

Dry deposition processes on urban and suburban surfaces: modelling and validation works

M. Giardina^{a,*}, A. Donato^b, P. Buffa^a, D. Contini^b, A. Cervone^c, C. Lombardo^c, F. Rocchi^c

^a Department of Engineering, University of Palermo, Viale delle Scienze, Edificio 6, 90128 Palermo, Italy

^b Institute of Atmospheric Sciences and Climate, ISAC-CNR, str. prov. Lecce-Monteroni, km 1,2, 73100 Lecce, Italy

^c Italian Agency for New Technologies, Energy and Sustainable Economic Development (ENEA), Via Martiri di Monte Sole, 4, 40129 Bologna, Italy

Abstract

Dry deposition process is one of the important pathways for the removal of particles from atmosphere. It is the result of a combination of different environmental and physical factors as atmospheric conditions, particle properties, characteristics of the canopy. For this latter factor, urban canopy represents unevenly combinations of different types of surface elements that increases the complexity of phenomena that influence deposition processes. Therefore, particle dry deposition on urban surfaces is not easy to configure and, although empirical or semi-empirical models in literature have been developed to address this aspect, there is not standardized and commonly accepted criteria. To overcome this issue, a comparison work has been performed between different models and measurement data by experimental campaigns covering for different surface roughness conditions, i.e. from the patchy Venice surface to the suburban area of Maglie and Lecce, or near urban zone for Bologna. This activity has also included other cities in the world. Giardina et al. (2017b) model, by using two different Brownian diffusion resistance, seems to capture the main dry deposition processes for the examined contexts.

1. Introduction

Urbanization and industrialization are the cause of the formation of pollutants that are deposited on trees, grass, crops, water bodies, and buildings, generating significant environmental impacts. In case of routine emissions, or releases due to nuclear power plant accidents or, more generally, incidents involving the release of hazardous substances a key challenges are the characterization of the specific pollutants that are released as well as studying the dispersion and deposition phenomena. Risk assessment for pollution from hazardous substances is necessary for mitigation strategy formulation and development of policies by decision makers for health protection at regional and national levels.

Dry deposition process, one of the important pathways for the removal of particles from atmosphere, is the result of a combination of different environmental and physical factors as atmospheric conditions, particle properties, characteristics of the canopy. For this latter aspect, the urban canopy represents unevenly combinations of different types of surface elements that increases the complexity of the involved phenomena that influence particle depositions.

Smooth surfaces tend to have lower deposition rates per unit area than rougher surfaces. Such relatively small deposition rates are reported by Roed (1983) for radioactive particles surface deposition as Cs137, on vertical walls, in Denmark. The Author reports the results of nine samples for a brick wall with a range of wet/dry deposition velocities 0.003 to 0.07 cm/s, four samples for a plastered wall with a range of 0.014 to 0.085 cm/s, and only one sample, in an area sheltered from wet deposition, with a dry deposition velocity of 0.003 cm/s.

Nicholson (1987) reported similarly small deposition rates for deposition of Cs134 and Cs137 particles to roof and building materials in England. Although the data set was small, the results were consistent with lower deposition velocities over smoother surfaces.

In (Papastefanou, 2008) an up-to-date summary of knowledge about depositions of radioactive aerosols is provided. Experimental deposition velocities highlighted for Be7 particles variations from 0.1 to 3.4 cm/s, Pb210 from 0.7 to 1.1 cm/s and Cs137 from 1.3 to 6.3 cm/s. These data refer mostly to temperate latitudes of the Northern Hemisphere, e.g. at Oak Ridge, Tennessee (Mahoney, 1984), Norfolk, Virginia (Todd et al., 1989), New Haven, Connecticut (Turekian et al., 1983), Detroit, Michigan (McNeary and Baskaran, 2003), Quillayte, Washington (Creelius, 1981), Munich, Germany (Rosner et al., 1996).

Tai et al. (1999) highlighted that particles less than 10 μm contributed to only a small fraction (<10% in mass) of the dry deposition fluxes, both in urban and non-urban locations. These calculations are extremely sensitive to the end-point particle diameter since the dry deposition velocity for a particle of 1 μm in diameter is approximately 100 times smaller than a 10 μm particle, so the concentration of 10 μm particles is 100 times smaller than the concentration of 1 μm particle, if the flux is the same (Holsen and Noll, 1992; Lin et al., 1994; Lee et al., 1996; Ki et al., 2007). Consequently, particles larger than 10 μm in diameter are

* Corresponding author: Mariarosa Giardina, Department of Engineering, University of Palermo, Italy. E-mail address: mariarosa.giardina@unipt.it.

60 responsible for the majority of dry deposition fluxes, but these particles have a smaller influence on the total
61 airborne concentration (Johnson and Davidson, 2019).

62 In urban areas, measurements of particle vertical turbulent fluxes are a powerful tool to characterize road
63 traffic pollution or, more in general, urban emissions and the dynamics of particles in the urban surface layer
64 (Contini et al., 2012; Deventer et al., 2015), which allows us to better characterize the exchange of particles
65 between a city and the atmosphere. Diurnal evolution of the particle number concentration and flux for small
66 particles (0.25 μm in diameter), as reported in Conte et al. (2018), shows a clear correlation with anthropic
67 activities (mainly road traffic). The particles fluxes were upward on average, even though small net
68 deposition was observed for larger particles. The city is a continuous source of particles of small size. Others
69 studies of urban deposition rates of hydrocarbons and metals show variations that largely reflect the
70 influence of local sources on ambient airborne contaminant concentrations (Azimi et al., 2005).

71 Roupsard et al. (2013) studied dry deposition velocities of submicron aerosols on horizontal and vertical
72 urban surfaces of glass, cement facing by using wind tunnel experiments. On horizontal surfaces, dry
73 deposition velocity varies from 0.0014-0.0045 cm/s on conventional glass to 2.8-1.25 cm/s on synthetic
74 grass.

75 The approach for determining deposition rates has several limitations for urban canopies. Total deposition on
76 the various urban surfaces scales with the available surface areas. Considering the total deposition per unit
77 horizontal area, grass and trees have relatively high deposition rates compared to smooth surfaces. To that
78 extent, modeling the variations in surface deposition require considering interactions of surface roughness
79 and local air circulation.

80 Therefore modelling of dry deposition phenomena over urban canopies is not easy to configure and, although
81 empirical or semi-empirical models have been developed to address this complex aspect, there is not
82 standardized and common accepted criteria proposed in literature (Droppo, 2006). Indeed, their application
83 remains valid for specific conditions and if the data meet all of the assumptions required by the data used to
84 define the models.

85 To overcome this issue, a comparison work has been performed between different deposition velocity
86 models and data from several experimental campaigns, covering different surface roughness conditions, i.e.
87 from the patchy Venice surface to the suburban area of Maglie and Lecce, or near urban zone for Bologna.
88 This activity has also included other cities in the world compared to Italy. Giardina et al. (2017b) model, by
89 using two different Brownian diffusion resistance, seems to capture the main dry deposition processes for the
90 examined urban canopies.

91

92 **2. The main phenomena in dry deposition processes**

93

94 In atmospheric models, the Surface Layer (SL) is the air layer over the surface whose properties are largely
95 controlled by the local surface fluxes. The strict definition of the surface layer is a fully turbulent layer over
96 homogenous surfaces under steady-state conditions. With this surface layer, a second layer is designated that
97 refers to the laminar, or near-laminar, flow that occurs immediately over the surfaces. This layer, in the
98 literature, is also referred to as the “quasi-laminar sublayer” or “deposition layer”.

99 The main transport processes, that occur within the above layers, are:

- 100 • transport due to atmospheric turbulence in SL. This process is independent of the physical and
101 chemical nature of the pollutant and it depends only the turbulence level;
- 102 • diffusion in the thin layer of air which overlooks the air-ground interface (i.e. quasi-laminar sublayer).
103 The dominant component becomes molecular diffusion for gasses, Brownian motion for particles and
104 gravity for heavier particles;
- 105 • transfer to the ground. This component exhibits a pronounced dependence on surface type with which
106 the pollutant interacts (i.e. urban context, grass, forest, etc..).

107 Therefore, dry deposition phenomena involve three sequential sets of processes:

- 108 1. through the turbulent SL, the particle moves by the combined effects of eddy diffusion (i.e., carried
109 by turbulent movements of air) and gravity.
- 110 2. in quasi-laminar surface layer, the particle can reach the surface by molecular diffusion, interception,
111 or impaction.
- 112 3. near the surface, retention or rebound depends on a combination of surface and impact properties.

113 It is worth to note that the deposition process changes quite a lot over the year, for example due to the
114 seasonal variation of vegetation (with or without leaf) or over the day in connection with meteorological
115 conditions (e.g. influence of temperature on leaf stoma).

116

117 **2.1 Transport for atmospheric turbulence in SL**

118

119 Eddy diffusion refers to the transport resulting from turbulent movements in the air that play a pivotal role in
120 determining the vertical transfer of momentum, heat and mass in the Atmospheric Boundary Layer (ABL),
121 that usually encompasses the lowest tens to hundreds of meters in the atmosphere over the earth's surface
122 (Garratt, 1992). It is widely studied in boundary-layer meteorological modeling, however its impact on dry
123 deposition has not been well understood and characterized.

124 In approximately the lowest 10% of the ABL (layer SL), the vertical fluxes of transferred quantities are
125 nearly constant with height and can be represented quite successfully by formulations based on the Monin-
126 Obukhov (M-O) similarity theory (Stull, 1988; Högström, 1988, 1996; Foken, 2006).

127 Under the assumption of steady state between generation, dissipation and transport of the turbulent eddies
128 across the SL, the M-O similarity theory describes relationships between vertical profiles and fluxes for
129 momentum and scalar quantities (i.e., heat and trace constituents), using a metric called Monin-Obukhov
130 length, L (Garratt, 1992):

131

$$132 \quad L = \frac{u_*^3 c_p \rho_a \bar{T}}{k g H} \quad (1)$$

133

134 where \bar{T} is the average temperature in SL (K); ρ_a air density (g/cm^3); c_p specific heat at constant pressure
135 [J/kg K]; g gravitational constant (cm/s^2); H sensible heat (W/m^2); k von Karman constant set to 0.4 .

136 The parameter L is a measure of atmospheric stability. It varies from small negative values in extremely
137 unstable atmospheric conditions to negative infinity, as the atmospheric stability approaches neutral from
138 unstable (neutral conditions $L = \infty$). In extremely stable conditions, L is small and positive.

139

140 **2.1.1 Friction velocity and wind speed**

141

142 The friction velocity is derived from the atmospheric boundary layer similarity theory proposed by Monin
143 and Obukhov (1954). The basic hypothesis of the similarity theory is that a number of parameters in the
144 atmospheric layer near the ground, including the wind profiles, should be universal functions of the friction
145 velocity, the length L , and the height z above ground.

146 Usually, the ratio z/L is related to atmospheric stability. When z/L is negative and large the atmosphere is
147 extremely unstable, if z/L is near zero, the atmosphere is neutral, if z/L is positive and large the atmosphere
148 is extremely stable.

149 Under the assumption of continued validity of the M-O flux-gradient relationships in the interfacial sublayer,
150 it is possible to evaluate wind speed profile as follows:

151

$$152 \quad \frac{u(z)}{u_*} = \frac{1}{k} \left[\ln \left(\frac{z}{z_0} \right) - \Psi_h(z) \right] \quad (2)$$

153

154 with $u(z)$ wind speed at deposition reference height (m/s); u_* friction velocity, (m/s); k von Karman constant;
155 z deposition reference height (m); z_0 surface roughness length (m); Ψ_h is a term that takes into account
156 effects of stability on the wind profile.

