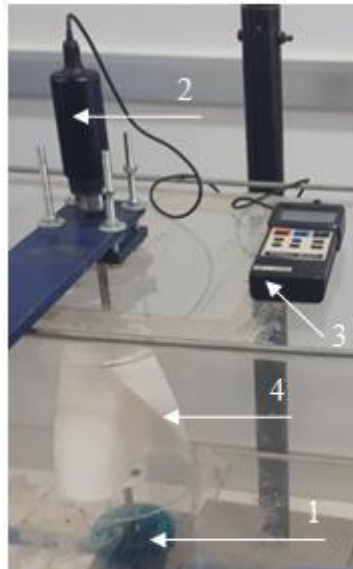


1 Settling chamber  
2 Diffuser

3 Test section  
4 Collector

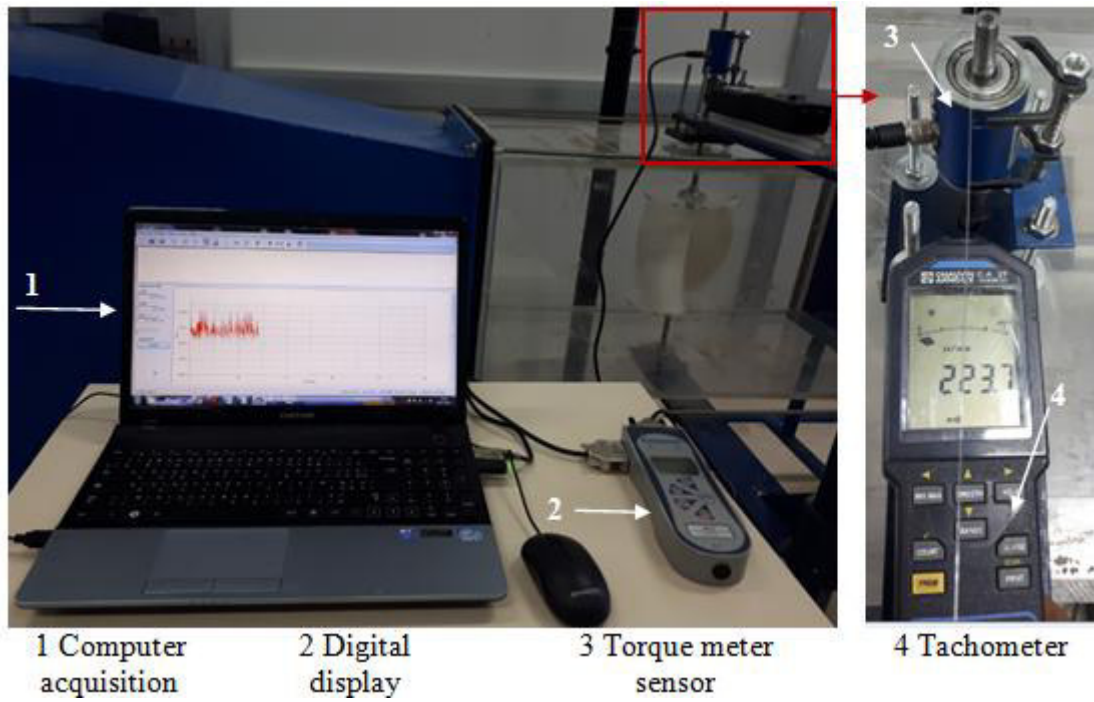
5 Drive section  
6 Structure

**Figure 2.** Wind tunnel test facility.

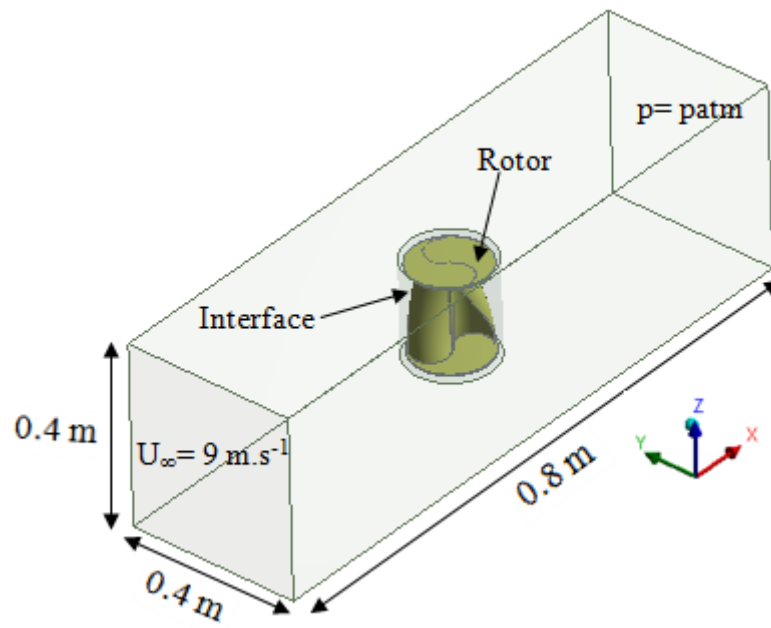


- 1 Gradual disc
- 2 Static torque meter sensor
- 3 Digital display
- 4 Rotor

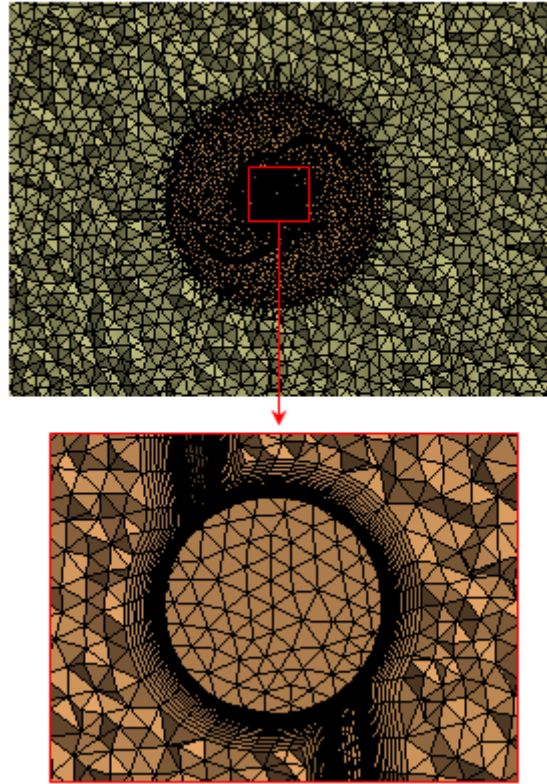
**Figure 3.** Static measurement apparatus.



