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A model of a linear synchronous motor based on distribution theory

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The fundamental idea of this paper is to use the distribution theory to analyze linear machines in order to include in the mathematical model both ideal and non ideal features. This paper shows how distribution theory can be used to establish a mathematical model able to describe both the ordinary working condition of a Linear Synchronous Motor (LSM) as well the role of the unavoidable irregularities and non ideal features. © 2012 American Institute of Physics. [doi:10.1063/1.3679046]

I. INTRODUCTION

The fundamental idea of this paper is to use the distribution theory to analyze linear machines in order to include in the mathematical model both ideal and non ideal features.

The main difference between rotating and linear machines is that the former have periodic structures and the latter do not. However, despite the fact that linear machines do not have a periodic structure, the mathematical tools generally used to describe and design them are based on Fourier analysis; see for example Refs. 1 and 2. This approach introduces a sort of periodic boundary condition in the description of the machine that limits in some way the analysis of this kind of machine. As an example, in this kind of analysis the role of the displacement of the armature coils from their ideal position as well as the error in the position of the field coils cannot be taken into account straightforwardly. On the contrary, an approach based on stochastic analysis can allow a deeper analysis.

This paper shows how distribution theory can be used to establish a mathematical model able to describe both the ordinary working condition of a linear synchronous motor (LSM) as well the role of the unavoidable irregularities and non ideal features.

II. THE LINEAR SYNCHRONOUS MOTOR

A linear synchronous motor (LSM) consists of two parts: the induct and the inductor.

The main magnetic field is generated by the inductor. The inductor may be assembled by using either traditional windings, permanent magnets, or superconducting magnets. The first solution is used, for example, in Transrapid transportation systems, the second has been proposed for an electromagnetic aircraft launch system (EALS), and the third solution is currently tested in the Japanese maglev system; this solution uses no iron in the magnetic circuit.

The induct is made of a polyphase system (usually a three-phase system) and generates a traveling field. Thrust is

generated by the interaction between the inductor and the induct (Fig. 1).

A LSM inductor can be made of hundreds of poles distributed along hundreds of meters (e.g., transportation systems; Fig. 2) and quite often modular topologies are used.

Several motor topologies have been previously proposed and developed. One between the inductor and the induct is always longer than the other. The simplest linear motor topology is based on a single side inductor and a single side induct but most of the technologies developed use multiple inducts and inductors: in transportation systems a double sided induct is often used and in EALS a single inductor-multiple induct motor has been proposed.³

As a result, LSM are motors that tend to be quite large and probably are the largest motors technologically available. Such huge dimensions tend to present meaningful deviations from the ideal designed topologies. Such deviations can be caused by error in the induct coils placement, error in the inductor coils placement, a non uniform air gap dimension, deformation in the structure of either the moving or the stationary part of the machine, angular displacement in the multi sided machine, and so on. These kinds of errors can naturally be present also in a rotating machine but the main difference between a rotating and a linear machine is that in the former case the structure of the machine is periodic (i.e., the rotating part faces the same geometrical structure for each performed turn) and in the latter case the structure of the machine is not periodic. As a result deviations from idealities in a rotating machine can be easily included in the mathematical analysis of the machine by using Fourier analysis and introducing some ad hoc coefficients; in linear machines this approach cannot be followed straightforwardly.

III. MATHEMATICAL MODEL OF LINEAR SYNCHRONOUS MOTOR

The traditional approach to the mathematical model of a LSM describes a LSM as a multiple of a simple unit structure. This approach, as already stated, makes it difficult to include in the model the deviations from the ideal design that unavoidably are included in an actual machine. Moreover, a superconducting linear synchronous motor (SLSM)

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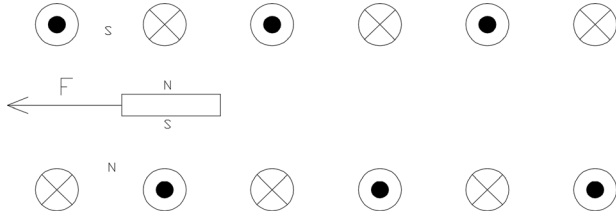


FIG. 1. Thrust generation in a LSM. Poles located in the inductor interact with the pole generated by the currents in the induct.

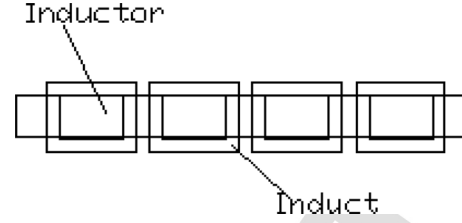


FIG. 3. General structure of a LSM.