157 To calculate the stability function Ψ_h in Eq. (2), Brandt et al. (2002) suggested the following relationship:

158

$$159 \quad \Psi_h = -5 \frac{z}{L} \quad \text{with } \frac{z}{L} > 0 \text{ (stable conditions)} \quad (3)$$

160

$$161 \quad \Psi_h = e^{\left\{ 0.598 + 0.390 \ln \left(-\frac{z}{L} \right) - 0.09 \left[\ln \left(-\frac{z}{L} \right) \right]^2 \right\}} \quad \text{with } \frac{z}{L} < 0 \text{ (unstable conditions)} \quad (4)$$

162

163 Under neutral atmospheric stability, Eq. (2) can be simplified as:

164

$$165 \quad \frac{u(z)}{u_*} = \frac{1}{k} \ln \left(\frac{z}{z_0} \right) \quad (5)$$

166

167 For an aerodynamically rough but relatively flat surface, an extrapolation of the mean wind speed profile
168 downward shows that it reaches zero at some distance above the physical surface. The height at which this
169 occurs is called roughness length, z_0 . The roughness length is positively correlated with the physical

170 roughness of the surface, although a strict functional relationship between physical roughness measures and
171 z_0 do not exist.

172 The surface roughness length over land depends on the surface cover and land use and is often difficult to
173 estimate (WMO, 2010). A way of determining z_0 is by a visual survey of the terrain around the wind station
174 with the help of the table below, the validity of which has been corroborated (Davenport et al., 2000). Tab. 1
175 reports Davenport classification of terrain roughness, z_0 , for different urban configurations. Other studies
176 have found that the surface roughness length tends to be about one-tenth of the dimensions of the surface
177 elements (Sutton, 1953; Stull, 1988).

178 In environments such as over agricultural crops and forest and urban canopies, the practice is to displace the
179 entire velocity profile upward such that the height at which velocity profiles reach zero is the sum of a
180 canopy roughness length, z_0 , and displacement height, d , defined as zero plane displacement. Fig. 1 reports a
181 generalized mean wind velocity profile in a developed urban area and the zero-plane displacement length, d ,
182 for this configuration.

183 Consequently, Eq. (5) is modified as follows:

184

$$185 \frac{u(z)}{u_*} = \frac{1}{k} \ln \left(\frac{z-d}{z_0} \right) \quad (6)$$

186

187 **2.2 Transport to Brownian motions**

188

189 Brownian motion is usually assumed to dominate the diffusion processes in the quasi-laminar surface layer.
190 However, there is the possibility that phoretic forces can locally influence dry deposition fluxes.

191 Particles in the range of 0.001 to 0.1 (μm) (ultrafine particles) move like gaseous molecules in flowing air
192 (i.e., they exhibit rapid random Brownian motion). The motion causes them to collide with any nearby
193 surfaces. Ultrafine particles tend to adhere to these surfaces as the result of intermolecular forces. This
194 mechanism tends to be an effective deposition process with very small particles depositing at rapid rates on
195 the nearest available surfaces. Under some circumstances, this diffusion mechanism can continue to be the
196 dominant deposition process for particles $>0.1 \mu\text{m}$.

197

198 **2.3 Gravitational settling**

199

200 Gravitational settling is the downward motion of particles that results from the gravitational attraction and it
201 is the dominant process for the dry deposition of the larger particles of about $10 \mu\text{m}$.

202 Particulate sizes, densities, and shapes largely define gravitational settling rates.

203 The settling velocity for particles, v_s , can be computed using a modified form of Stokes Law (Hanna et al.,
204 1982):

205

$$206 v_s = \frac{C_c d_p^2 g (\rho_p - \rho_a)}{18 \mu_a} \quad (7)$$

207

208 where ρ_p is the particle density, (g/cm^3); μ_a is the air dynamic viscosity ($\text{g}/\text{cm s}$); d_p is the particle diameter,
209 (cm); and C_c is the Cunningham factor.

210 The Cunningham factor can be expressed as (Seinfeld and Pandis, 1998):

211

$$212 C_c = 1 + \frac{\lambda_a}{d_p} \left(2.514 + 0.8 e^{-\frac{0.55 d_p}{\lambda_a}} \right) \quad (8)$$

213 where λ_a is the mean free path of air (cm).

214 Non-spherical particles fall at slower rates. For materials with equivalent densities, the change in settling
215 velocities is less than 30% for ellipsoid and cylinder shapes. To account for shape effects, an
216 aerodynamically equivalent diameter (Jennings and Parslow, 1988; Reid et al., 1994) is frequently used to
217 define the settling velocity of a particle.

218

219 **3. Urban deposition models**

220

221 An urban area represents a complex context for assessment of potential exposures from an atmospheric
222 release and, in this field, dry deposition models are too simple for application to the urban environment

223 (Cherin et al., 2015). Their classical approaches (Wesely and Hicks, 2000; Petroff et al., 2008), which are
224 inherited from semi-empirical models, were developed for deposition over vegetated surfaces, bare soil or
225 water, and therefore they fail to represent the complexity of the dry deposition processes over an urban
226 canopy.

227 A review of dry (and wet) deposition computational methods was conducted for radioactively contaminated
228 particles (in the range 0.1 to 10 μm) by the Atmospheric Dispersion Modeling Liaison Committee (NRPB,
229 2001). They are recommend values and methods to estimate deposition rates and special parameter limits to
230 extrapolate the dry deposition model to an urban environment. However, neither of these reviews addressed
231 the issues of the applicability of the dry deposition models to non-ideal conditions such as the
232 aerodynamically very rough surfaces encountered in an urban environment.

233 Resistance-based approaches are widely used as a basis for dry deposition formulations. This approach,
234 explained in more detail below, has the advantage of providing a means of combining a number of the
235 processes controlling dry deposition into a single formulation.

236 In one of the early implementations, Sehmel and Hodgson (1978) proposed an empirical model based on
237 curve fits to wind tunnel deposition results for a range of soil surface covers. Their model combined
238 empirical data with the theory for molecular diffusion of very small particles and gravitational settling rates
239 for larger particles. The Authors also demonstrated the importance of considering the density of the particles
240 in the dry deposition computation. Subsequent applications have included air quality (e.g., chemicals and
241 trace metals), health physics (radionuclides), and acid rain models.

242 Detailed models, that address the processes leading to exposures in an urban environment, have been
243 developed for radiological exposures (Jones et al., 2006).

244 Eged et al. (2006) used a Monte Carlo approach to evaluate potential radiological doses in urban
245 environments. The results show that urban dose computation models provide some results that are the same
246 and some that are not.

247 Based on their experimental results, Chen et al. (2012) developed a relationship between TSP (Total
248 Suspended Particulate Matter) deposition velocity and meteorological parameters. The experimental
249 campaigns were conducted in locations near Guangzhou, China, during dry season. The dry deposition
250 velocity equation is a function of the wind speed, the relative humidity, and air temperature.

251 Noll et al. (2001) correlate the particle deposition velocity with Stokes settling velocity, friction velocity,
252 dimensionless inertial deposition velocity, and dimensionless Brownian diffusion deposition velocity by
253 using 20 atmospheric samples collected at flow Reynolds numbers ranging from 9000 to 30,000 and related
254 to particle size of range 1–100 μm . Some samples were collected on the roof of a four-story building (12 m
255 height) located in a mixed institutional, commercial, and residential area on the south of Chicago. A model
256 was developed using a least square procedure to fit a sigmoid curve to ambient data, similarly to one
257 developed for deposition in a vertical pipe (Muyschondt et al., 1996).

258 A parameterization of particle dry deposition has been developed in (Zhang et al., 2001) for different
259 surfaces and meteorological variables. It includes deposition processes, such as, turbulent transfer, Brownian
260 diffusion, impaction, interception, gravitational settling, particle rebound and also particle growth under
261 humid conditions. The parameters in dry deposition velocity calculations are different on the basis of 12 land
262 use categories, including urban area, and five seasonal categories.

264 **3.1 Resistance approach to describe dry deposition processes**

266 In mechanistic or process-based dry deposition models, an electrical resistance-based approach is widely
267 used to parameterize the dry deposition velocity (Venkatram and Pleim, 1999; Zhang et al., 2001; Giardina et
268 al. 2017a; Giardina and Buffa, 2018).

269 By considering that the reciprocal of the dry deposition velocity, v_d , is the overall resistance to the mass
270 transfer, the influence of the various phenomena on the deposition velocity can be expressed in terms of an
271 electrical analogy.

272 Giardina et al. (2017b) proposed a new approach to evaluate the total resistance r_t for urban rough surfaces.

273 In SL the aerodynamic resistance, r_a , is connected in series with the resistance across the quasi-laminar
274 sublayer, r_{ql} , to take into account mechanisms of diffusion by Brownian motion and turbulent phenomena
275 (Fig. 2). Accordingly, the overall resistance, r_t , is evaluated by the using the following equations:

$$277 \quad r_t = r_a + r_{ql} \quad (9)$$

278 Relationships for aerodynamic resistances r_a are based on surface layer parameterizations from M-O
 279 similarity theory:
 280

$$281 \quad r_a = \frac{\ln\left(\frac{z}{z_0}\right) - \psi_h}{ku_*} \quad (10)$$

282
 283 In Eq. (10), Eq.s (3) and (4) can be used to calculate the stability function Ψ_h .
 284 The resistance r_{qi} is evaluated by considering two resistances in parallel, that is: resistance r_{bd} , which
 285 represents the Brownian diffusion process; and resistance r_i , which allows to treat turbulent and impaction
 286 phenomena. The resistance r_i is evaluated by considering two resistances in series: r_{ii} to take into account
 287 inertial impact conditions; r_{ti} to consider effects resulting from turbulent impaction. This last assumption
 288 allows to take into consideration the reciprocal influence of the two impact processes on dry deposition
 289 efficiency.

290 The Authors assumed the expression reported in Eq. (11) for the resistance r_{bd} as suggested in (Wesely and
 291 Hicks, 2000; Paw, 1983; Hicks et al., 1987; Pryor et al., 2009; Kumar and Kumari, 2012; Giardina and
 292 Buffa, 2018):
 293

$$294 \quad r_{bd} = \frac{1}{u_* Sc^{-2/3}} \quad (11)$$

295 where Sc is the Schmidt number evaluated as follows:
 296

$$297 \quad Sc = \frac{v_a}{D} \quad (12)$$

298 with v_a the air kinematic viscosity (m^2/s) and D the Brownian diffusivity of air (m^2/s).
 299 To evaluate the resistance for inertial impact process r_{ii} , it is used the following relationship valid for rough
 300 surfaces:
 301

$$302 \quad r_{ii} = \frac{1}{u_* \left(\frac{St^2}{St^2+1}\right) R} \quad (13)$$

303 where St is the Stokes number:
 304

$$305 \quad St = \frac{u_*^2 v_s}{v_a} \quad (14)$$

306
 307 The parameter R in Eq. (13) is the particle rebound evaluated as follow:
 308

$$309 \quad R = e^{(-b\sqrt{St})} \quad (15)$$

310 where b is assumed equal to 2, as suggested in (Nemitz et al., 2002)
 311 Empirical relations of turbulent deposition are typically presented in terms of the dimensionless particle
 312 relaxation time τ_+ :
 313

$$314 \quad \tau_+ = \tau \frac{u_*^2}{v_a} \quad (16)$$

315 where τ is the particle relaxation time defined, for a spherical particle, as follows:
 316

$$317 \quad \tau = \frac{d_p^2 \rho_p C_c}{18\mu_a} \quad (17)$$

318
 319 Various models predict a functional dependence of resistance turbulent impact phenomena, r_{ti} , on τ_+ as
 320 follows:
 321

$$322 \quad r_{ti} = \frac{1}{u_* m \tau_+^n R} \quad (18)$$

323 Various works have been performed to evaluate m and n parameters, and these constants have been revised
 324 by fitting some data reported in literature for roughness surfaces (Giardina et al., 2018). The results were
 325 $m=0.1$ and $n=0.5$.

326 Venkatram and Pleim (1999) proposed the following expression for dry deposition velocity of particles that
 327 is consistent with the mass conservation equation:

$$328 \quad v_d = \frac{v_s}{(1 - e^{-r_t v_s})} \quad (19)$$

330 In the light of the above considerations about the total resistance r_t , Eq. (19) can be rewritten as follows:
 331

$$332 \quad v_d = \frac{v_s}{1 - e^{-\left\{v_s \left[r_a + 1 / \left(\frac{1}{r_{db}} + \frac{1}{r_{ii} + r_{ti}} \right) \right] \right\}}} \quad (20)$$

334 where r_{db} , r_{ii} , and r_{ti} are evaluated by Eq.s (10), (11), (13), and Eq. (18) with $m=0.1$ and $n=0.5$.