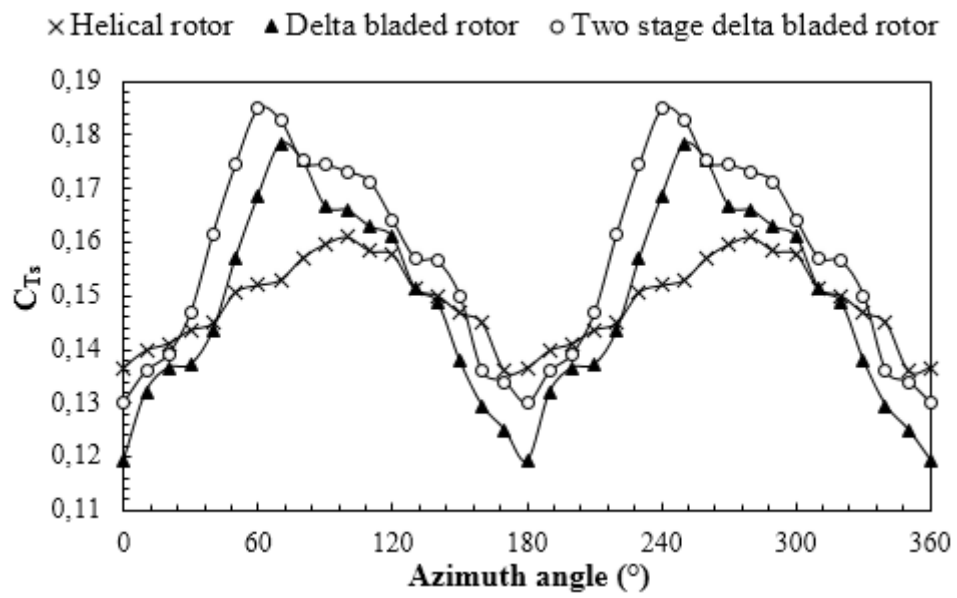
**Figure 4.** Dynamic measurement apparatus.



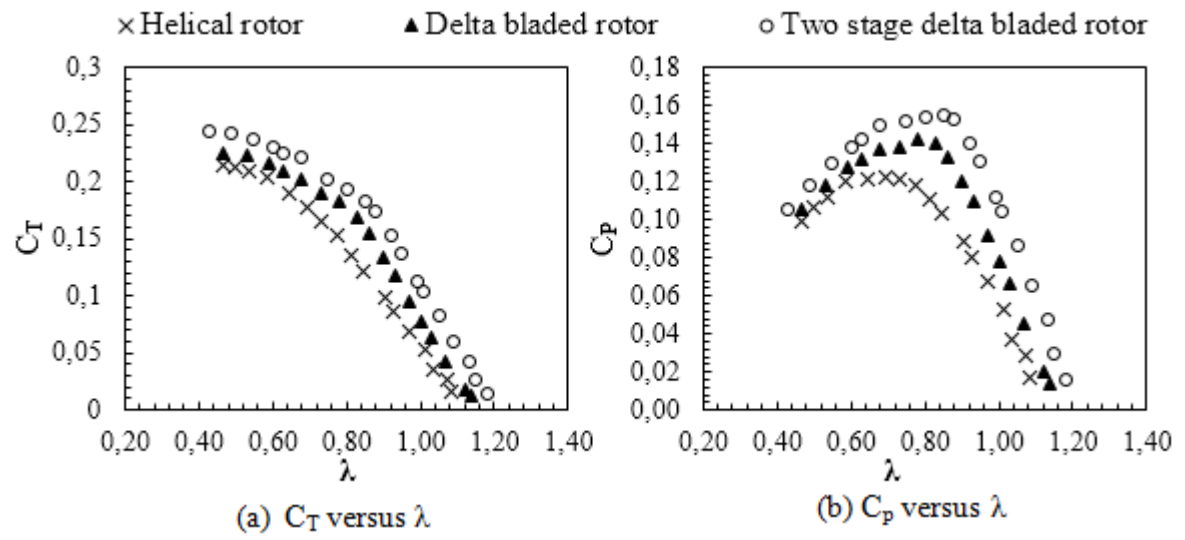
**Figure 5.** Overall domain and boundary condition details.



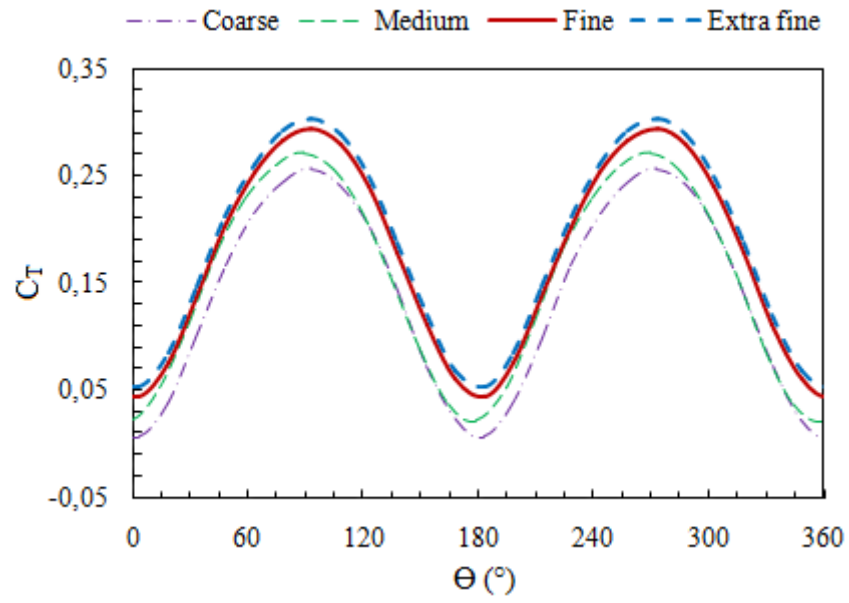
**Figure 6.** Grid generation.