$B_g(z, z_0, t)$ can be expressed as follows: 116

$$B_g(z, z_0, t) = \Lambda(z, z_0) \left(H_a(z, t) + H_{PM}(z, z_0, t) \right), \quad (3) \quad 118$$

where Λ is the position-dependent permanance (however, in an ironless machine the position dependence can be neglected) and where H_{PM} is the armature magneto motive force (MMF). 119 AQ6 120 121 122

The induct coils can be connected in several ways; however, the most studied are the following: the first one contains a half pole per phase (HPP), the second one contains one phase per pole (OPP). In Ref. 3 the corresponding MMFs are reported. In Fig. 4 the thrust for a constant current in the induct is shown. This figure can be interpreted as the thrust provided by the motor when the displacement of the rotor from the ideal position varies according to the coordinate z and the synchronicity between the traveling wave and the rotor is guaranteed. 123 124 125 AQ7 126 127 128 129 130 131 132

As a result, in order to evaluate the thrust in Eq. (1), one must be able to calculate H_{PM} , which can be evaluated in several ways: if it is generated by a PM it can be evaluated by using the physical parameters of the PM used, if it is generated by a bulk superconductor it can be found using the “sand-pile model” along with the Biot-Savart law, and if it is generated by a superconducting coil it can be evaluated by modeling the coil through the Biot-Savart law. 133 134 135 136 137 138 139 140

The total thrust provided by the motor along the z direction is simply obtained by multiplying the thrust provide by Eq. (1) by the number of pole pairs. The same simple approach, which describes motors with multiple inductors along the y direction, is already contained in Eq. (1) and is represented by the multiplication by Y factor inside the integral. This approach simply assumes that each pole pair is 141 142 143 144 145 146 147

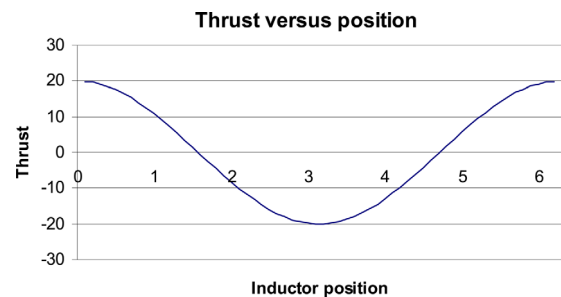


FIG. 4. (Color online) The thrust for a constant current in the induct is shown. This figure can be interpreted as the thrust provided by the motor when the displacement of the rotor from the ideal position varies according to the coordinate z and the synchronicity between the traveling wave and the rotor is guaranteed.

89 differs from a LSM because of the use of superconductors to
90 obtain the inductor field. This permits an ironless induct, but
91 if iron losses, fringing, or skin effect are neglected the mag-
92 netic analysis of the motor is identical for superconducting
93 and PM linear synchronous motor.

94 In what follows, the traditional approach is preliminarily
95 presented and subsequently the new proposed approach is
96 illustrated. As an example, the analysis is carried out for the
97 geometry of a SLSM for EMALS proposed in Refs. 3 and 5.

98 A. The 1D traditional approach to the model of LSM

99 The traditional approach to the model of a LSM consists
100 in modeling the LSM in terms of one pole and the total thrust
101 is calculated by adding the thrust provided by each pole.^{3,4} If
102 one assumes that the general structure of a LSM is the one
103 reported in Fig. 3, the thrust generated by a pole pair can be
104 expressed as:

$$F(z_0, t) = \int_0^{2\tau_p} A(z, t) B_g(z, z_0, t) Y dz, \quad (1) \quad 106$$

107 where F is the thrust provided per pole pair, z_0 is the dis-
108 placement of the inductor with respect to the induct, t is
109 time, τ_p is the polar pitch, $A(z, t)$ is the armature current sheet,
110 $B_g(z, z_0, t)$ is the air gap flux density, and Y is the motor length
111 along the y direction.

112 $A(z, t)$ can be expressed as follows:

$$A(z, t) = \frac{dH_a(z, t)}{dt}, \quad (2) \quad 113$$

115 where H_a is the armature magneto motive force (MMF).



FIG. 2. (Color online) A photograph of a Transrapid, which shows clearly how the LSM is placed along the whole vehicle. As a result the length of the motor is equal to the length of the vehicle.

from the magnetic point of view in exactly the same condition of all the other pole pairs. In a very large machine, under extreme dynamic conditions this can be not true and a mathematical way to describe this situation is useful in order to better model LSM.

B. Non idealities and the proposed approach

Deviations from the ideal design can be caused by errors in the induct coils placement, error in the inductor coils placement, a non uniform air gap dimension, deformation in the structure of either the moving or the stationary part of the machine, angular displacement in the multi sided machine, and so on. These kinds of errors can naturally be present also in a rotating machine but the main difference between a rotating and a linear machine is that in the former case the structure of the machine is periodic (i.e., the rotating part faces the same geometrical structure for each performed turn) and in the latter case the structure of the machine is not periodic. As a result, deviations from idealities in a rotating machine can be easily included in the mathematical analysis of the machine by using Fourier analysis and introducing some ad hoc coefficients; in linear machines this approach cannot be followed straightforwardly. Moreover, such deviations are intrinsically stochastic so an attempt to describe them by a deterministic analysis is not very fruitful.