335

336 **3.2 Considerations about Brownian diffusion resistance**

337

338 In literature, various models have been developed to predict functional dependence of Brownian diffusion
 339 resistance, r_{bd} , from Sc number and friction velocity, u_* , however they cannot be extended for urban
 340 conditions.

341 Some Authors defined a Brownian diffusion resistance dependency on roughness Reynolds number defined
 342 as $Re_* = u_* z_o / \nu$. For example, in (Chamberlain et al., 1984) experiments were carried out to study diffusive
 343 transfer to bluff roughness elements that can be compared to urban and suburban conditions.

344 Radioactive gases and labelled particles were used in wind tunnels to measure the effects of Re_* and
 345 Sc numbers on transport to surfaces with widely spaced roughness elements. It is highlighted that deposition
 346 of larger particles is dominated by the effects of bounce off, which depends on surface conditions.

347 In this field, research activities have been carried out to define a new formulation that is capable to perform
 348 predictions for different typologies of urban area and small particle diameters (Giardina et al., 2018).

349 Fitting approach was performed on parameters that associate Brownian diffusion resistance to Re_* and Sc .
 350 This analysis was carried out by using experimental data of dry deposition velocity for different rough
 351 surfaces (Chamberlain et al., 1984, Pryor et al., 2009; Donateo and Contini, 2014).

352 The following relationship was obtained:

353

$$354 \quad r_{bd} = \frac{1}{u_* Sc^{-0.5} Re_*^{-0.05}} \quad (21)$$

355

356 **4. Validation works of dry deposition velocity modeling for Italian cities**

357

358 The validation works have been performed by using experimental campaigns carried out by researchers from
 359 Institute of Atmospheric Sciences and Climate (ISAC), Italian National Research Council (CNR), unit of
 360 Lecce, covering for different surface roughness conditions of urban-suburban Italian areas (Donateo and
 361 Contini, 2014).

362 For the experiments, performed between 2004s and 2009s, the displacement height, d , and the roughness
 363 length, z_o , are reported in Tab. 2. These data have been evaluated from micrometeorological measurements,
 364 following a method reported in (Toda and Sugita, 2003) which uses similarity relationship for sonic
 365 temperature and vertical wind component.

366 Measured dry deposition velocities have been compared with predictions of models reported in (Giardina et
 367 al., 2017b), (Chen et al., 2012), (Noll et al., 2001), and (Zhang et al., 2001), by adopting the roughness
 368 lengths reported in Tab. 2. For Chen et al. (2012) model, average values of wind velocity, temperature and
 369 relative humidity measured during the experimental campaigns were used, where available. The model of
 370 Zhang et al. (2001) was applied by using the parameters imposed for urban land use categories, as for
 371 roughness length, z_o , imposed by the authors to 1 m.

372 In all figure captions, NNR index (Attilio et al., 1993), which allows comparisons of model performances
 373 independently of the shape of the set of measured data, was evaluated as follows:

374

375
$$\text{NNR} = \frac{\sum_i (1 - \hat{k}_i)^2}{\sum_i \hat{k}_i} \quad (22)$$

376 where

377

378
$$\hat{k}_i = e^{-|\ln k_i|} \quad (23)$$

379

380
$$k_i = \frac{v_{dp_i}}{v_{do_i}} \quad (24)$$

381

382 being v_{dp_i} and v_{do_i} predicted and observed deposition velocities, respectively.

383 NNR index is the normalized mean square error of the distribution of the normalized ratios \hat{k}_i compared with
384 the ideal distribution $\hat{k}_i = 1$.

385

386 **4.1 Maglie suburban site (South-Eastern Italy)**

387

388 The measurement site, located in NE boundary of the town of Maglie (LE) in the Apulia region of Italy
389 (40°07'38.39''N, 18°17'59.50''E), can be considered a suburban site influenced by an industrial area. The
390 town is extending mainly in the sector of wind direction between SE and SW and the country side is in the
391 sector between NNO and E. The site is characterized by presence of small buildings (1-2 floors) and roads
392 with relatively high traffic volume (Donateo et al., 2006).

393 PM_{2.5} monitoring campaigns were performed mainly from November to January in 2004, 2006 and 2007.
394 Fig.s 3 and 4 show the deposition velocity as function of friction speeds u_* and unstable and stable
395 conditions, respectively. Comparisons between experimental data reported in (Donateo and Contini, 2014),
396 carried out by ISAC-CNR unit of Lecce, and predictions obtained by applying models reported in (Giardina
397 et al., 2017b), (Chen et al., 2012), (Noll et al., 2001), and (Zhang et al., 2001), for particle diameter d_p of 2.5
398 μm , are shown.

399 Giardina et al. (2017b) model has been tested by using Eq. (20) and relationships reported in Eq.s (11) and
400 (21) for the Brownian diffusion resistances.

401 Analyzing results, Noll et al. (2001) and Zhang et al. (2001) models overestimate experiments with wind
402 friction velocities higher than 0.2 m/s, whereas Chen et al. (2012) model overestimates experiments for all
403 wind friction velocities. Giardina et al. (2017b) model shows a good agreement with the experimental data if
404 it is applied with the Brownian diffusion resistances reported in Eq.s (11) and (21), and the two predictions
405 result very close both for unstable and stable conditions. This result is confirmed by NNR index values.

406

407 **4.2 Venice urban site (North-Eastern Italy)**

408

409 Measurements were performed at a background site placed on Venice lagoon (Donateo et al., 2012). The
410 measurement site (45°29'09.5''N, 12°24'12.7''E) was a field located at about 8 km NE of the Venice town.
411 This site was located very close (about 5 m) to the water lagoon at the W-SW side, while, in the other
412 directions (selected for this work), it was characterized by land for about 1-2 km with short vegetation, some
413 small trees, and one or two-floor houses, although channels and water were also present in this area.

414 Monitoring campaigns of PM_{2.5} concentrations and flux measurements were performed during the summer
415 2004, winter 2005 and spring 2006.

416 Fig.s 5 and 6, that report the deposition velocity as function of friction speed, refer to instability and stability
417 atmospheric conditions, respectively.

418 These figures shown comparisons among experimental data carried out by ISAC-CNR and elaborated in this
419 work, and predictions obtained by using models reported in (Giardina et al., 2017b), (Chen et al., 2012),
420 (Noll et al., 2001), and (Zhang et al., 2001), for particle diameter d_p of 2.5 μm .

421 (Giardina et al., 2017b) model has been tested by using Eq.s (11) and (21).

422 Very high differences can be highlighted between dry deposition velocities experimental data and the
423 predictions of Zhang et al. (2001) and Noll et al. (2001) models, both for unstable and stable conditions and
424 wind friction velocities higher than 0.2 m/s. Chen et al. (2012) model overestimates the observations data for
425 all wind friction velocities.

426 The predictions obtained by Giardina et al. (2017b) model, applied with Eq.s (11) and (21), show a good
427 agreement with experimental data.

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4.3 Bologna industrial district (Central Italy)

The measurement site was near the incinerator plant for the city of Bologna, (44°31'17.59''N, 11°25'53.48''E). The data refer to experimental campaigns performed in summer 2008 and winter 2009.

Fig.s 7 and 8 report the deposition velocity experimental data as function of friction velocity for instability and stability atmospheric conditions, respectively.

Comparisons between experiments reported in (Donateo and Contini, 2014), and results obtained by using models reported in (Giardina et al., 2017b), (Chen et al., 2012), (Noll et al., 2001), and (Zhang et al., 2001) are shown.

The experiments reported in (Donateo and Contini, 2014) were carried out from June to July 2008, and January to March 2009 for particle diameter d_p of 0.045 μm .

High differences can be highlighted between dry deposition velocities experimental data and the predictions of Noll et al. (2001) and Chen et al. (2012) models.

As NNR index values pointed out, the predictions of Zhang et al. (2001) and Giardina et al. (2017b) models, this last applied with Eq. (21), show very good agreement for all experiments and unstable conditions (Fig.7). For stable conditions (Fig. 8) the best predictions are attributed to (Giardina et al., 2017b) model applied both with Eq.s (11) and (21).

4.4 Lecce suburban site

Experimental campaigns were performed during spring/summer of 2005. $\text{PM}_{2.5}$ concentrations and fluxes were measured.

The site was the experimental field of the Lecce Unit of ISAC-CNR placed inside the University Campus (40°20'10.8''N, 18°07'21.0''E) and located at about 3.5 km SW from the town of Lecce. It is a rectangular field with a major side of about 200 m characterized by short vegetation, with two contiguous sides surrounded by small trees.

The urban background area is characterized for at least 1 km in all directions by the presence of patches of trees (8–10 m tall) and small two-story buildings and some roads with no industrial releases nearby.

Due to the proximity of urban areas, the site can be categorized as an urban background area. Measurements were taken at 10 m above the ground.

Fig.s (9) and (10) show the deposition velocity as function of friction speeds for unstable and stable conditions reported in (Donateo and Contini, 2014), and predictions obtained by applying the models reported in (Giardina et al., 2017b), (Noll et al., 2001), and (Zhang et al., 2001).

Analyzing results, Noll et al. (2001) model allows very high underestimations rates of experimental data with friction velocities higher than about 0.2 (m/s), both unstable and stable conditions.

Giardina et al. (2017b) model, applied with Eq.s (11), and (21) show a good agreement with all experimental data. The curves result very close each other, both for unstable and stable conditions.

5. Validation work of dry deposition velocity on United States urban areas

Dry deposition phenomena modelling on urban canopies is limited and, although empirical or semi-empirical models have been developed to address this complex issue, there is no universal acceptance criteria. Therefore, it can be said that the models proposed in literature are not capable of representing particle dry deposition for several categories of pollutants and different urban surface conditions.

In order to overcome this issue, validation works have been extended to experimental campaigns on different urban canopies compared to Italy.

In this section the results obtained on typical United States urban areas are shown and discussed.

It is to be noted that the roughness length z_0 for the examined urban areas were evaluated by using Devenport classification reported in Tab. 1 if not reported by the Authors. (Zhang et al., 2001) model was applied by using the parameters imposed for urban land use categories and $z_0 = 1$ m as suggested by the Authors.

5.1 Experimental data of Noll et al. (2001)

Four sampling campaigns conducted at Chicago are described in (Noll et al., 2001). The data were collected on the roof of a four-story building (12 m height) located in a mixed institutional, commercial, and residential area on the south of Chicago.

485 The building is located on the IIT (Illinois Institute of Technology) campus, which is located 5.6 km south of
486 Chicago's center and 1.6 km west of Lake Michigan. The IIT campus consists of predominately low rise
487 buildings, landscaped areas, and asphalt parking lots. The atmospheric particle mass size distribution and dry
488 deposition flux were measured simultaneously with a wide range aerosol classifier (WRAC) and a smooth
489 greased surface.

490 The 20 sets of atmospheric measurements that were used to develop Noll et al. (2001) model were grouped
491 into 4 Reynolds number classes of 9,000–13,000 (class I), 13,000–17,000 (class II), 17,000–21,000 (class
492 III), and >21,000 (class IV). The average wind velocities for the 4 classes were 3.42, 4.64, 5.82, and 7.87
493 m/s.

494 Fig. (11) reports the experimental dry deposition velocity as function of particle diameters in log-log graph.
495 Comparisons with predictions obtained by using models reported in (Giardina et al., 2017b), (Noll et al.,
496 2001) and (Zhang et al., 2001) are also shown.

497 The models reported in (Giardina et al., 2017b) and (Noll et al., 2001) were applied by using the class “Very
498 rough” for the terrain roughness reported in Tab. 1 ($z_0=0.5$). Moreover Noll et al. (2001) model is used with
499 Reynolds number of 15000. The friction velocity was imposed equal to 20 cm/s.