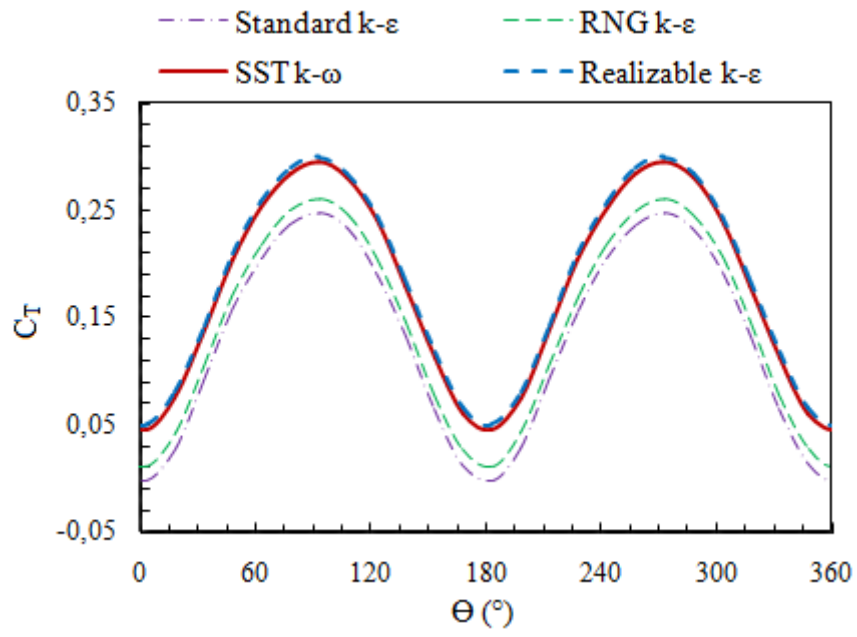
**Figure 7.** Static torque coefficient versus azimuth angle.



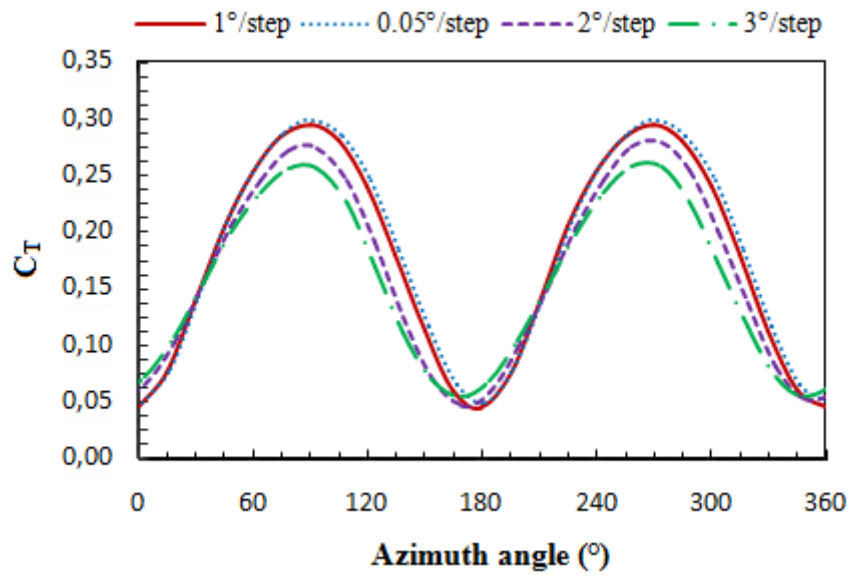
**Figure 8.** Performance characteristics.



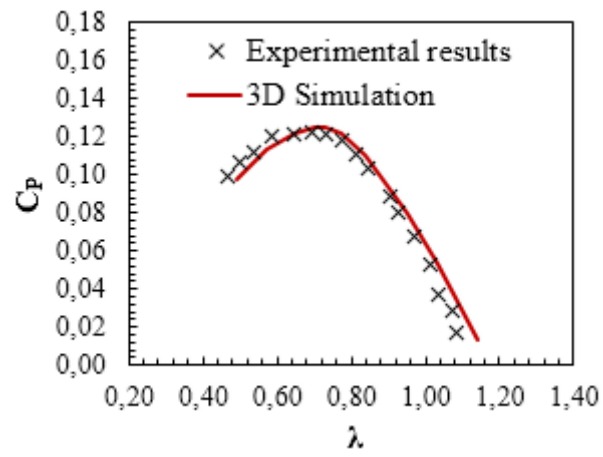
**Figure 9.** Grid form test.