The stochastic aspect of this issue can be modeled by assuming that the motor consists of N pole pairs whose displacement with respect the magnetic field generated by the induct can be described by a variable z_{0i} (where i spans from 1 to N); as a result the total thrust F_{tot} provided by the motor can be expressed as follows:

$$F_{tot} = \sum_{i=1}^N F(z_{0i}, t) = \sum_{i=1}^N \int_0^{2\tau_p} A(z, t) B_g(z, z_{0i}, t) Y dz. \quad (4)$$

If one assumes that z_{0i} is a stochastic variable then the thrust also becomes a stochastic expression, which can be described as $F_i = F(z_{0i}, t)$. The total thrust becomes, therefore,

$$F_{tot} = \sum_{i=1}^N F_i. \quad (5)$$

If one can guess the stochastic distribution of the thrust and assume that N is very large, Eq. (5) can be re-cast as follows:

$$F_{tot} = N \int_{-\infty}^{+\infty} F^*(z) P(z) dz, \quad (6)$$

where N is the normalization constant and is equal to the number of poles, F^* is the expression of the generated thrust for a z displacement, and P is the probability distribution of z .

Equation (5) and Eq. (6) show clearly that the traditional approach (i.e., the calculation of the total thrust as the multiplication of the total number of pole pairs by the thrust provided by a single pole pair) can be followed only if P is a delta function, that is, if no deviations from the ideal design are presented in the actual motor; in any other case the actual thrust is less than the ideal one.

As a matter of fact, Eq. (6) provides more information: It permits also the calculation of the reduction of the performances if one assumes to know P .

IV. CALCULATION OF THE LOSS OF PERFORMANCES OF A LSM AFFECTED BY NON IDEALITIES

As already stated, looking at Fig. 4 one can see that the thrust provided by each pole varies as a cosinusoid. As a result, if one assumes a suitable probability function the total provided thrust can be calculated using Eq. (6). In what follows, the deviations from idealities are assumed to be described by a Gaussian law, whose average is zero and the standard deviation is equal to σ . As a result, Eq. (6) can be explicitly calculated and reads as follows:

$$F_{tot} = N \int_{-\infty}^{+\infty} F^*(z) P(z) dz = \frac{N}{\sqrt{\pi\sigma^2}} \int_{-\infty}^{+\infty} T_{\max} \cos(bz) e^{-\frac{z^2}{2\sigma^2}} dz = N e^{-\left(\frac{b^2\sigma^2}{4}\right)} \quad (7)$$

where b is the inverse of the polar pitch. Equation (7) expresses the reduction of the provided thrust when deviations from idealities are present. It can be seen how the provided thrust is affected by non idealities and the reduction follows an exponential law. Equation (7) holds along the z direction but an analogous result can be found also along the Y direction

Equation (7) can be used to establish an upper limit on the error that can be tolerated in order to achieve some desired performances and is useful in order to give an estimation of the actual thrust provided by a LSM.

V. VALIDATION

To verify the results expressed by Eq. (7), a numerical model of a LSM has been constructed. In this model the magnets have been supposed to be placed in ideal positions but the armature winding positions have been supposed to be affected by errors that can be described through a Gaussian distribution. The motor has been supposed to consist of 20

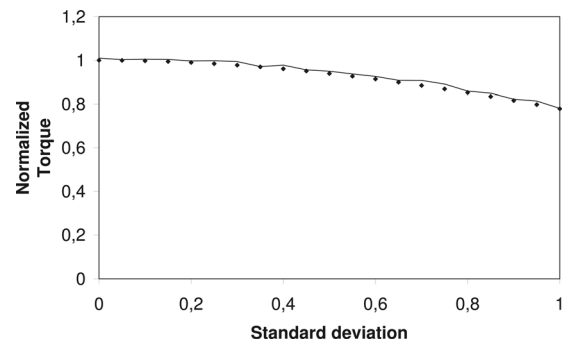


FIG. 5. A comparison between the normalized thrust obtained by the numerical model of the LSM and the thrust obtained using Eq. (7). The continuous line represents the normalized thrust calculated by the numerical model, and the diamonds represent the thrust calculated according to Eq. (7). Normalized thrust is expressed in adimensional units, σ in the unit of polar pitch.

excitation magnets whose position was distributed according to a Gaussian law whose standard deviation can be varied numerically. The total thrust has been calculated by summing the thrust obtained from each magnet. Several runs with different standard deviation of this mathematical model have been performed. Figure 5 shows a comparison between the normalized thrust obtained by the numerical model of the LSM and the thrust obtained by using Eq. (7). The distribution of the results obtained with this mathematical approach followed that expressed by Eq. (7) with a deviation smaller than 2%.

VI. CONCLUSION

In this paper an approach to the calculation of the thrust provided by a LSM has been presented. This approach models the induct of a LSM as a series of coils interacting with inductor coils. In the calculation of the total thrust the

deviation of the position of each coil from its ideal position is taken into account by using a probability distribution function. This explicitly gives the dependence of the total thrust from the standard deviation of this distribution function. It has been shown that the reduction of the provided thrust follows an exponential law, which decays as the standard deviation of the distribution function.

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