500 As expected, Noll et al. (2001) predictions show a good agreement with the experimental data reported in
501 Fig. (11), while the model of Giardina et al. (2017b), applied with Eq.s (11) and (21), underestimate the
502 experimental data if the particle diameter is greater than 10 μm .

503 (Zhang et al., 2001) predictions give a good agreement with the experimental data. However, it shows the
504 highest deposition velocity values among models' applications for particle diameters d_p less than about 10
505 μm .

506

507 **5.2 Experimental data of McNeary and Baskaran (2003)**

508

509 The depositional fluxes in the bulk and dry fallout as well as the concentrations of Be7 and Pb210 in aerosols
510 were measured for a period of 17 months at Detroit, Michigan. The concentrations contributed 2.1–19.8%
511 and 3.6– 48.6% of the bulk depositional fluxes.

512 The bulk depositional fluxes varied between 3.11 and 63.0 dpm/cm²yr and 0.35 and 10.3 dpm/cm²yr,
513 respectively, and this variability is attributed to the frequency and amount of precipitation and seasonal
514 variations in the depositional fluxes.

515 The sampling site is one of the air monitoring network stations operated by the Wayne County Air Quality
516 Management Division and jointly operated by the Wayne County and the Michigan Department of
517 Environmental Quality (MDEQ), under cooperative agreement with the U.S. Environmental Protection
518 Agency (EPA).

519 A bulk rain collector (200-L polyethylene drum with surface area of 2800 cm²) was deployed in September
520 1999 at a site in the southwest area of Detroit, Michigan (42° 250' N; 83° 10' W; 175 m above mean sea
521 level) at about 1 m above the ground to prevent the resuspension of dust particles getting into the collector.
522 The lid of the bulk collector was deployed as the dry collector in October 1999 on the roof of a building at
523 the same site at about 4 m above ground.

524 The bulk rain samples were collected after each major precipitation event or once in about a month and after
525 about 10 days of dry weather for the dry collector.

526 The models reported in (Giardina et al., 2017b) and (Noll et al., 2001) were applied by using the class
527 “Skimming” for the terrain roughness reported in Tab. 1.

528 Fig.s (12), (13) and (14) report comparisons between dry deposition experimental data, during the period
529 from October 1999 to January 2001, and predictions of models proposed in (Giardina et al., 2017b), by using
530 Brownian diffusion resistance Eq.s (11) and (21), and (Zhang et al., 2001), for particle diameters $d_p = 0.1$ and
531 1 μm , respectively.

532 The wind speeds used for model applications are related to measurements performed by meteorological
533 station located near Detroit airport (Windsor) for the examined periods. This station is the closer site to the
534 dry deposition flux measurement point.

535 As we can see in all figures, as expected the calculated dry deposition rate increases with increasing of the
536 particle diameter.

537 For some periods, (Giardina et al., 2017b) model with Eq. (11) and particle diameter of 0.1 μm (Fig. 12)
538 underestimates the dry deposition velocity, whereas the predictions are improved if the model is used with
539 Eq. (21) (Fig. 13).

540 The predictions, obtained by using (Zhang et al., 2001) model and reported in Fig. 14, overestimate the
541 experimental data for particle diameter of 1 μm and allow a good agreement for particle diameter of 0.1 μm .

542 It should be highlighted that the experimental measurements were carried out also during rainy days, so these
543 weather conditions can only increase deposition processes compared to dry deposition phenomena.
544 Therefore, the above described results shall be assessed also in relation to this last consideration.

545 **6. Validation work of dry deposition velocity on France suburban-urban areas**

546 Experimental campaigns reported in (Connan et al., 2018) were conducted from December 2010 to August
547 2011 in the car park of Cherbourg, France (49.6348 ° N-1.6456 ° W), on a site characteristic of an urban
548 zone. The sonic anemometer was placed 1 m from the ground. The friction wind u_* was calculated from the
549 measurements at the time of the experiment that was taken systematically (i.e. over a period of about 10 to
550 30 min).

551 Measurements made for different relative contact surfaces. The classical sample with a normal density was
552 defined as a grass with a relative contact surface (RCS) of 1; samples with 2 times smaller RCS (RCS = 1/2)
553 and 3 times smaller RCS (RCS = 1/3) were also prepared by removing 1 strand of grass out of 2 or 1 strand
554 of grass out of 3 for a given area (20 x 20 cm).

555 The experimental results are reported in terms V_d/u_* parameter as function of particle diameters as shown in
556 Fig. 15.

557 Experimental results show that v_d/u_* varies in a range of 2 to 6 10^{-3} for particle diameters from 0.2 to 1.2
558 μm . Beyond 1.2 μm the v_d/u_* ratio increases rapidly with a value of 5.4 10^{-2} to 7.8 μm .

559 Fig. (15) also shows the predictions obtained by applying the models reported in (Giardina et al., 2017b) and
560 (Zhang et al., 2001). Giardina et al. (2017b) model has been tested by using Eq. (20) and relationships
561 reported in Eq.s (11) and (21) for the Brownian diffusion resistances.

562 Good predictions are obtained by using (Giardina et al., 2017b) model, with Eq. (21), and (Zhang et al.,
563 2001) model, even if this last shows the best performance.

564

565 **7. Conclusion**

566

567 Dry deposition process is the result of a combination of different factors as atmospheric conditions, particle
568 properties, characteristics of deposition methods as natural surfaces, urban area and water surfaces. Among
569 the many possible configurations, urban canopy represents an uneven combinations of different types of
570 surface elements that increases the complexity of the involved phenomena.

571 Quantifying the amount of dry deposition is critical since these deposition processes determine the pollutant
572 species' lifetime in air and their input to various ecosystems.

573 The approach to determine deposition rates has several limitations for urban area. Total deposition on the
574 various urban surfaces scales with the deposition rates and available surface areas. Considering the total
575 deposition per unit horizontal area, grass and trees have relatively high deposition rates compared to smooth
576 surfaces.

577 Therefore, modeling the variations in surface deposition requires to considere interactions of surface
578 roughness and local air circulation. Dry deposition models are too simple for application to the urban
579 environment (Cherin et al., 2015). Their classical approaches (Wesely and Hicks, 2000; Petroff et al., 2008),
580 which are inherited from semi-empirical models, were developed for deposition over vegetated surfaces, bare
581 soil or water, and therefore they fail to represent the complexity of the dry deposition processes over an
582 urban canopy.

583 It follows that the modelling of dry deposition phenomena within urban canopies is not easy to configure
584 and, although empirical or semi-empirical models have been developed to address this complex aspect, there
585 is not a standardized and common accepted criteria proposed in literature (Droppo, 2006). Indeed, their
586 application remains valid for specific conditions and if the data in that application meet all of the
587 assumptions required by the models.

588 To overcome this problem, a validation work of different models has been performed. For this activity, we
589 used the data obtained in experimental campaigns carried out by researchers from ISAC-CNR unit of Lecce,
590 covering for different surface roughness conditions, but also in other parts of the world compared to Italy.

591 In particular (Giardina et al., 2017b), (Chen et al., 2012), (Noll et al., 2001), and (Zhang et al., 2001) models
592 were tested.

593 Giardina et al. (2017b) model, by using Eq.s (11), (21) for the Brownian diffusion resistances, seems to
594 capture the main dry deposition processes for the examined urban canopies. However the use of Eq. (21)
595 shows the best agreements for Italian cities, if NNR index is investigated. Moreover, this latter approach can

596 improve deposition velocity predictions for particles with diameter smaller than 1 μm , as shown by
597 comparisons with the experimental data of (Connan et al., 2018) as reported in Fig. 15.
598 A good agreement is also obtained by using Zhang et al. (2001) model, especially for particles diameter
599 greater than 2.5 μm and friction velocity of about 0.2 m/s, as shown by dry deposition velocity predictions of
600 experiments reported in (Noll et al., 2001) (Fig. 11). However, overestimations are obtained if it is used for
601 Italian urban areas and friction speeds higher than 0.2 m/s, save for suburban areas such as Bologna and
602 particle diameter of 0.045 μm (Figs 7 and 8).
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Declaration of interests

The authors declare that they have no known competing financial interests or personal relationships that could have appeared to influence the work reported in this paper.

The authors declare the following financial interests/personal relationships which may be considered as potential competing interests:

Highlights (for review)

- a pathway to remove pollutants in atmosphere is the dry deposition process
- modelling of dry deposition phenomena over urban canopies is performed
- comparisons were performed between deposition velocity models and experiments covering for different surface roughness conditions
- a new parametrization of Brownian diffusion resistance seems to capture the main dry deposition processes for the examined urban canopies.

Table 1. Davenport classification of terrain roughness (Davenport et al., 2000).

Class	z_0 (m)	Landscape description
Roughly open	0.10	Moderately open country with occasional obstacles (e.g. isolated low buildings or trees) at relative horizontal separations of at least 20 obstacle heights.
Rough	0.25	Scattered obstacles (buildings) at relative distances of 8 to 12 obstacle heights for low solid objects (e.g. buildings).
Very rough	0.5	Area moderately covered by low buildings at relative separations of 3 to 7 obstacle heights and no high trees.
Skimming	1.0	Densely built-up area without much building height variation.
Chaotic	2.0	City centres with mix of low and high-rise buildings.

Table 2[Click here to download Table: Table 2_docx format.docx](#)

Table 2. Summary of experimental sites for aerosol sampling performed by ISAC-CNR. Measurement height (z), displacement height (d), and roughness length (z_0) are reported.

Site	Height z (m)	Displacement height d (m)	Roughness length z_0 (m)
Venice	10	5.1 ± 0.5	0.11 ± 0.03
Bologna	10	4.8 ± 0.5	0.35 ± 0.02
Lecce	10	6.1 ± 0.4	0.53 ± 0.02
Maglie	10	6.0 ± 0.5	0.52 ± 0.02

Figure 1

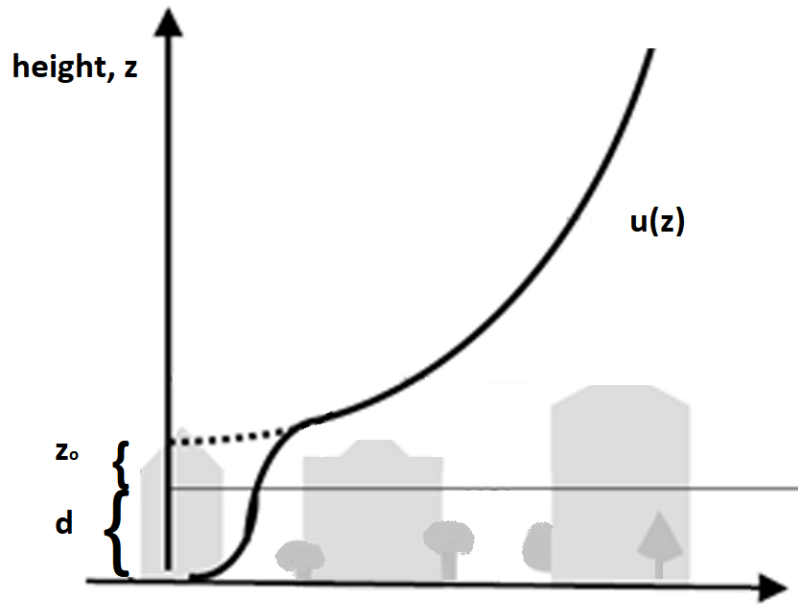


Figure 1. Profile of wind velocity, $u(z)$, in a developed urban area. The heights are the roughness length (z_0) and the zero-plane displacement length (d).

Figure 2

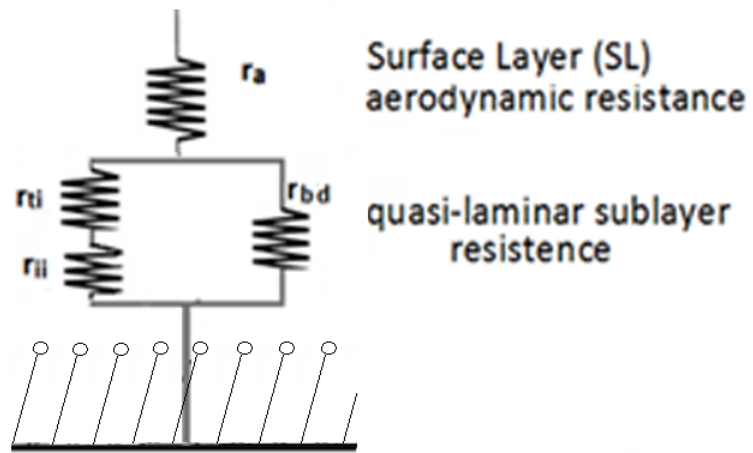


Figure 2. Electrical schematization for parametrization of particles deposition velocity reported in (Giardina et al., 2017b).