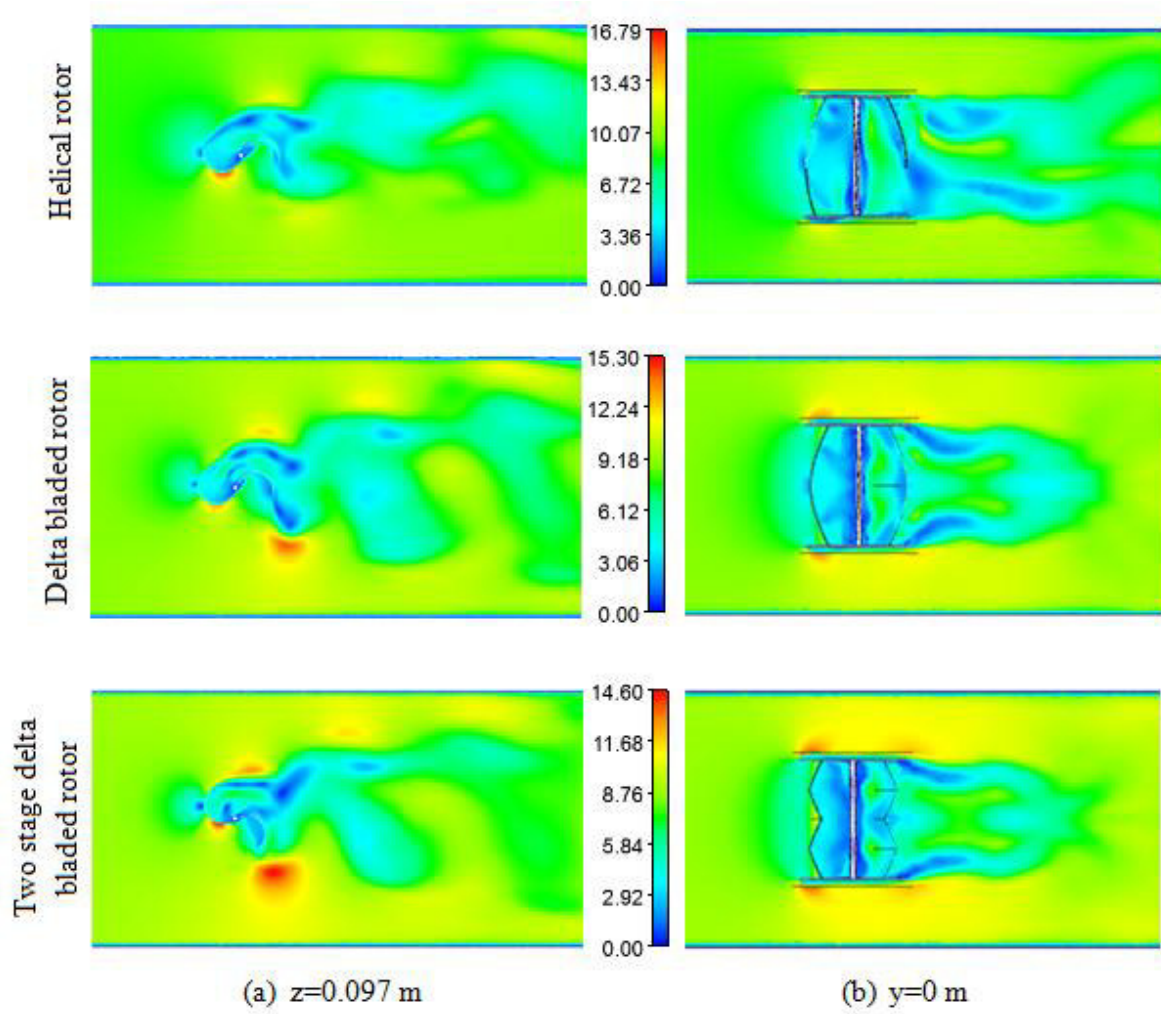
**Figure 10.** Turbulence model test.



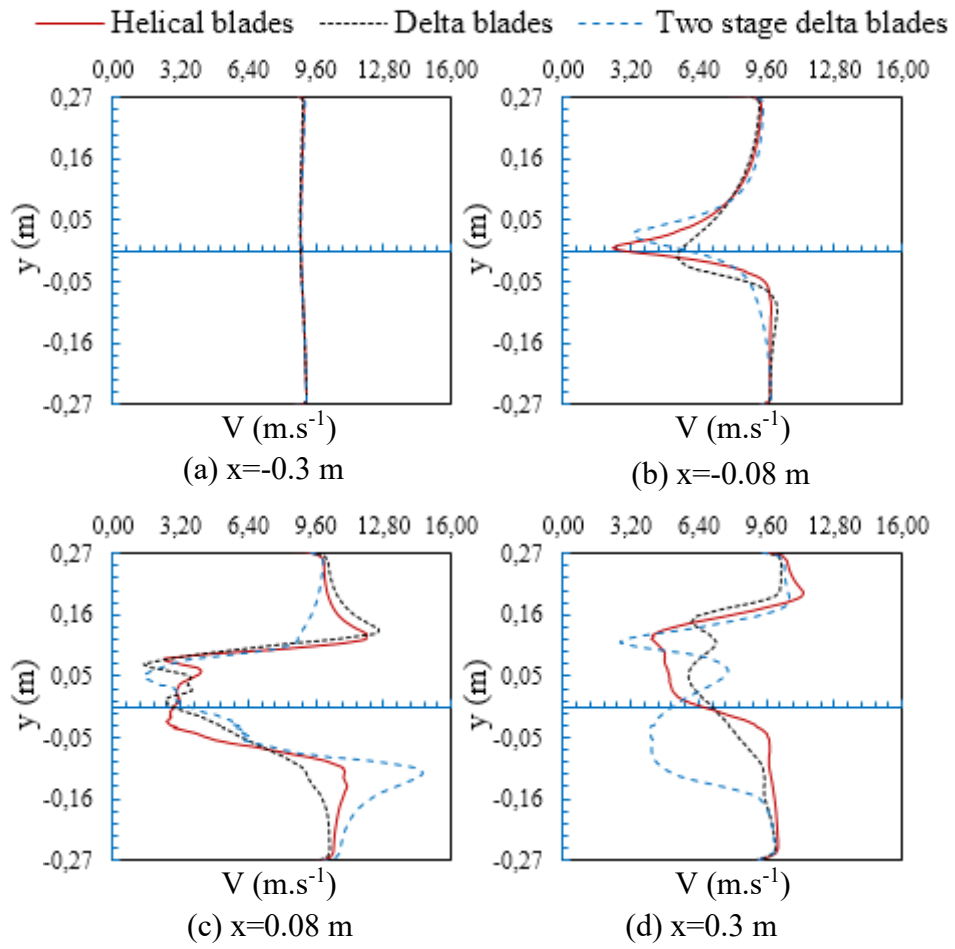
**Figure 11.** Time step independency test.