Figure 3

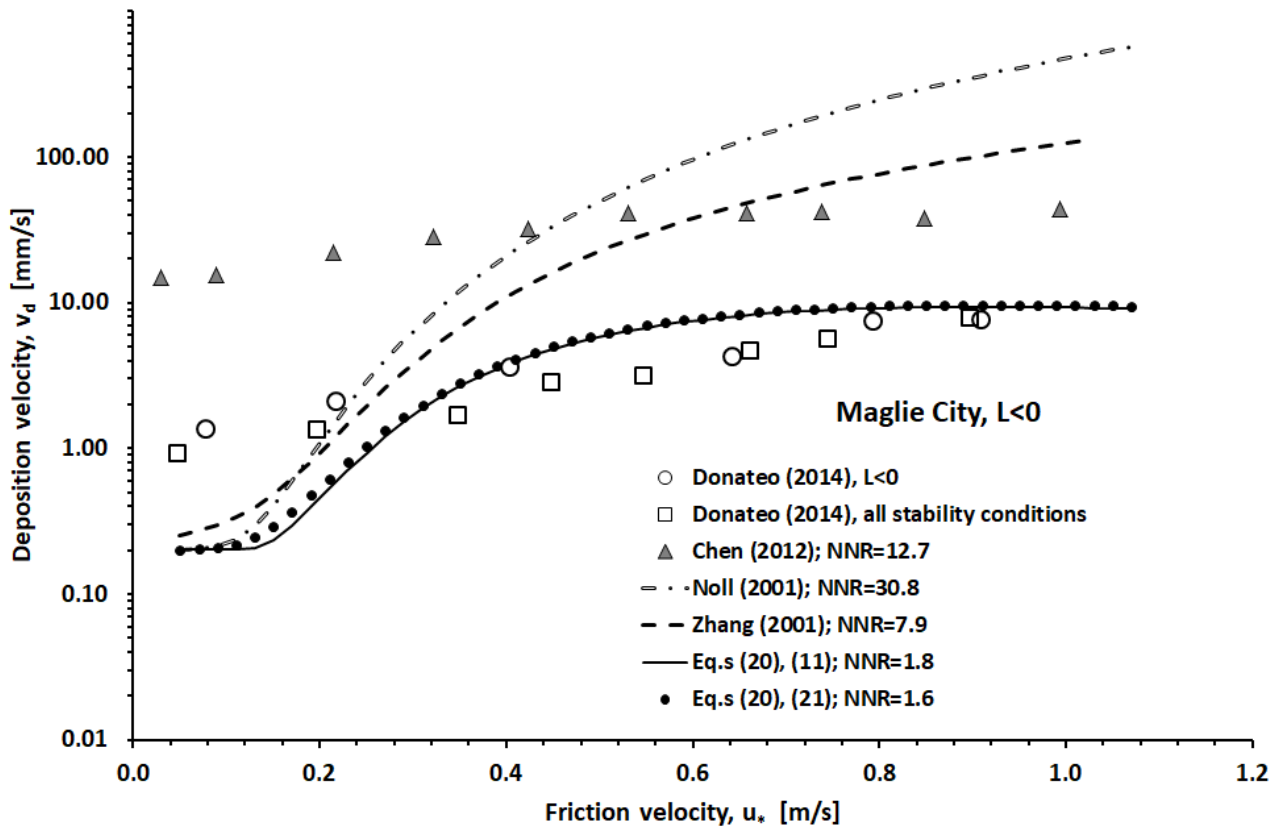


Figure 3. Comparisons among dry deposition velocity experimental data reported in (Donateo and Contini, 2014) for Maglie city, with $L < 0$ and all stability conditions, and predictions obtained by using models reported in (Giardina et al., 2017b), (Chen et al., 2012), (Noll et al., 2001), and (Zhang et al., 2001) for $d_p = 2.5 \mu\text{m}$.

Figure 4

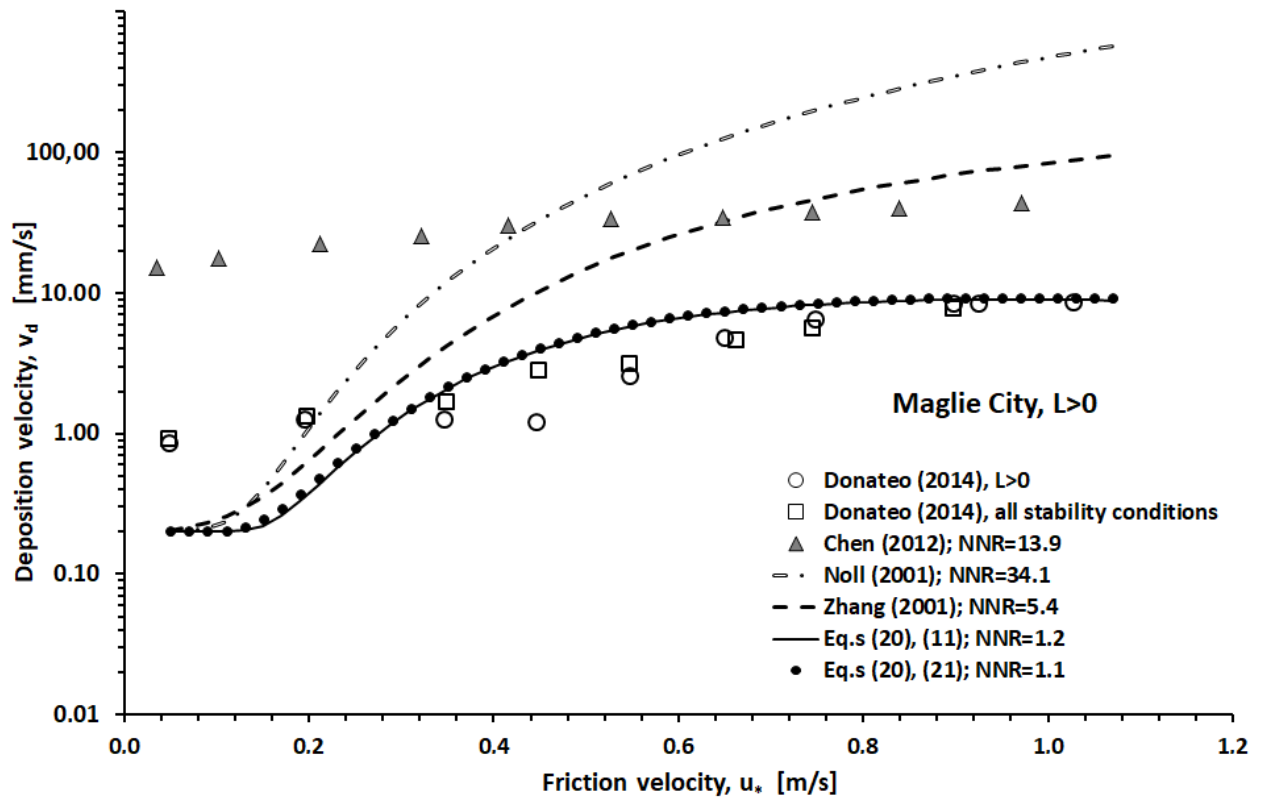


Figure 4. Comparisons among dry deposition velocity experimental data reported in (Donateo and Contini, 2014) for Maglie city, with $L > 0$ and all stability conditions, and predictions obtained by using models reported in (Giardina et al., 2017b), (Chen et al., 2012), (Noll et al., 2001), and (Zhang et al., 2001) for $d_p = 2.5 \mu\text{m}$.

Figure 5

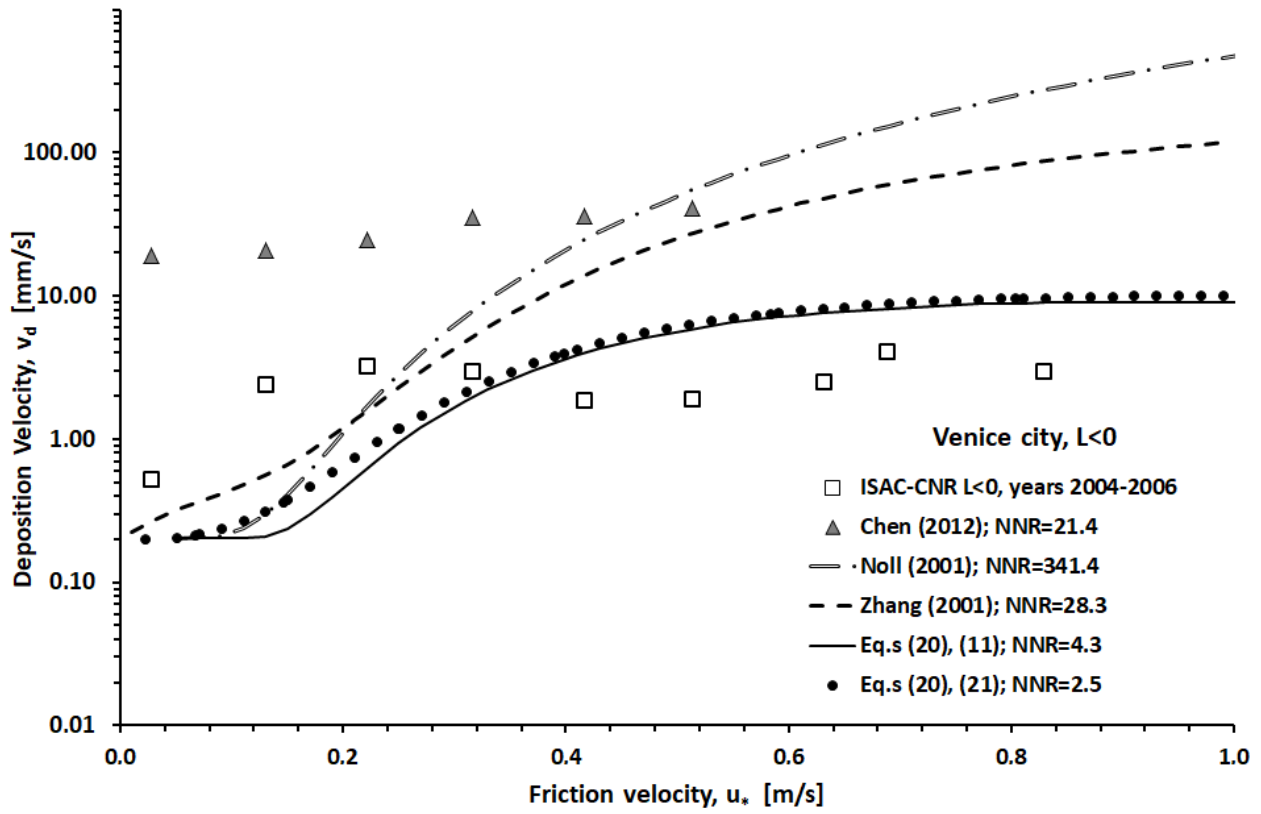


Figure 5. Comparisons between dry deposition velocity experimental data measured by ISAC-CNR for Venice, with $L < 0$, and predictions obtained by using models reported in (Giardina et al., 2017b), (Noll et al., 2001), and (Zhang et al., 2001) for $d_p = 2.5 \mu\text{m}$.