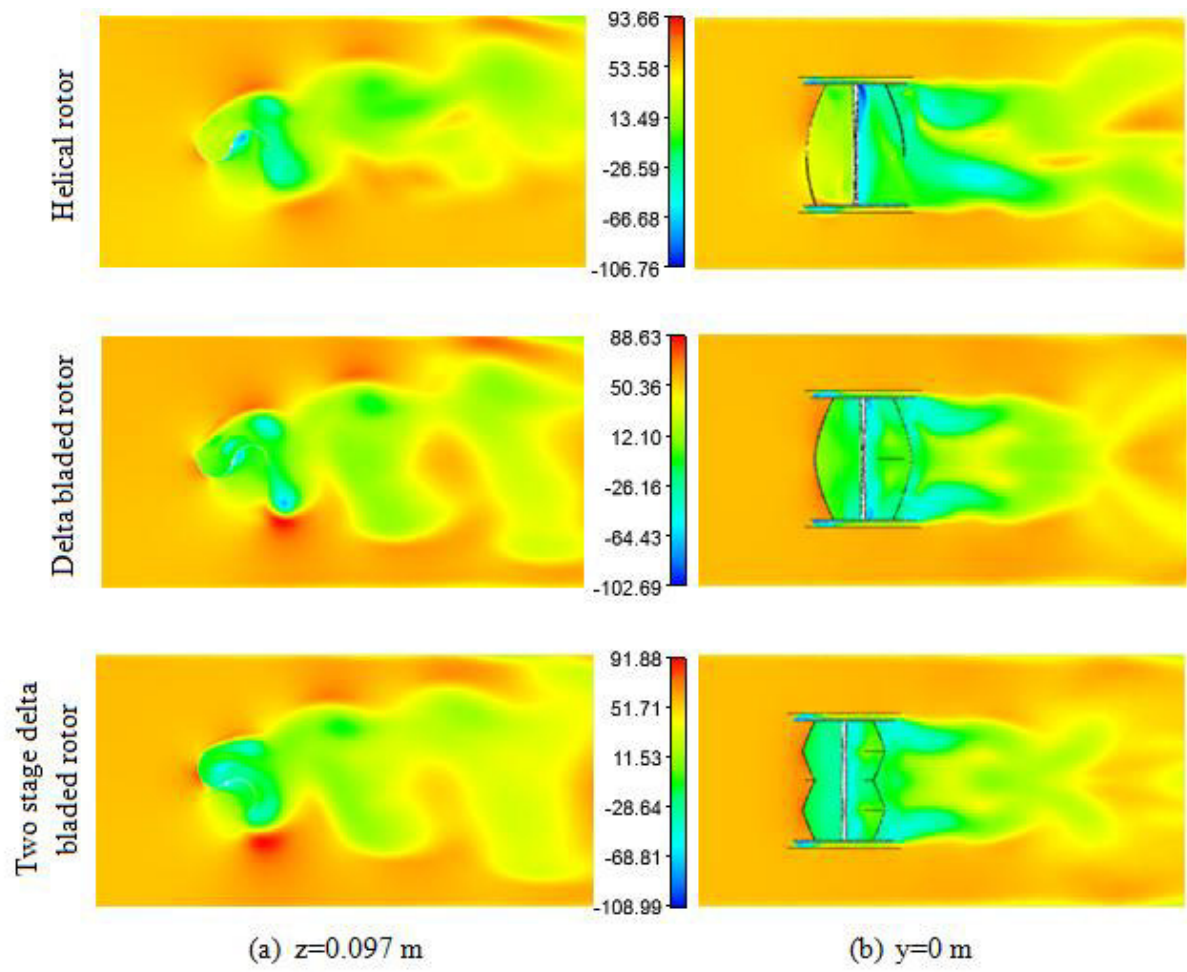
**Figure 12.** Numerical model validation.



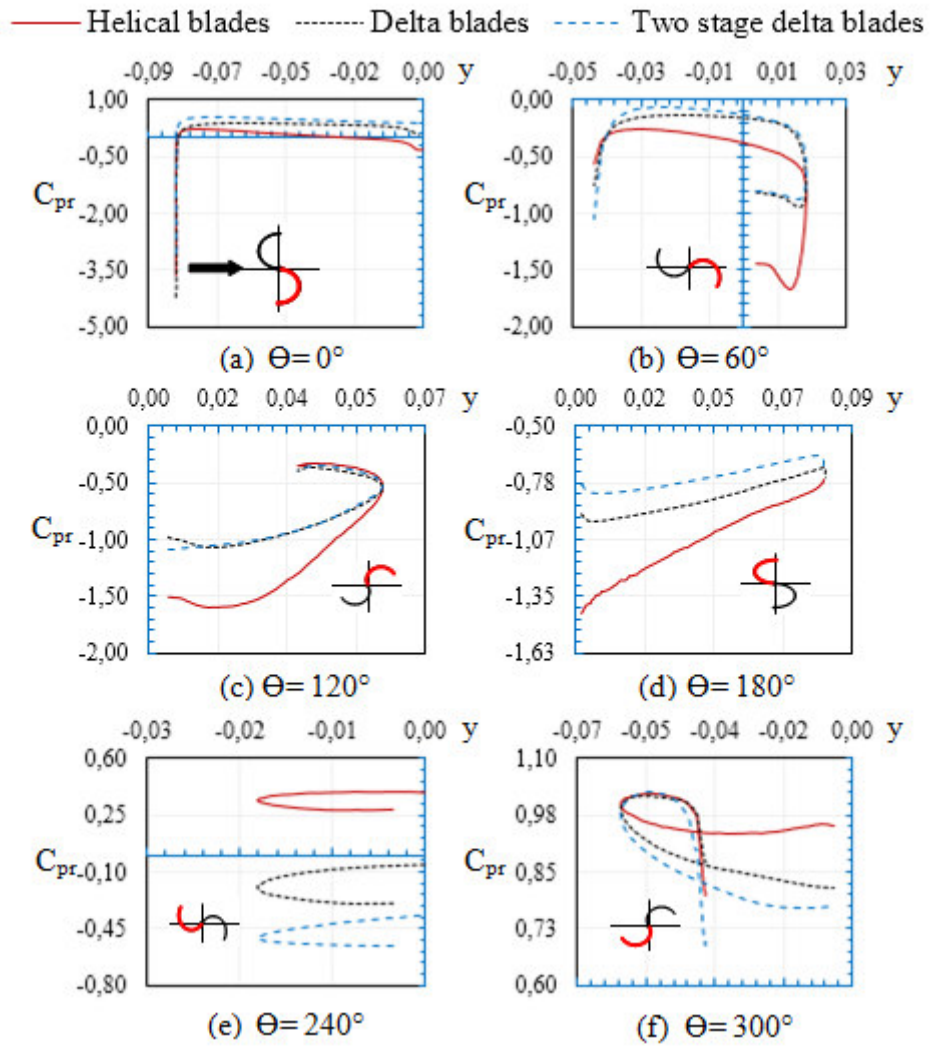
**Figure 13.** Magnitude velocity ( $\text{m.s}^{-1}$ ) distribution.



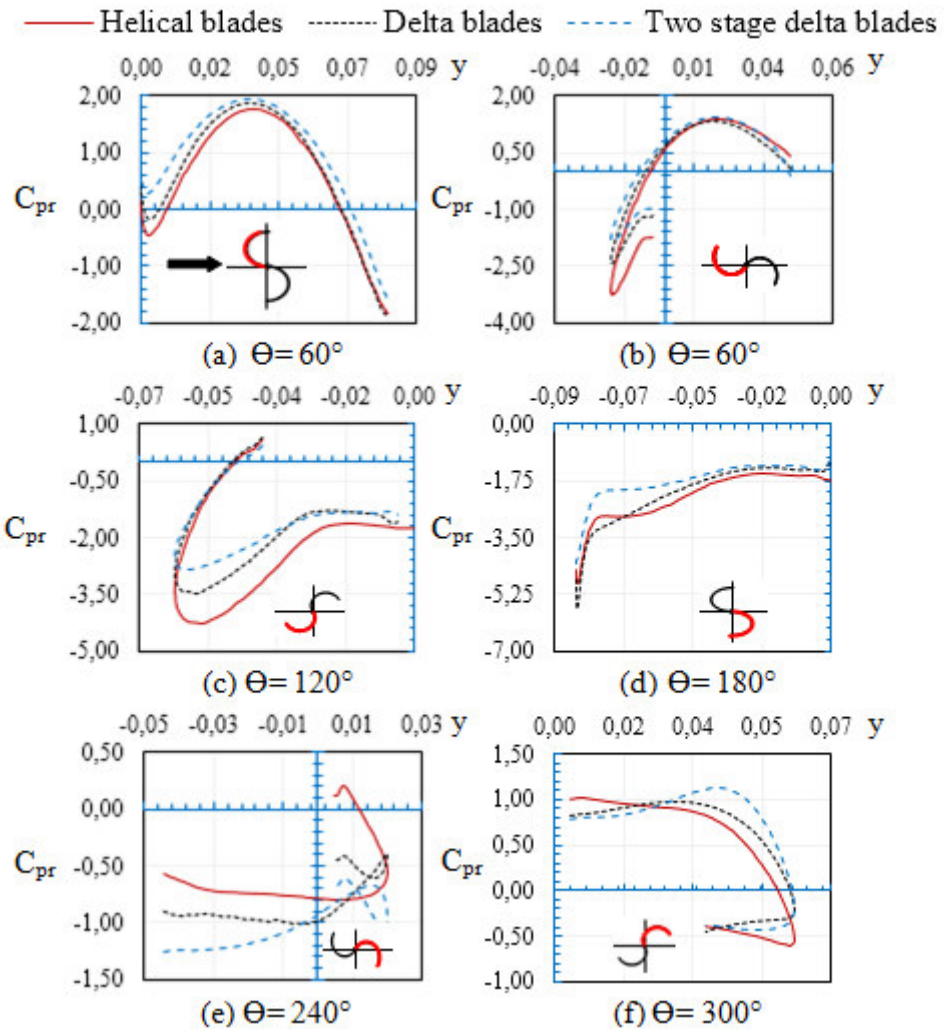
**Figure 14.** Velocity (m.s<sup>-1</sup>) profiles.