Figure 6

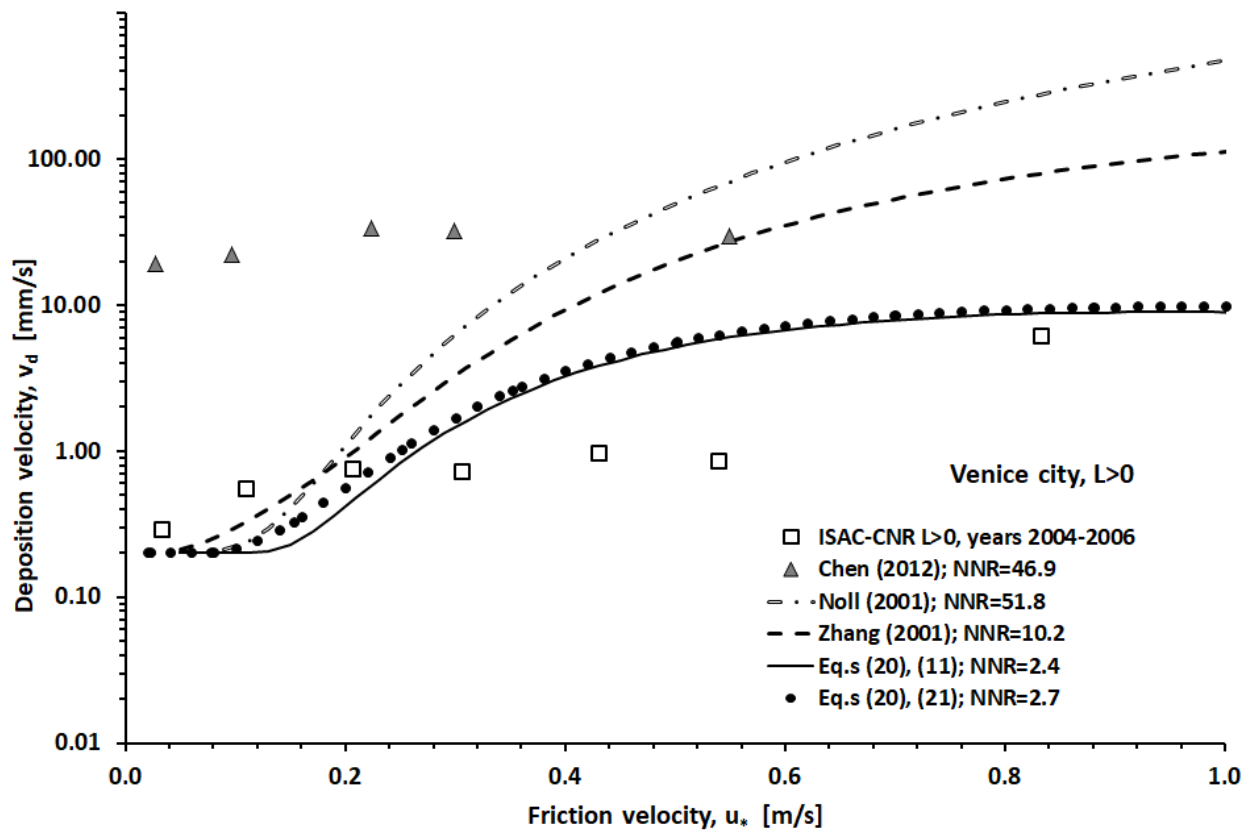


Figure 6. Comparisons between dry deposition velocity experimental data measured by ISAC-CNR for Venice, with $L>0$, and predictions obtained by using models reported in (Giardina et al., 2017b), (Noll et al., 2001), and (Zhang et al., 2001) for $d_p = 2.5 \mu\text{m}$.

Figure 7

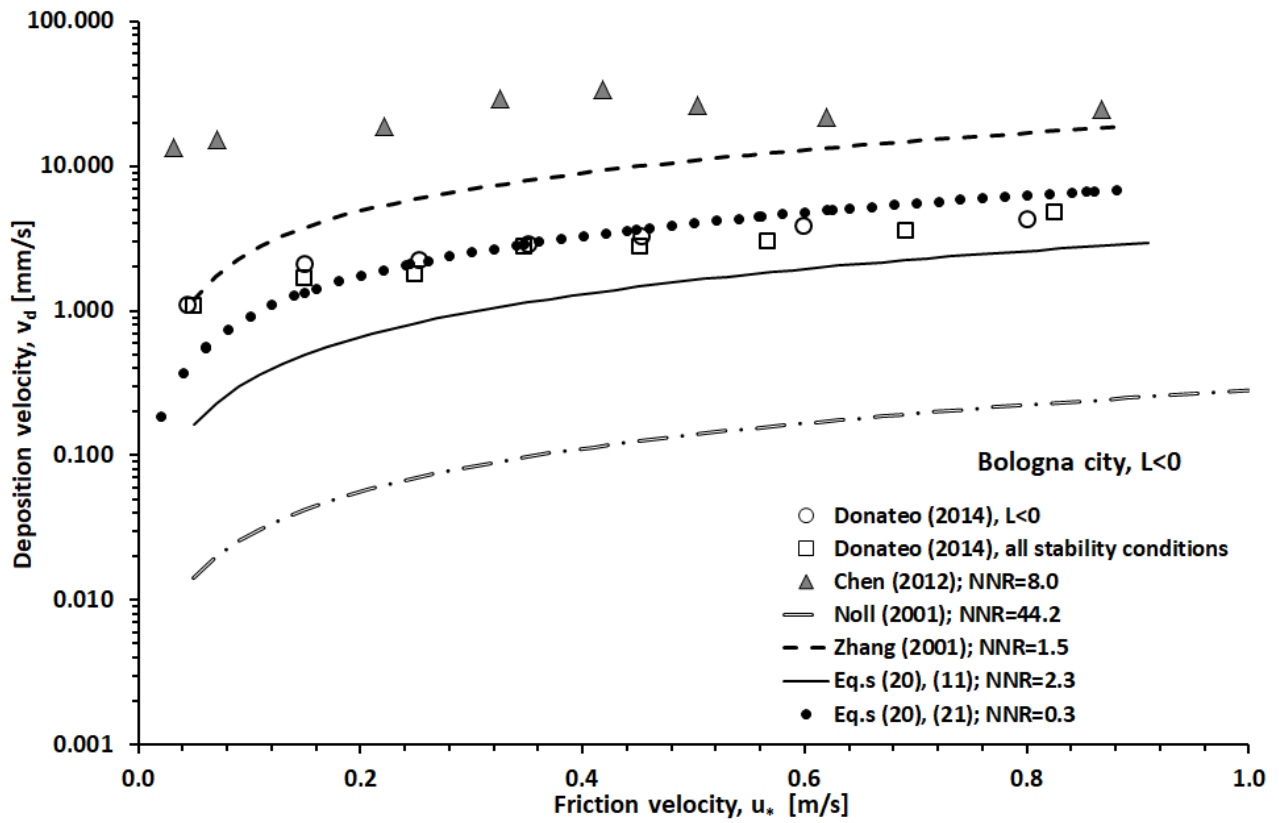


Figure 7. Comparisons among dry deposition velocity experimental data reported in (Donateo and Contini, 2014) for Bologna city outskirts, with $L < 0$ and all stability conditions, and predictions obtained by using models reported in (Giardina et al., 2017b), (Chen et al., 2012), (Noll et al., 2001), and (Zhang et al., 2001) for $d_p = 0.045 \mu\text{m}$.

Figure 8

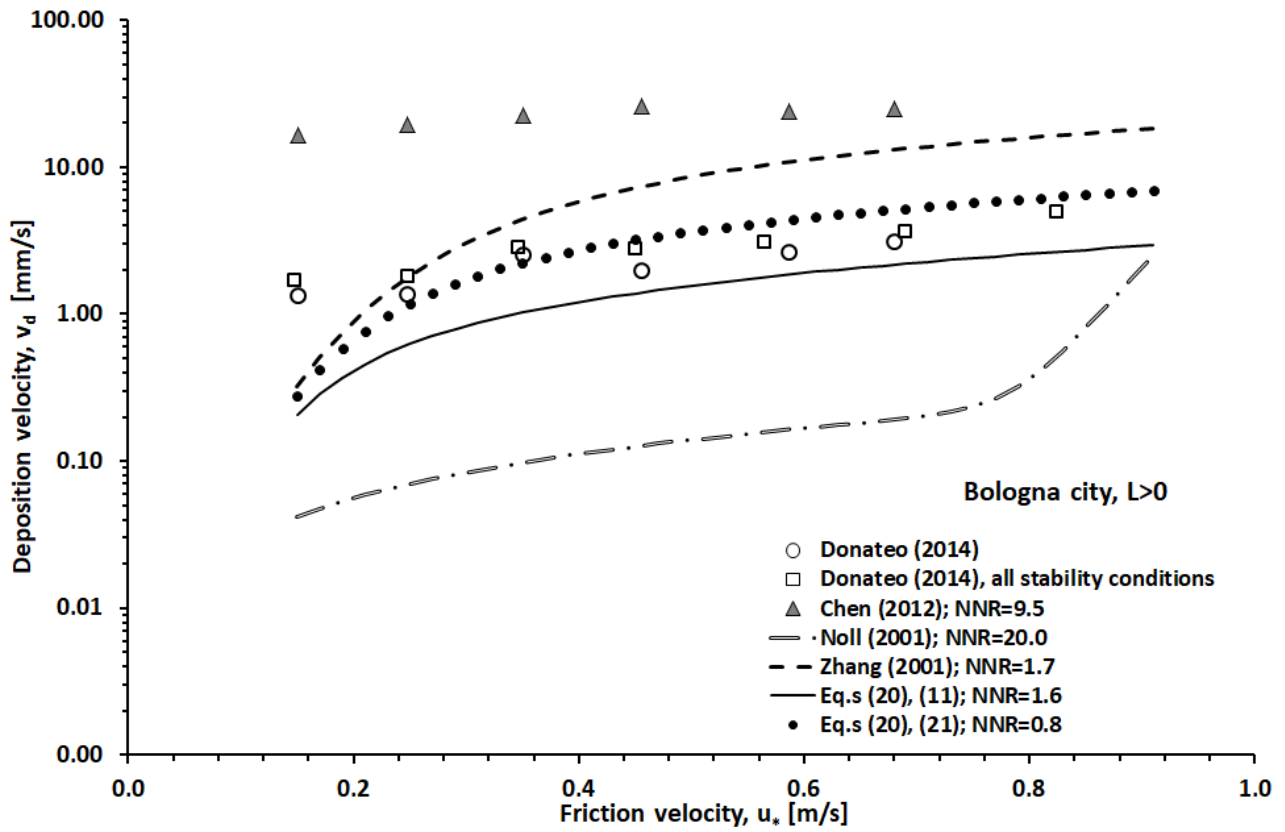


Figure 8. Comparisons among dry deposition velocity experimental data reported in (Donateo and Contini, 2014) for Bologna city outskirts, with $L > 0$ and all stability conditions, and predictions obtained by using models reported in (Giardina et al., 2017b), (Chen et al., 2012), (Noll et al., 2001), and (Zhang et al., 2001) for $d_p = 0.045 \mu\text{m}$.