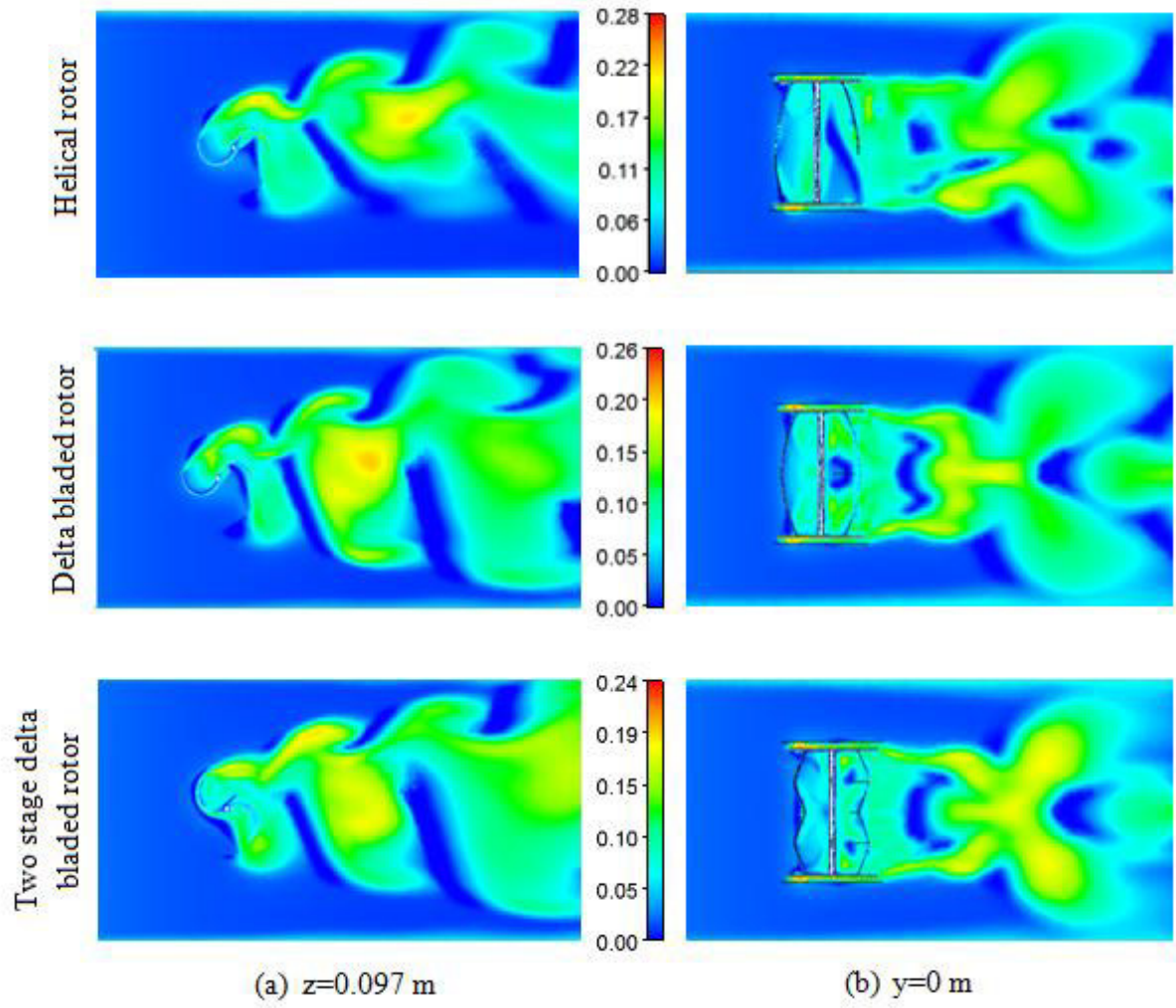
**Figure 15.** Total pressure (Pa) distribution.



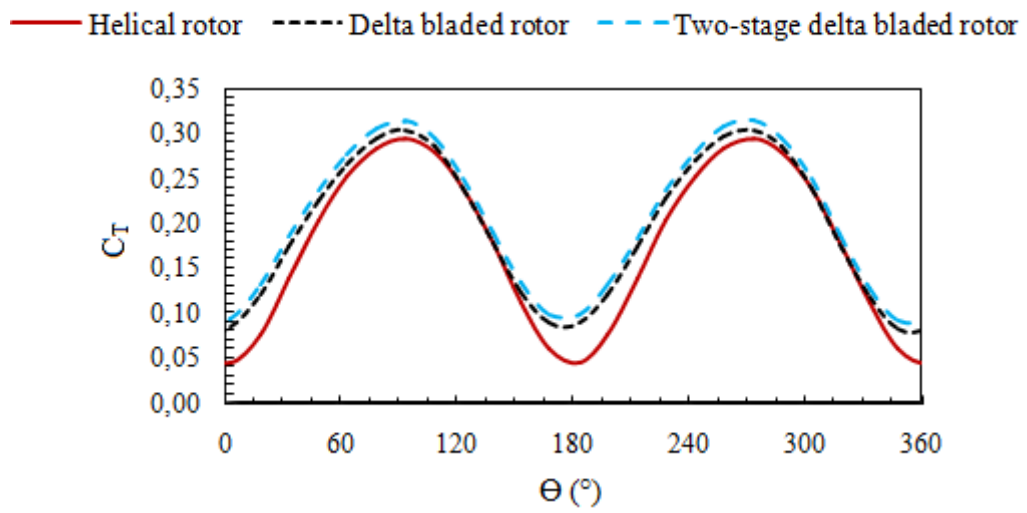
**Figure 16.** Pressure coefficient on the concave side of the advancing blade during one revolution.



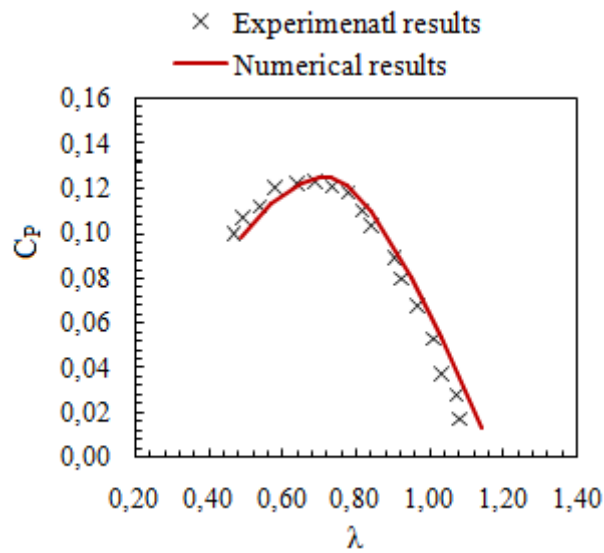
**Figure 17.** Pressure coefficient on the convex side of the returning blade during one revolution.



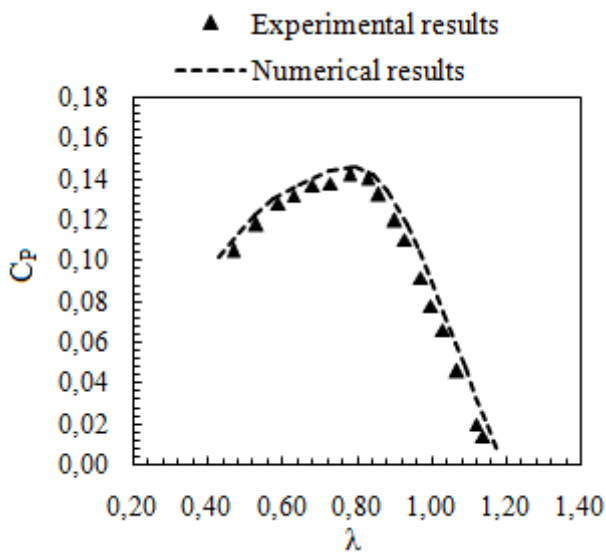
**Figure 18.** Turbulent intensity distribution.



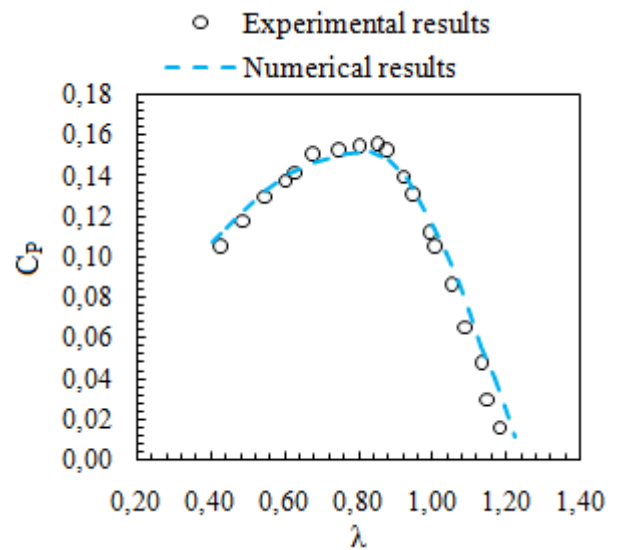
**Figure 19.** Torque coefficient at  $\lambda= 0.73$  for the different rotor blade shapes.



(a) Helical rotor



(b) Delta bladed rotor



(c) Two-stage delta bladed rotor

**Figure 20.** Comparison between experimental and numerical results for different rotor blade shapes.

**Table 1.** Geometrical parameters of the studied helical Savonius rotors.

		Geometrical parameters						
Savonius rotors	Blade profile	S (mm)	H (mm)	D (mm)	D <sub>e</sub> (mm)	e (mm)	d (mm)	Ψ (°)
Helical	Semi-circular	10	200	160	165	0	80	90
Delta bladed	Semi-circular	10	200	160	165	0	80	90
Two-stage delta bladed	Semi-circular	10	200	160	165	0	80	90

**Table 2.**Ananometer specifications.

Description	Type AM 4204
Manufacturer	Lutron
Probe type	Thermal
Measurement range	0.2 up to 20 m.s <sup>-1</sup>
Precision	±0.1 %
Resolution	0.1 m.s <sup>-1</sup>

**Table 3.** Static digital torque meter specifications.

Description	Type TQ 8800
Manufacturer	Lutron
Measurement range	0 up to 147.1 N.Cm
Precision	$\pm 1.5\%$
Resolution	0.1 N.Cm
Dimensions	180*72*32 mm

**Table 4.** Tachometer specifications.

Description	Type CA-1727
Measurement range	6 up to 10000 RPM
Precision	$10^{-4}$
Resolution	0.0006 RPM
Stability	$\pm 6$ points
Autonomy	250 measurements in 5 min
Dimension	216*72*47 mm

**Table 5.** Dynamic digital torque meter specifications.

Description	Type AFTI
Manufacturer	Mecmesin
Measurements range	0-2 up to 150 N.m
Precision	$\pm 1\%$
Dimensions	32*70*201mm

**Table 6.** Experimental uncertainties on  $\lambda$ ,  $C_{Ts}$ ,  $C_p$  and  $C_T$ 

Parameter	Uncertainty		
	Helical rotor	Delta bladed rotor	Two-stage Delta bladed rotor
$\lambda$	2.8 %	2.5 %	2.1 %
$C_{Ts}$	4.1 %	3.8 %	3.5 %
$C_T$	2.2%	2.5 %	2.3 %
$C_p$	2.9%	2.7 %	2.8 %

**Table 7.** Experimental error for the used measurement instruments.

Measurement instrument	Experimental error
Anemometer	1 %
Tachometer	3.6 %
Static torque meter	4 %
Dynamic torque meter	2.8 %

**Table 8.** Comparison with previous studies.

Reference	Methodology	Blade profile	Twist angle (°)	$C_{p \max}$	$\lambda$ at $C_{p \max}$	Wind speed (m.s <sup>-1</sup> )	$A_r$	$O_r$
Lee et al. <sup>27</sup>	Numerical and experimental	Semi-circular	45	0.13	0.54	10	1.3	0
Saad et al. <sup>28</sup>	Numerical	Semi-circular	45	0.232	0.8	10	1	0
Damak et al. <sup>41</sup>	Numerical and Experimental	Bach	90	0.2	0.75	12	0.7	0
Zadeh et al. <sup>42</sup>	Numerical and experimental	Bach developed	135	0.14	0.6	10	2	0
Kamoji et al. <sup>43</sup>	Experimental	Semi-circular	90	0.175	0.9	14	0.88	0
Anbarsooz et al. <sup>44</sup>	Numerical and experimental	Semi-circular	45	0.115	0.7	9	1	0
Current study	Numerical and experimental	Semi-circular	90	0.124	0.73	9	0.825	0
		Delta-bladed	90	0.148	0.8	9	0.825	0
		Two stage delta-bladed	90	0.152	0.88	9	0.825	0

## **Highlights**

- 1- Savonius wind turbines suffers low power efficiency.
- 2- This research paper proposes novel blade shapes of the helical Savonius rotor.
- 3- In this study, three rotors have been investigated numerically and experimentally.
- 4- The blade shape affects the performance of the helical Savonius rotor.
- 5- The maximum power coefficient was increased by 29.5% over the helical Savonius rotor.