Figure 9

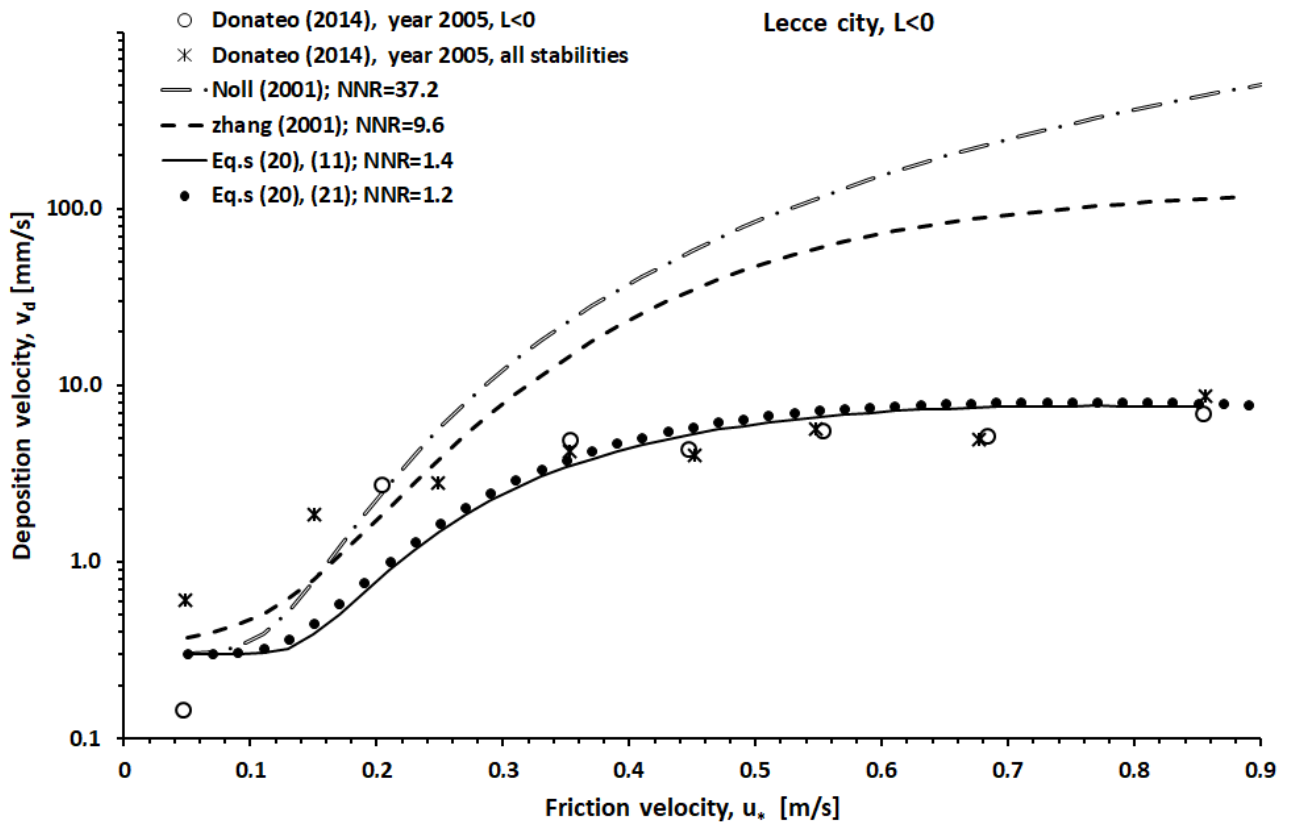


Figure 9. Comparisons among dry deposition velocity experimental data reported in (Donateo and Contini, 2014) for Lecce suburbs city, with $L<0$ and all stability conditions, and predictions obtained by using models reported in (Giardina et al., 2017b), (Noll et al., 2001), and (Zhang et al., 2001) for $d_p = 2.5 \mu\text{m}$.

Figure 10

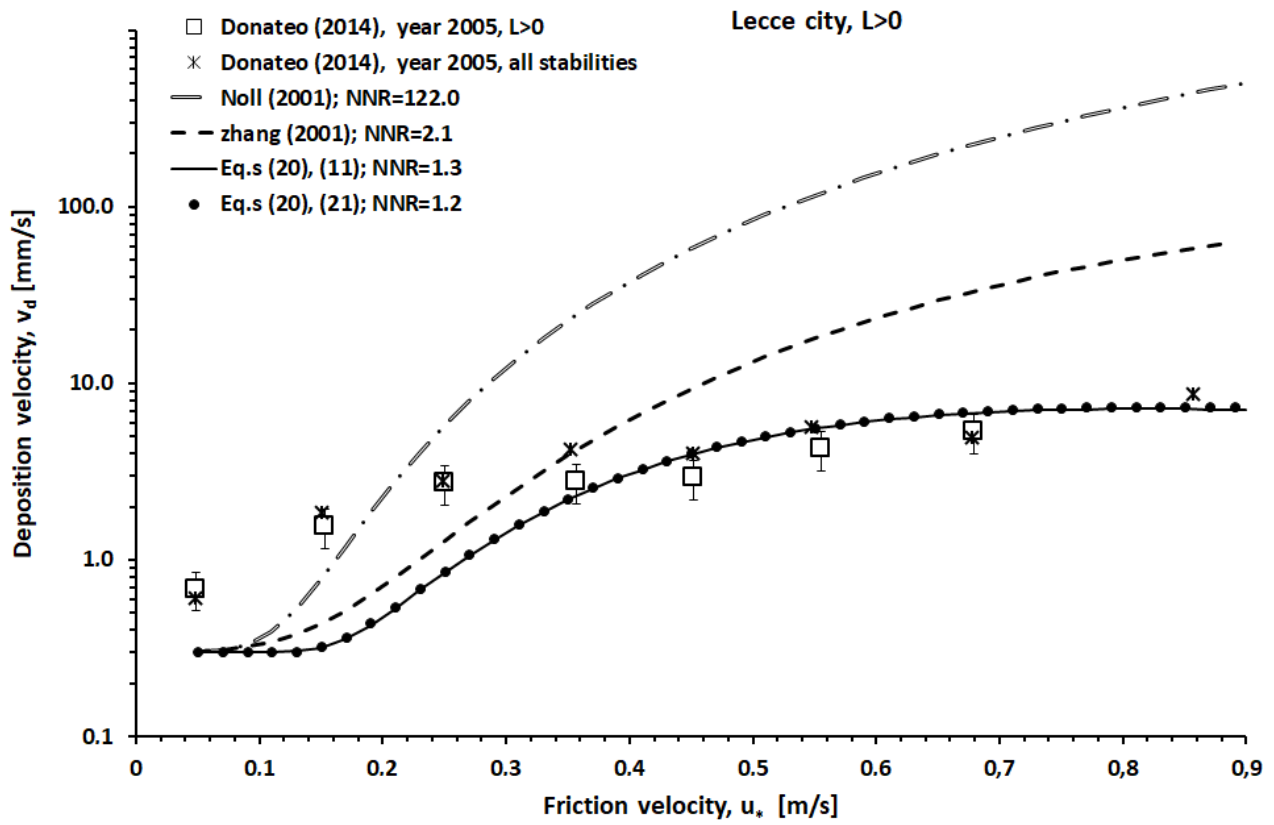


Figure 10. Comparisons among dry deposition velocity experimental data reported in (Donateo and Contini, 2014) for Lecce suburbs city, with $L>0$ and all stability conditions, and predictions obtained by using models reported in (Giardina et al., 2017b), (Noll et al., 2001), and (Zhang et al., 2001) for $d_p = 2.5 \mu\text{m}$.

Figure 11

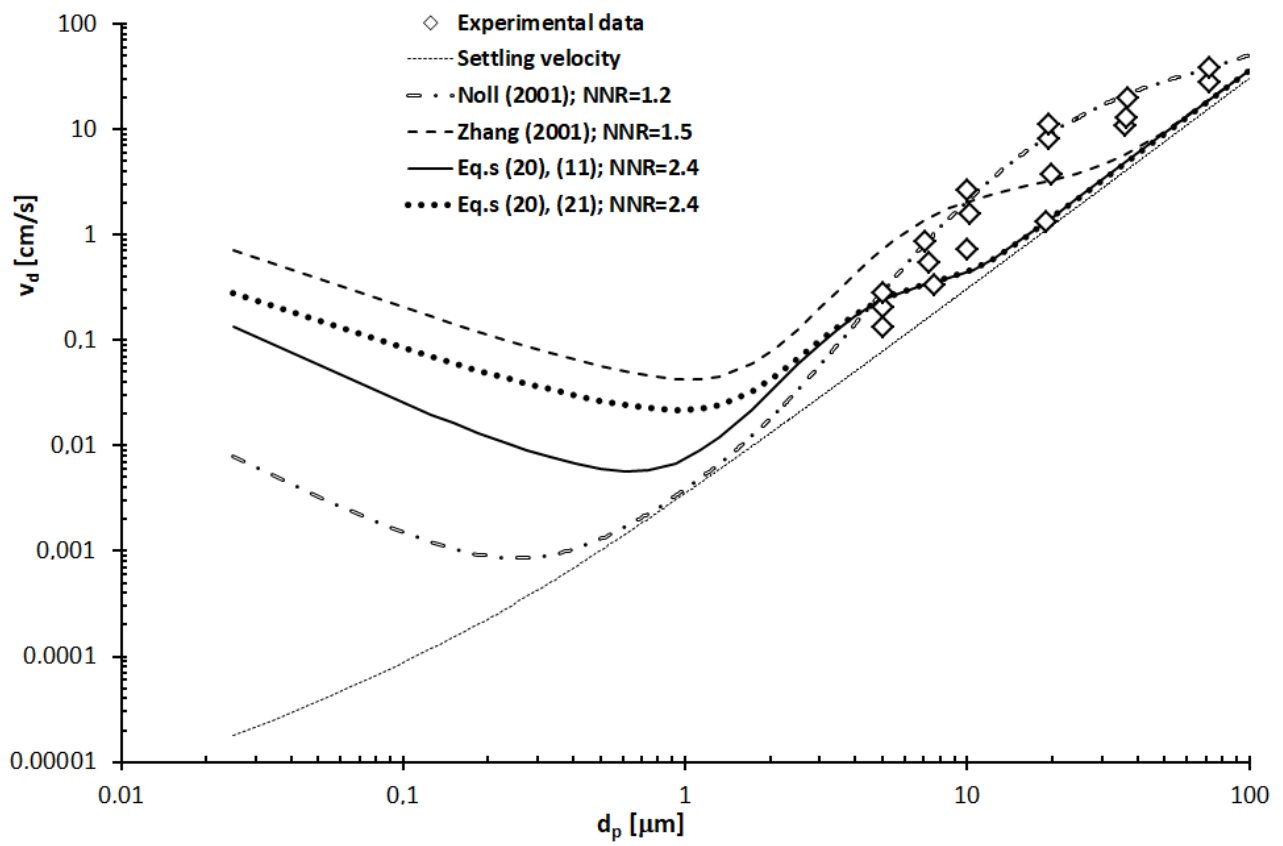


Figure 11. Comparisons between experimental dry deposition velocities reported in (Noll et al., 2001) as function of particle diameters and predictions of models reported in (Giardina et al., 2017b), (Noll et al., 2001) and (Zhang et al., 2001).

Figure 12

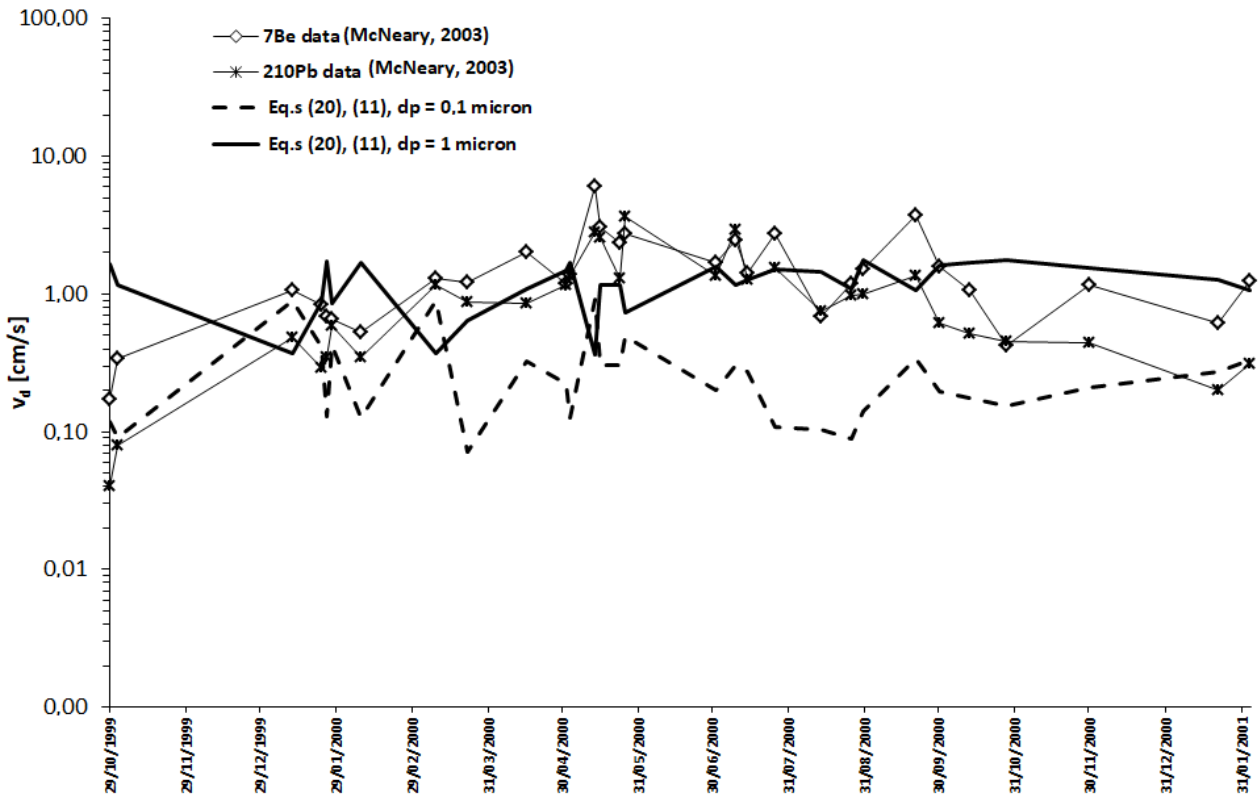


Figure 12. Comparisons between dry deposition experimental data reported in (McNeary and Baskaran, 2003) and predictions of the model proposed in (Giardina et al., 2017b) by using Brownian diffusion resistance Eq. (11) for particle diameters d_p 0.1 and 1 μm

Figure 13

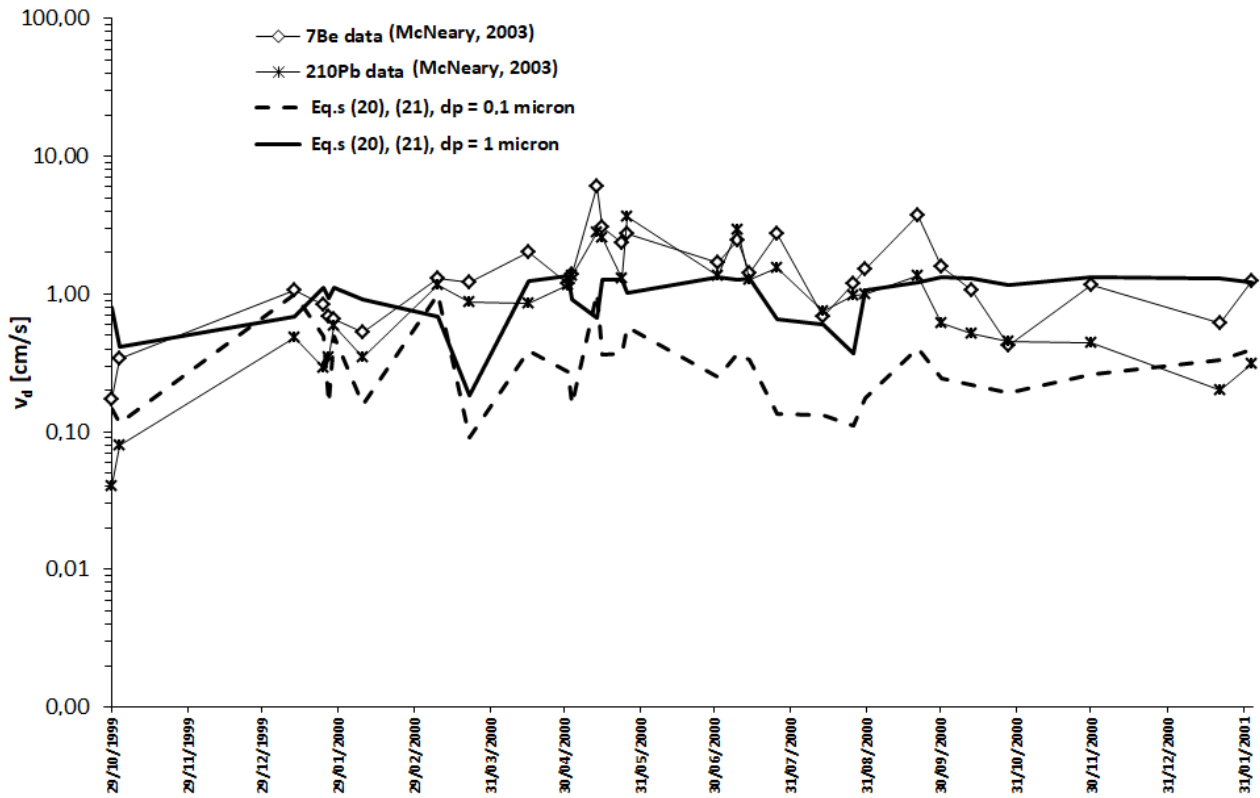


Figure 13. Comparisons between dry deposition experimental data reported in (McNeary and Baskaran, 2003) and predictions of the model proposed in (Giardina et al., 2017b) by using Brownian diffusion resistance Eq. (21) and particle diameters d_p 0.1 and 1 μm .

Figure 14

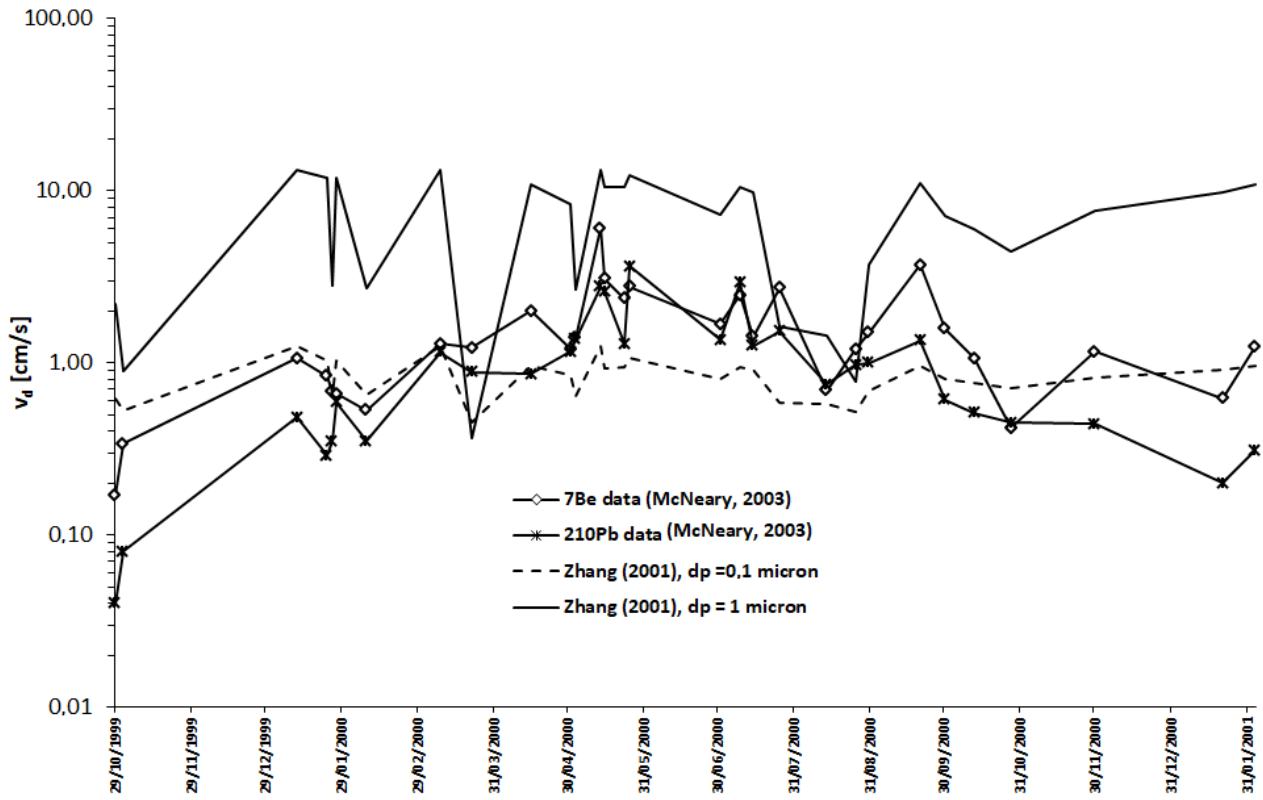


Figure 14. Comparisons between dry deposition experimental data reported in (McNeary and Baskaran, 2003) and predictions of the model proposed in (Zhang et al., 2001) for particle diameters d_p 0.1 and 1 μm .

Figure 15

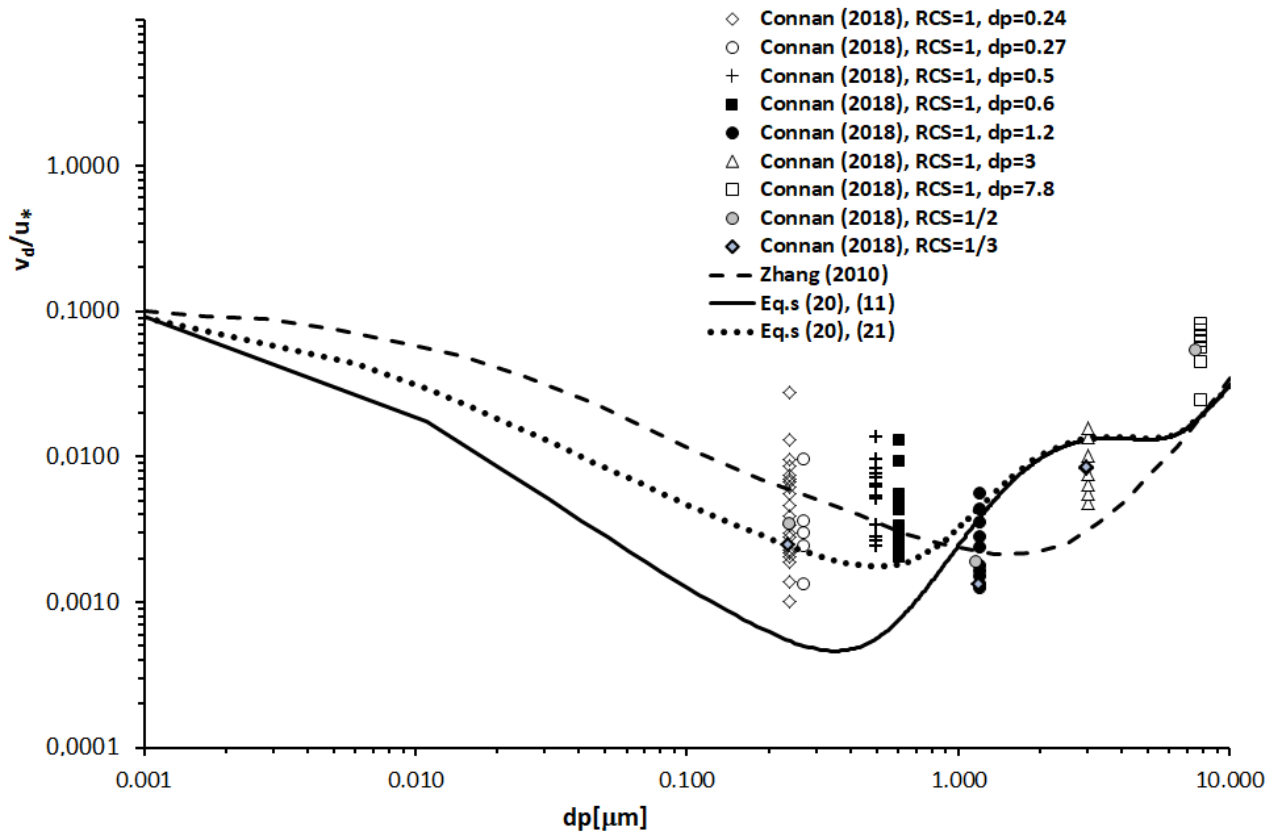


Figure 15. Comparisons between experimental data reported in (Connan et al., 2018) and predictions obtained by using (Giardina et al., 2017b) and (Zhang et al., 2001) models.